

# Probing low energy quantum matter

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UNIVERSITY OF  
**BATH**

**E. DaComo**

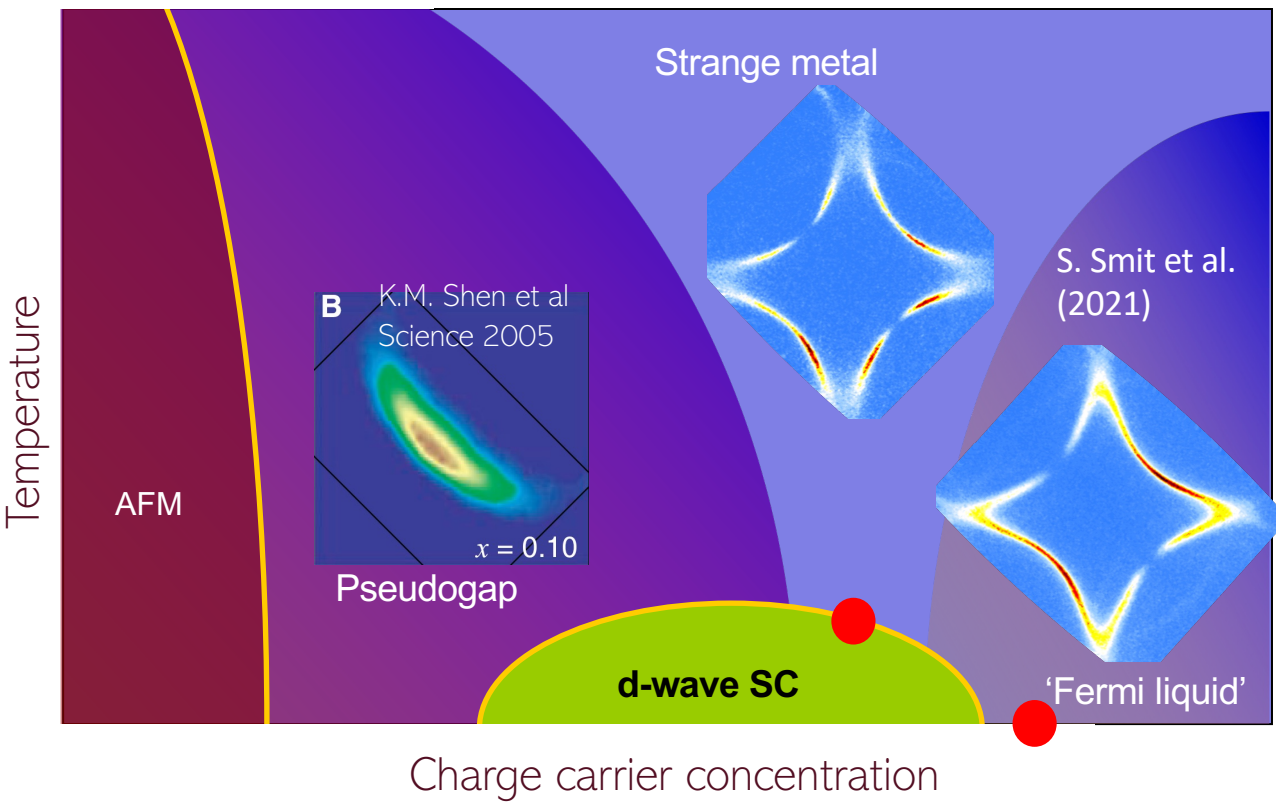
Collaborators

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Quantum  
Materials  
Amsterdam

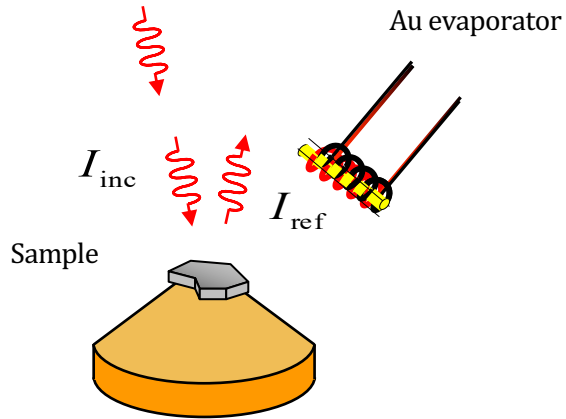
# From Fermi liquid to strange metal



Key question:  
What distinguishes the strange metal phase from a Fermi liquid?

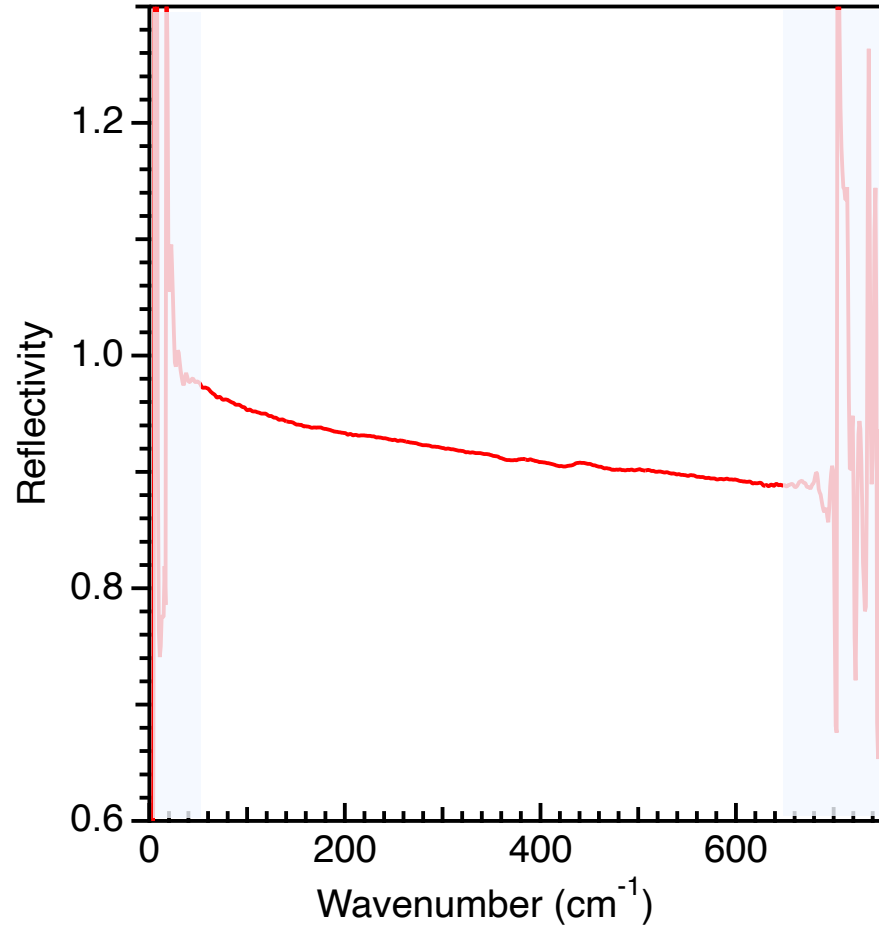
Optical experiments:  
Frequency dependence of the conductivity is a valuable source of information on collective response.

# How do you measure the optical response?



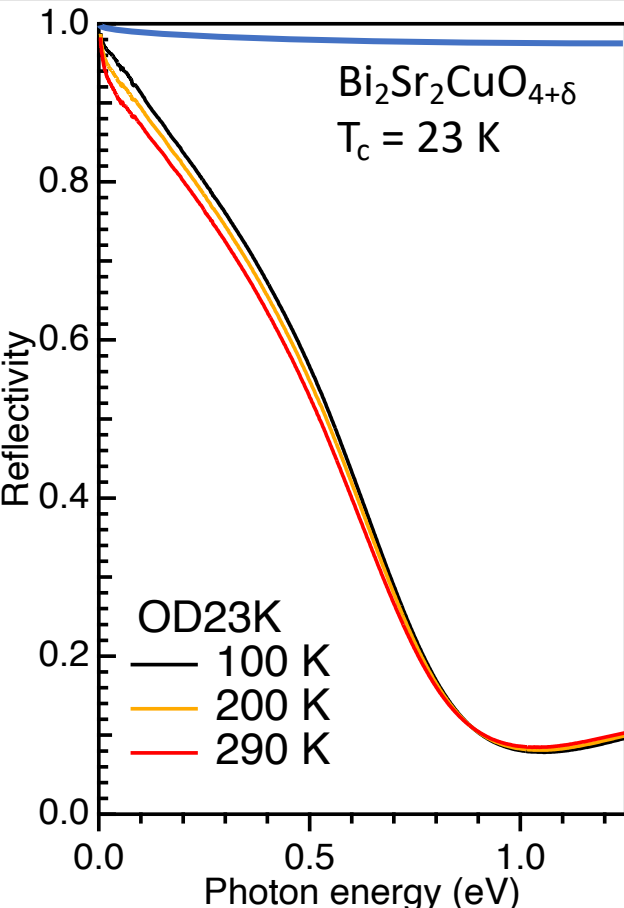
Measures  $R(\omega)$

$$R(\omega) = \frac{I_{sample}}{I_{ref.}}$$

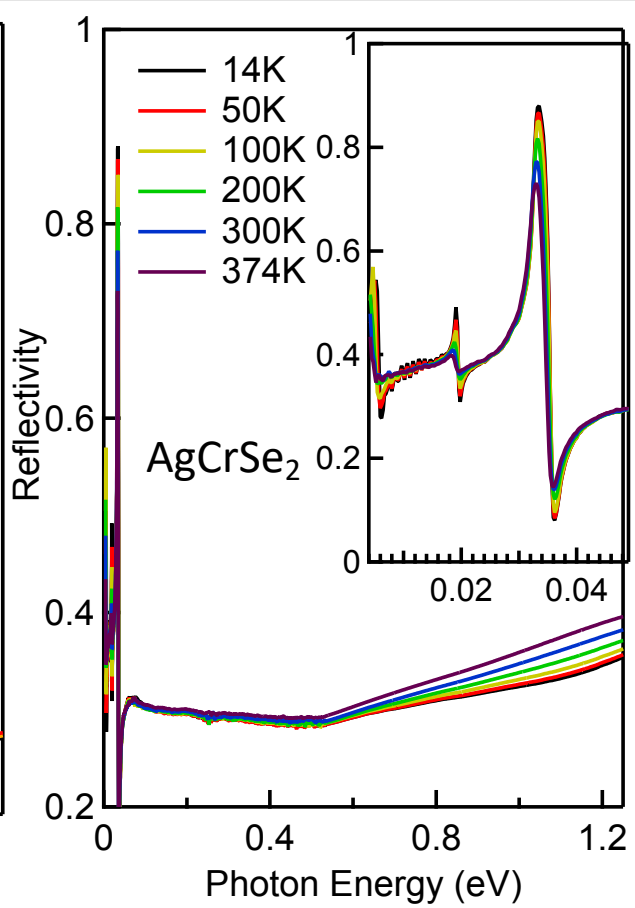


# Metals, insulators and spontaneous symmetry breaking

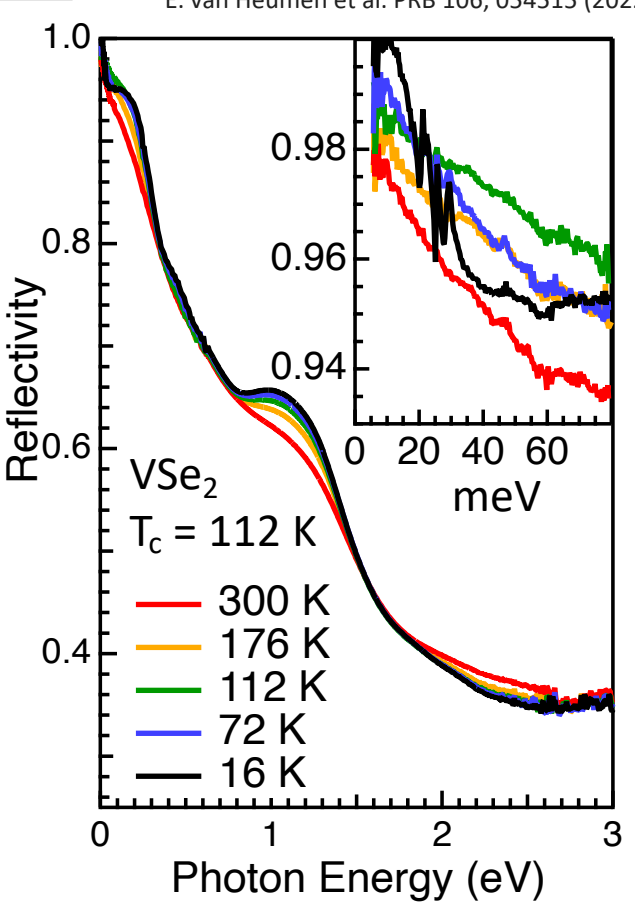
X. Feng et al. PRB 104, 165134 (2021)  
J. Groefsema et al. PRM 6, 115402 (2022)  
E. van Heumen et al. PRB 106, 054515 (2022)



Bad/strange Metal



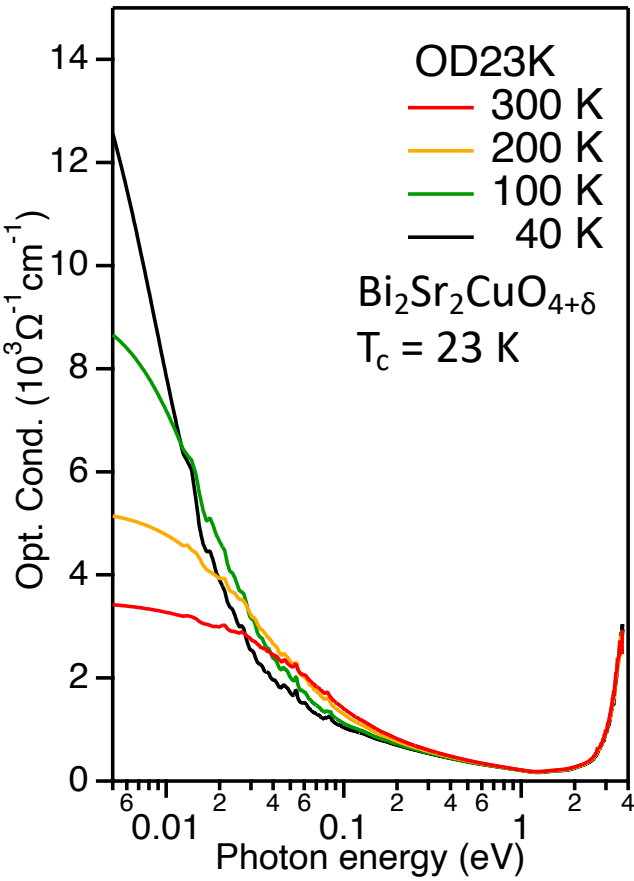
Insulator



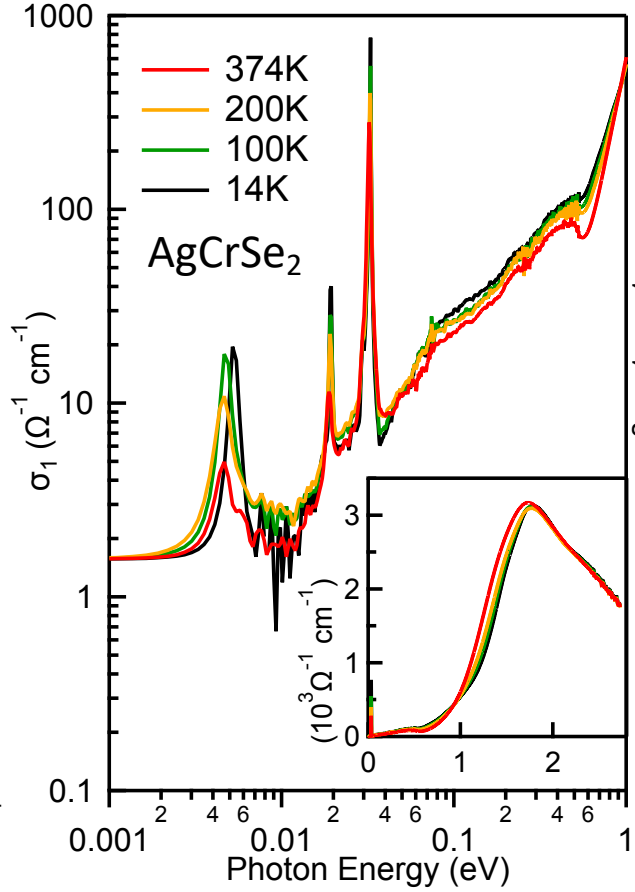
Charge density wave

# Metals, insulators and spontaneous symmetry breaking

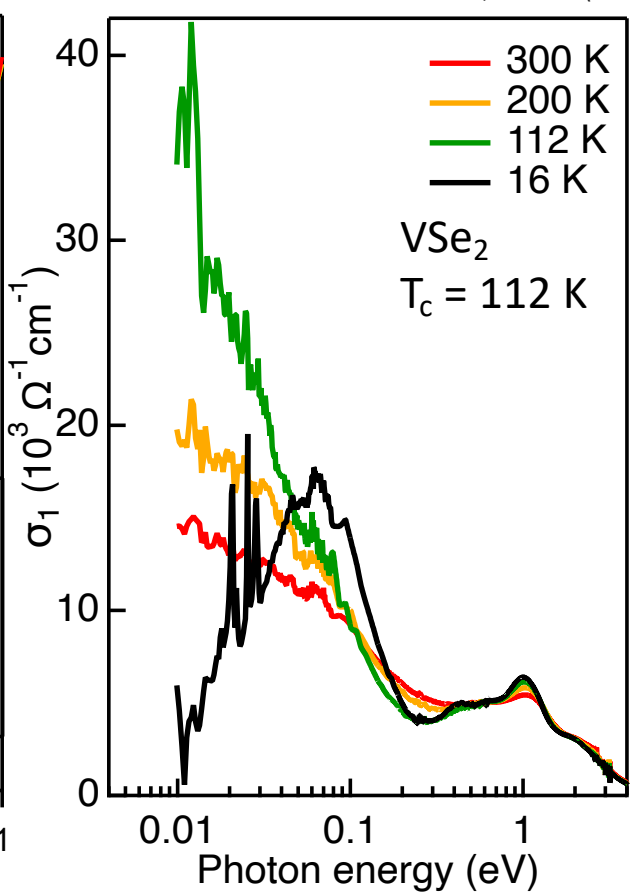
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Bad/strange Metal



Insulator

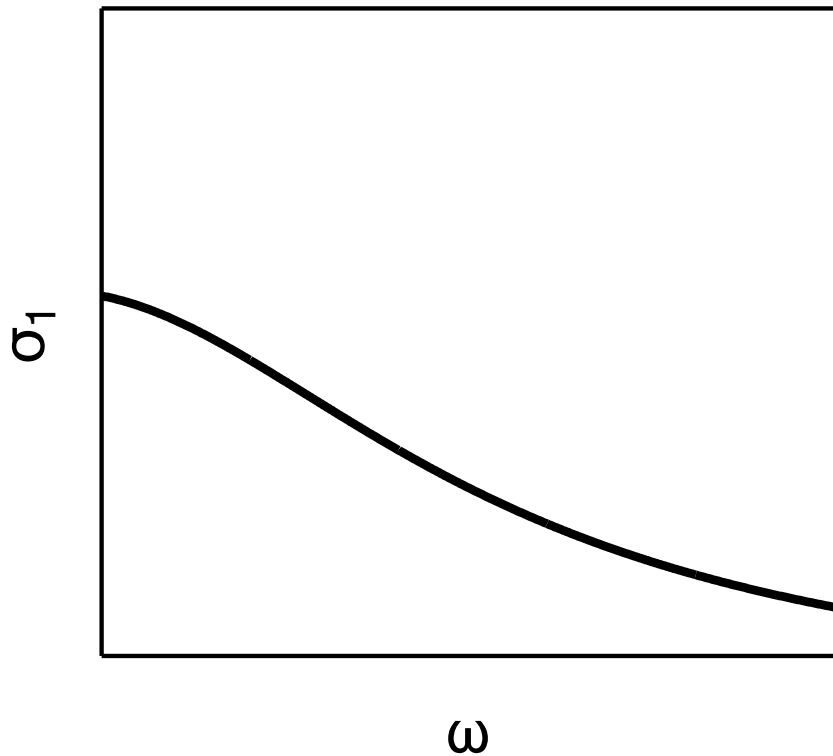
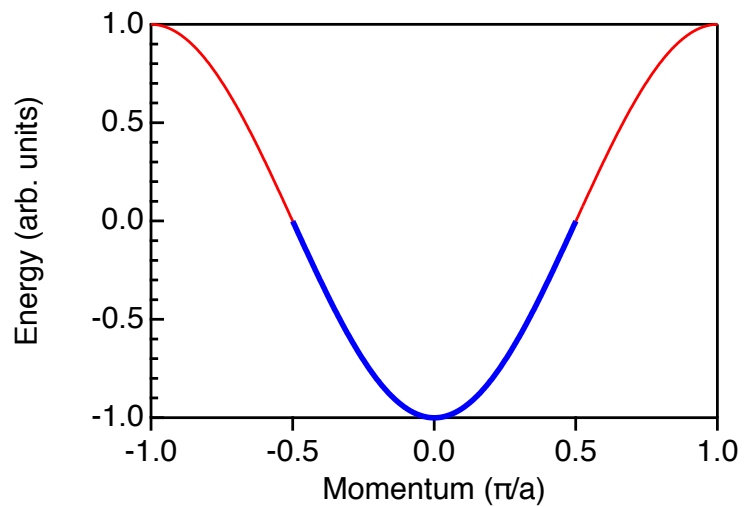
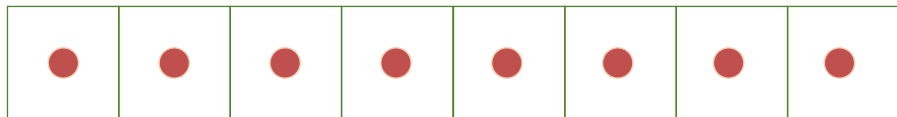


Charge density wave

# Broken symmetry states

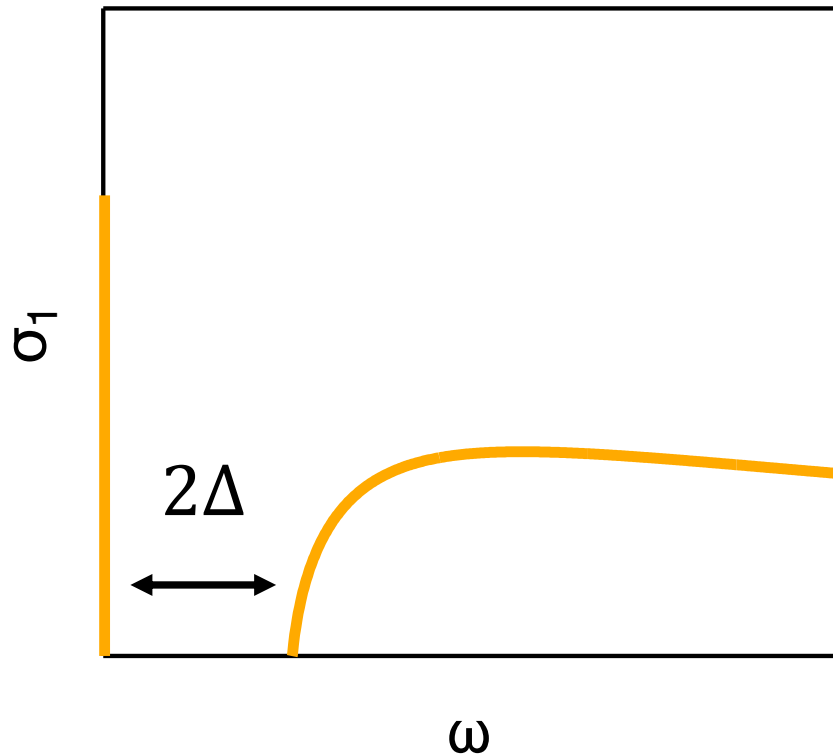
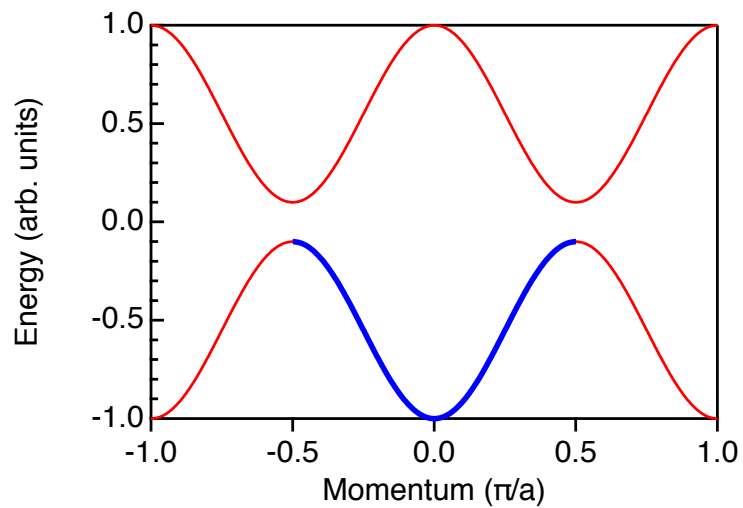
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## Naive picture of CDW formation



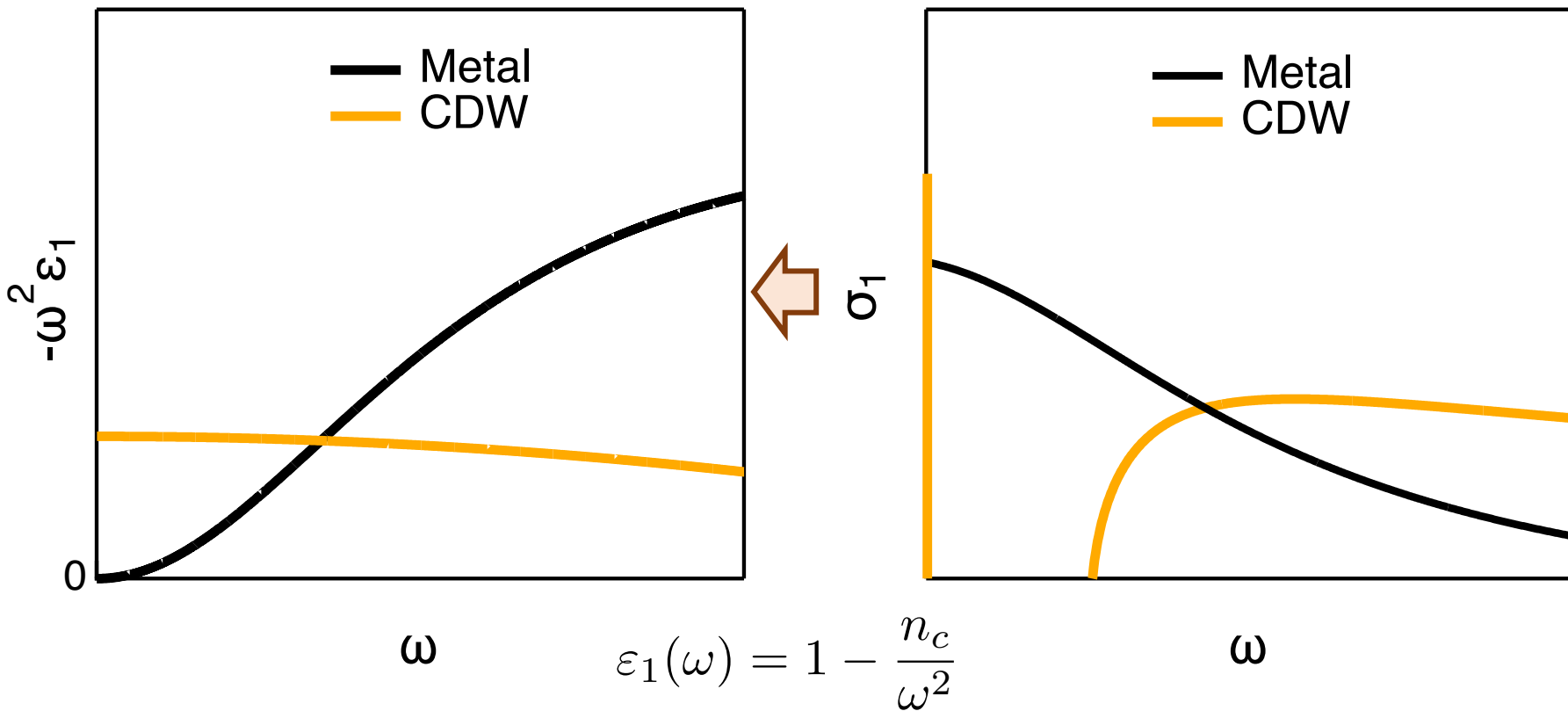
# Broken symmetry states

## Naive picture of CDW formation



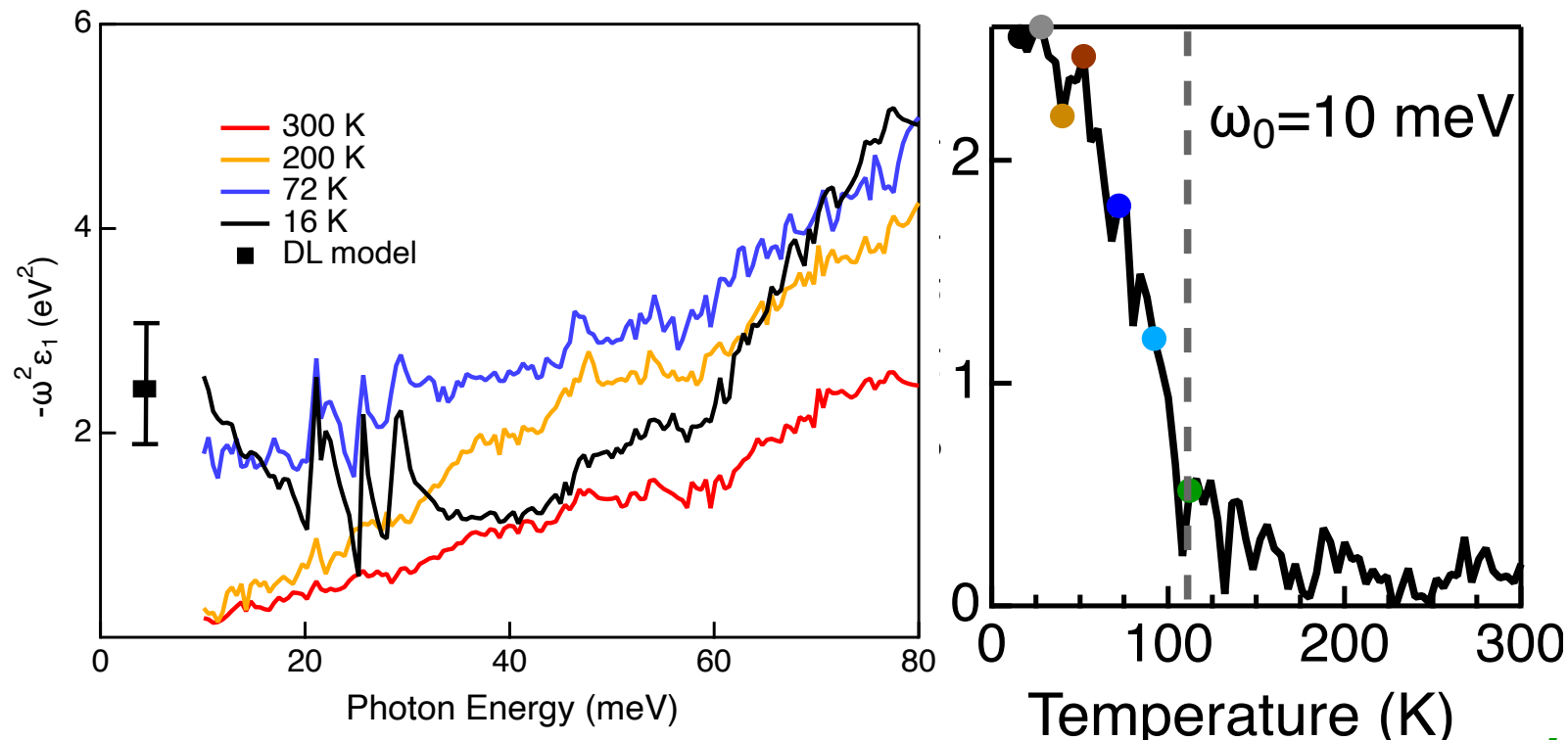
# Probing the condensate through causality

Kramers Kronig transformation of  $\delta$ -function response:

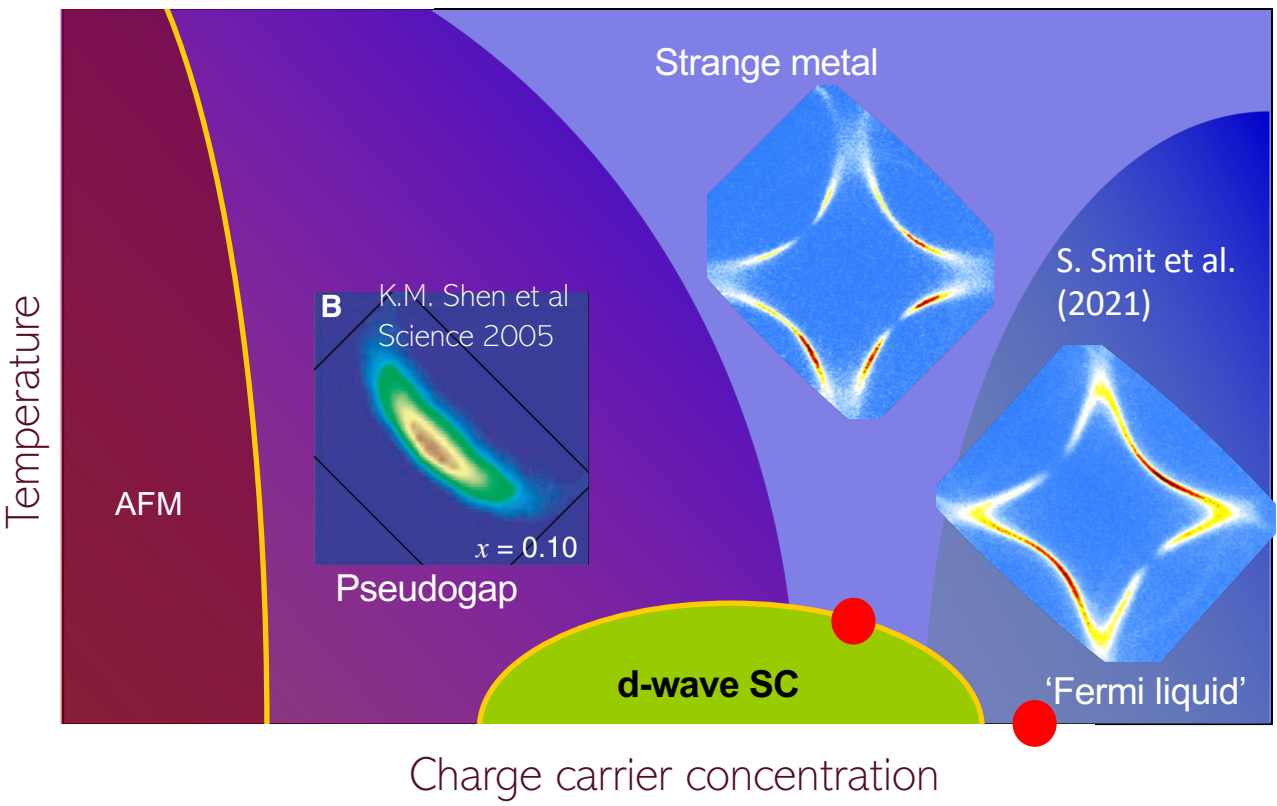




# The condensate density in $VSe_2$



# From Fermi liquid to strange metal

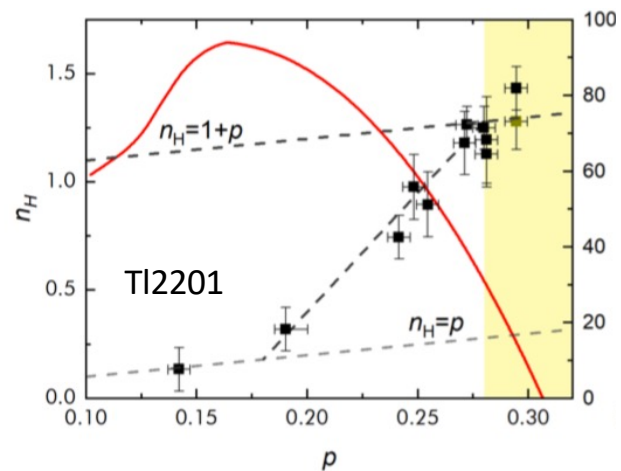
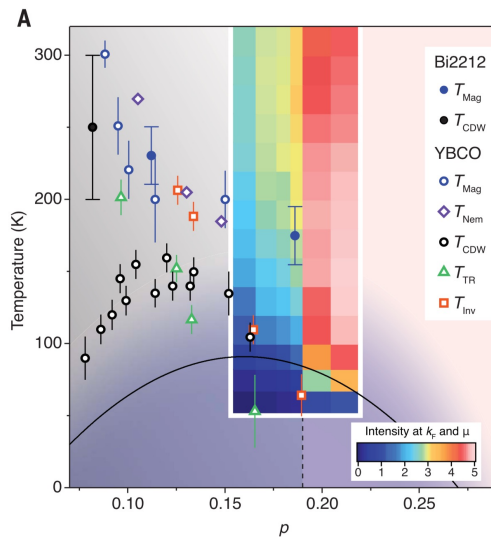
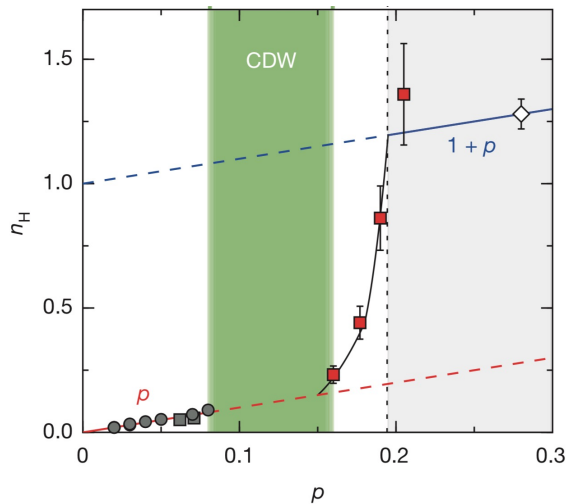


Key question:  
What distinguishes the strange metal phase from a Fermi liquid?

Optical experiments:  
Frequency dependence of the conductivity is a valuable source of information on collective response.

# First Question

- Is there a hint of a critical doping near 0.2 holes/Cu in optics?



S.-D. Chen et al., Science **366**, 1099 (2019).

Badoux et al. Nature **531**, 210 (2016)

(also Legros Nature Phys **15**, 142 Michon Nature **567**, 218 (2019))

Putzke et al., Nature Phys **17**, 826 (2021)

Ayres et al., Nature **595**, 661 (2021)

# cuprate superconductors: One or two components?

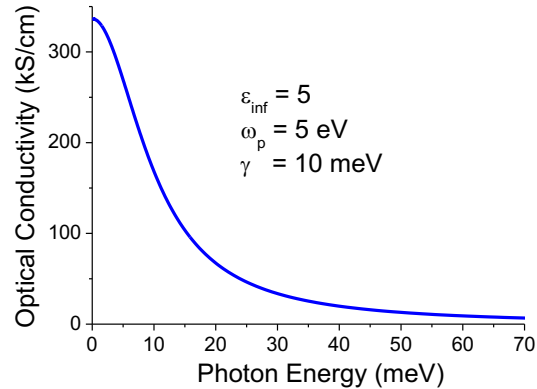
## single component

quasi-particles

residual interactions

(energy dependent) lifetime broadening

$$\sigma(\omega, T) = \frac{iD}{\omega + \hat{M}(\omega, T)}$$



## two (or more) components

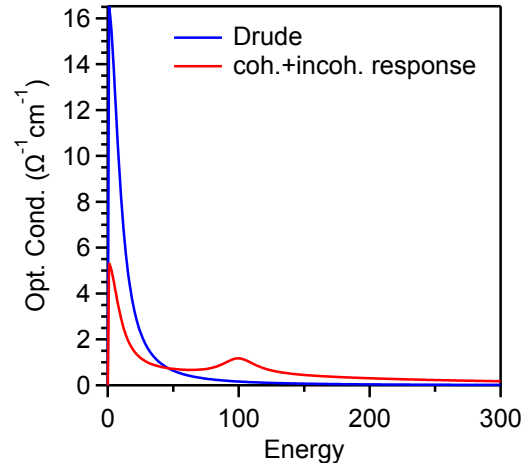
no quasi-particles

(residual) interactions

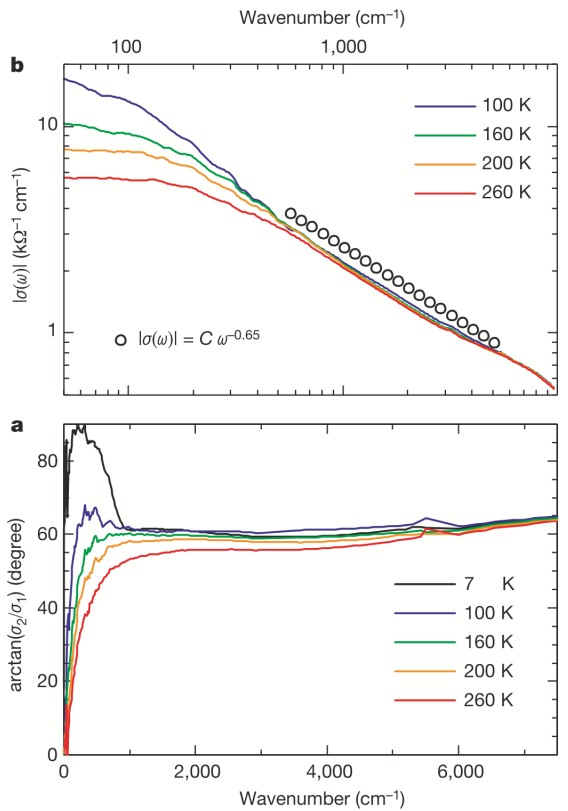
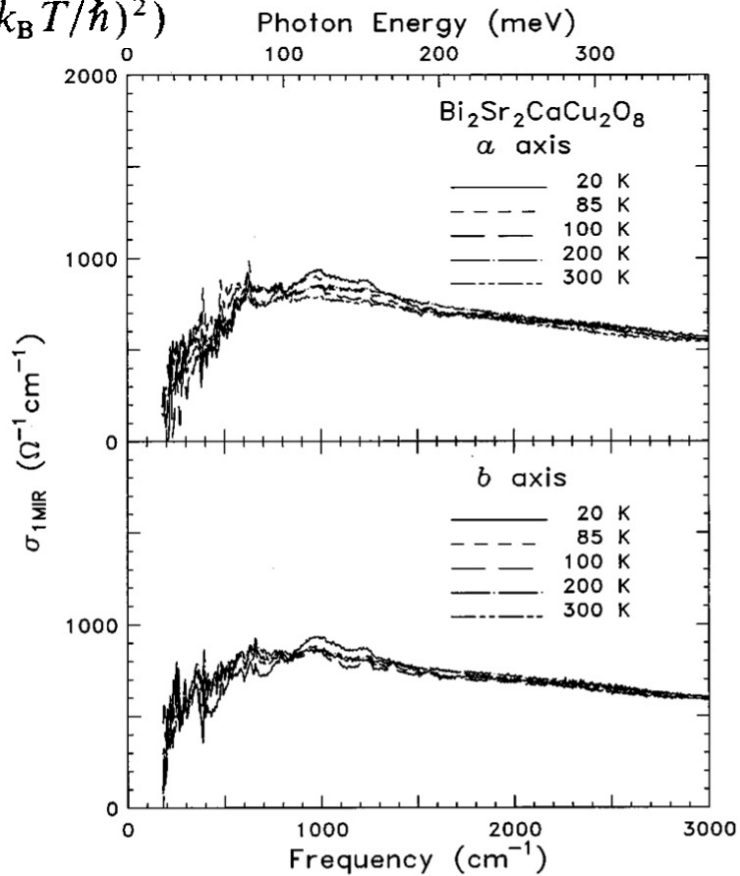
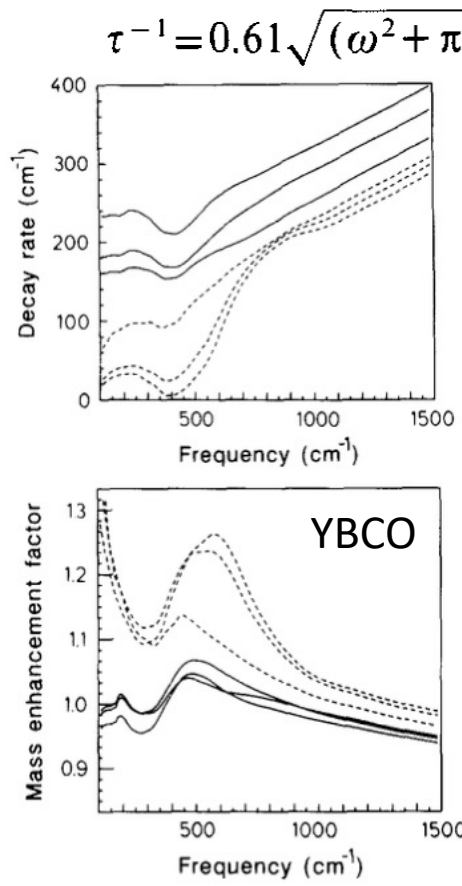
mid-infrared bands

lifetime broadening

$$\sigma(\omega, T) = \sigma_a(\omega, T) + \sigma_b(\omega, T)$$



# State of the art over the years



D. van der Marel et al., Physica C, 176 (1991)

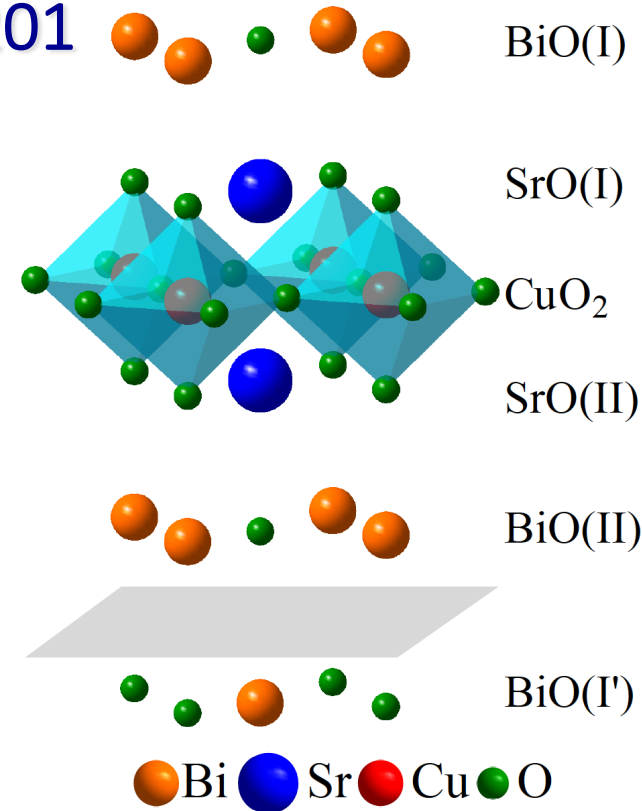
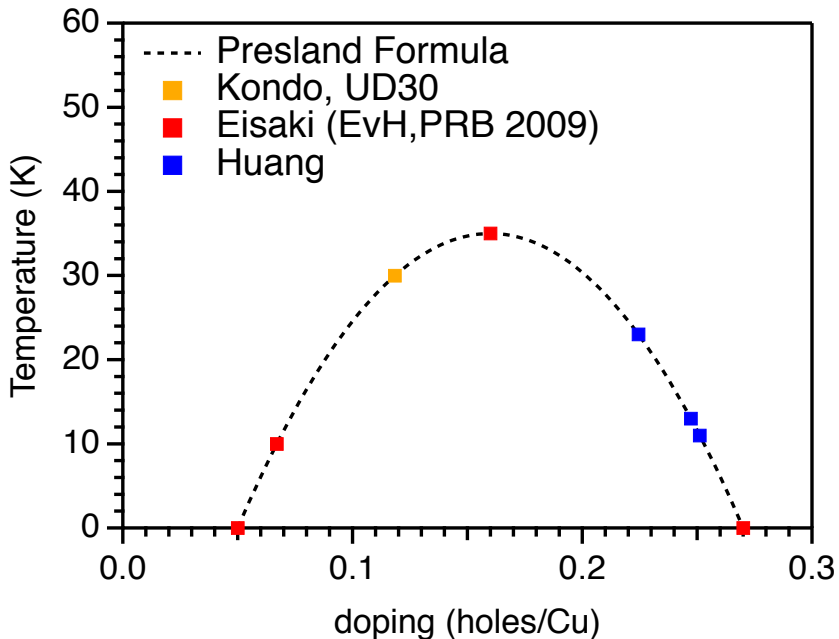
S. Uchida, et al., Phys. Rev. B **43**, 7942 (1991)

D. van der Marel, et al. Nature **425**, 271 (2003)

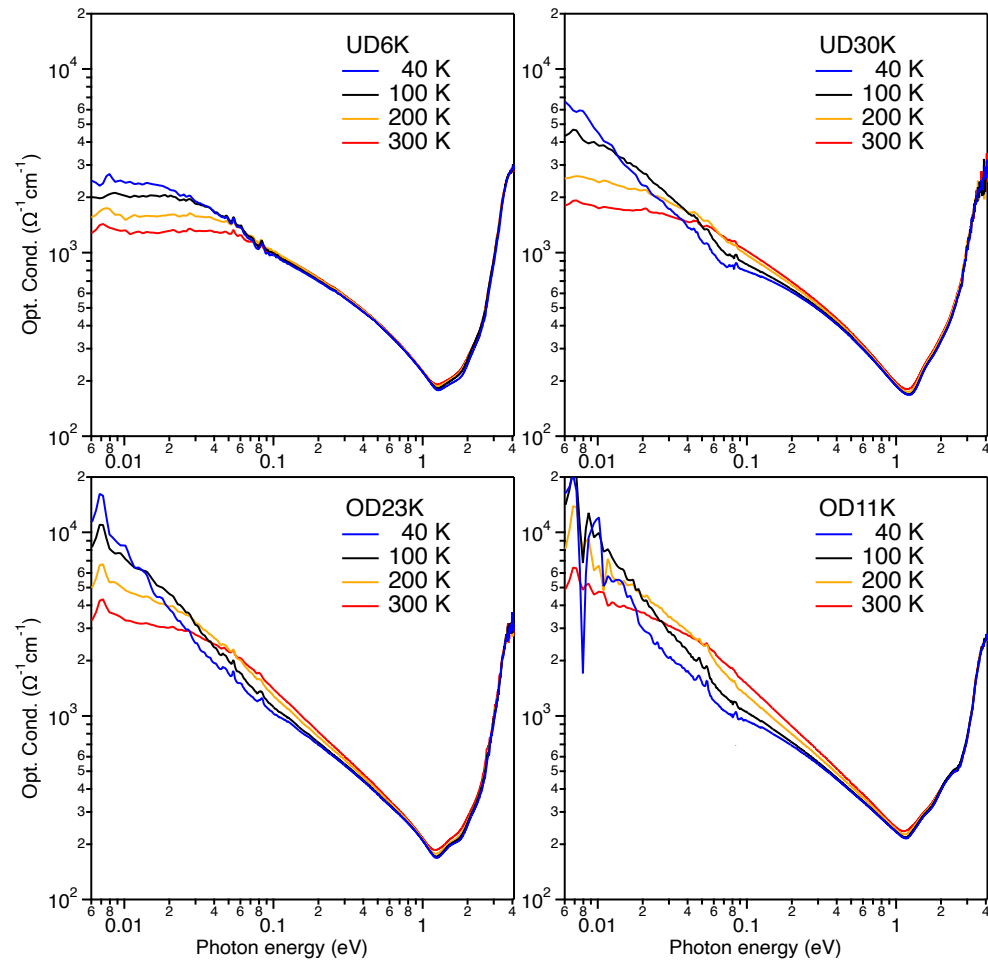
# The ideal choice: $\text{Bi}_{2-x}\text{Pb}_x\text{Sr}_{2-y}\text{La}_y\text{CuO}_{6+\delta}$

## The material of choice: single layer Bi2201

- One  $\text{CuO}_2$  layer per unit cell.
- Low critical temperature ( $T_{c,\text{max}} = 35 \text{ K}$ )
- Good cleavage surface



# The optical response of cuprates

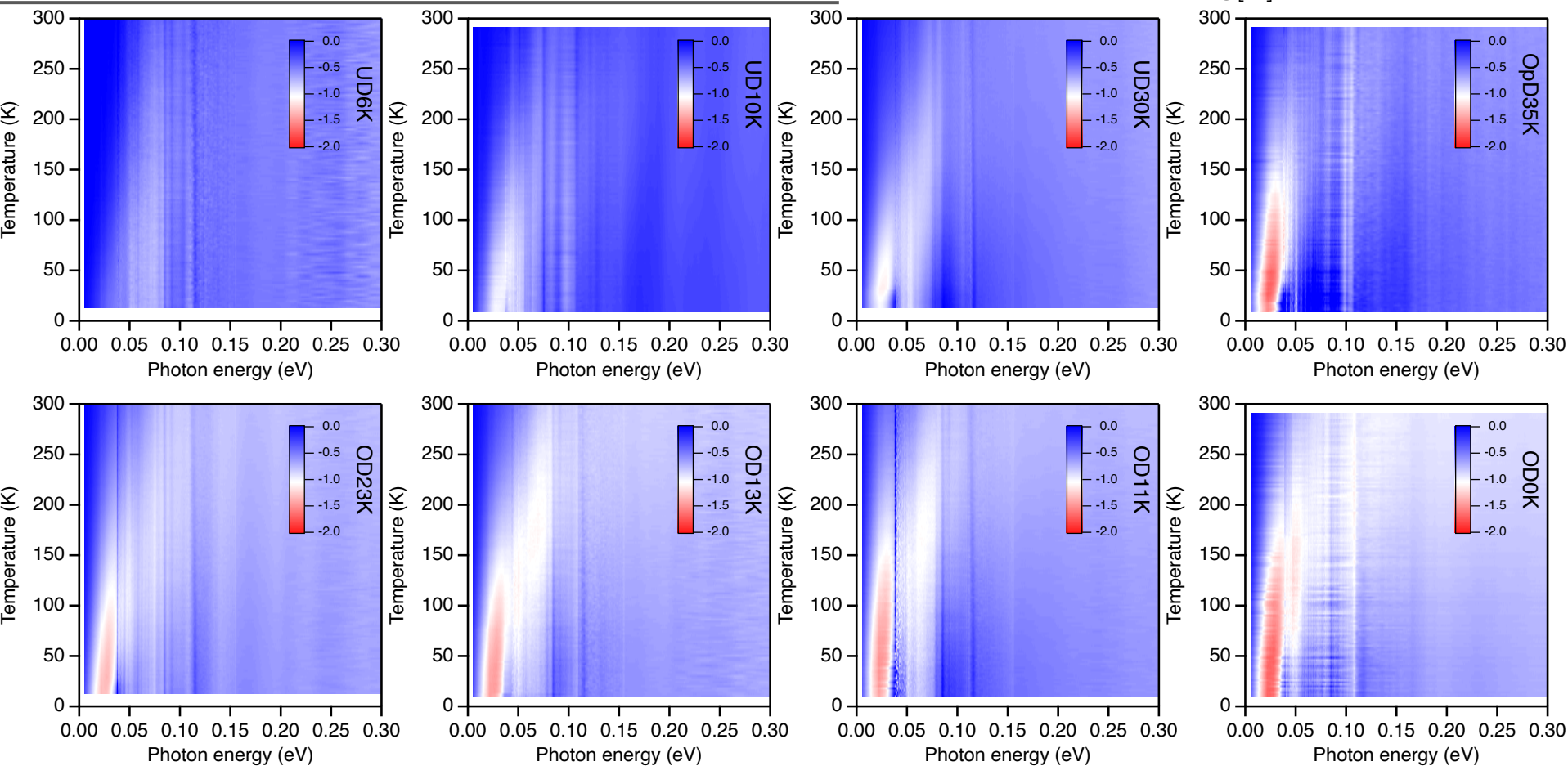


- Extrapolation to  $\omega=0$  consistent with  $\rho_{\text{DC}}$ .
- Two components in the ‘intraband’ optical response.
- Crossover regime around 0.1 eV

$$\sigma_1(\omega) = \frac{\omega_p^2}{4\pi} \frac{\Gamma}{\Gamma^2 + \omega^2}$$

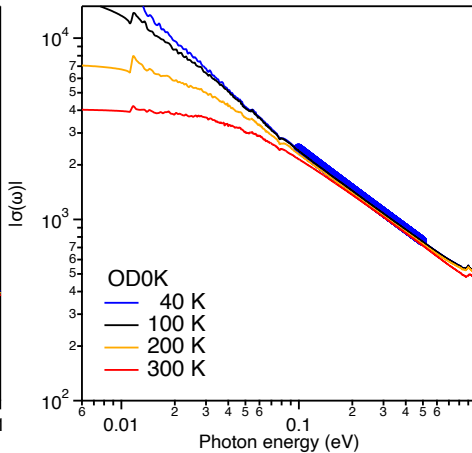
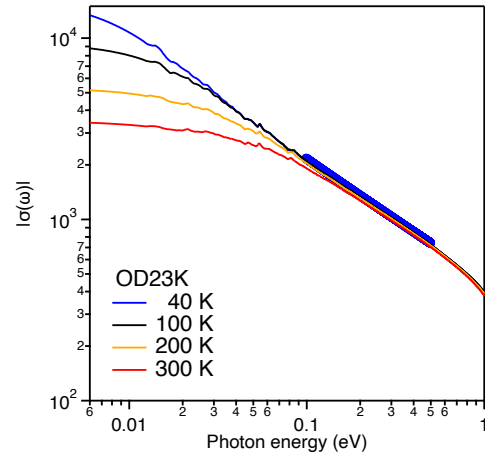
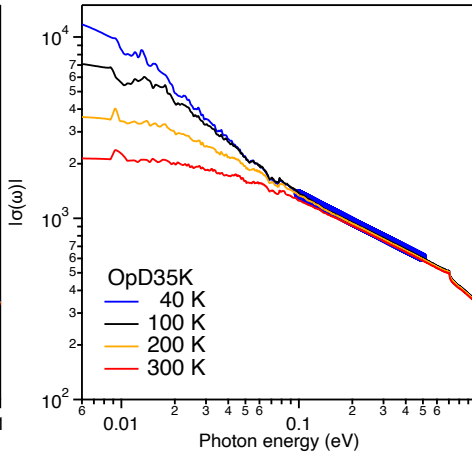
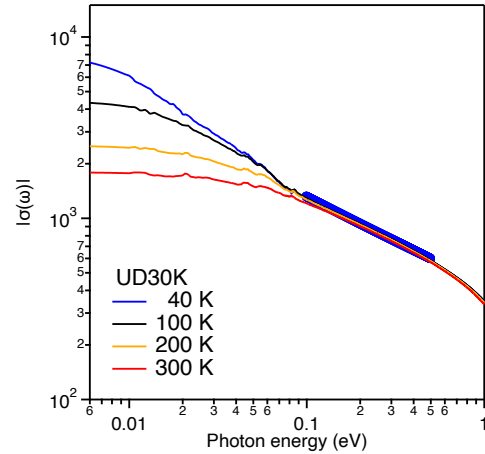
# Component 1: a Drude response

$$\frac{\partial \log[\sigma_1(\omega)]}{\partial \log[\omega]}$$





# Component 2: a high energy powerlaw



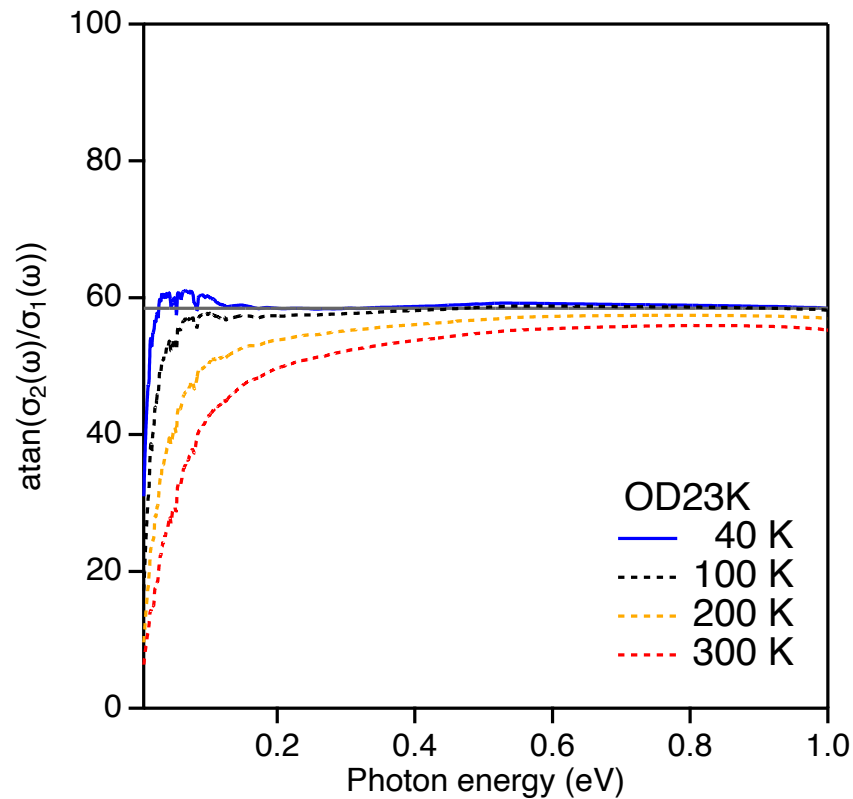
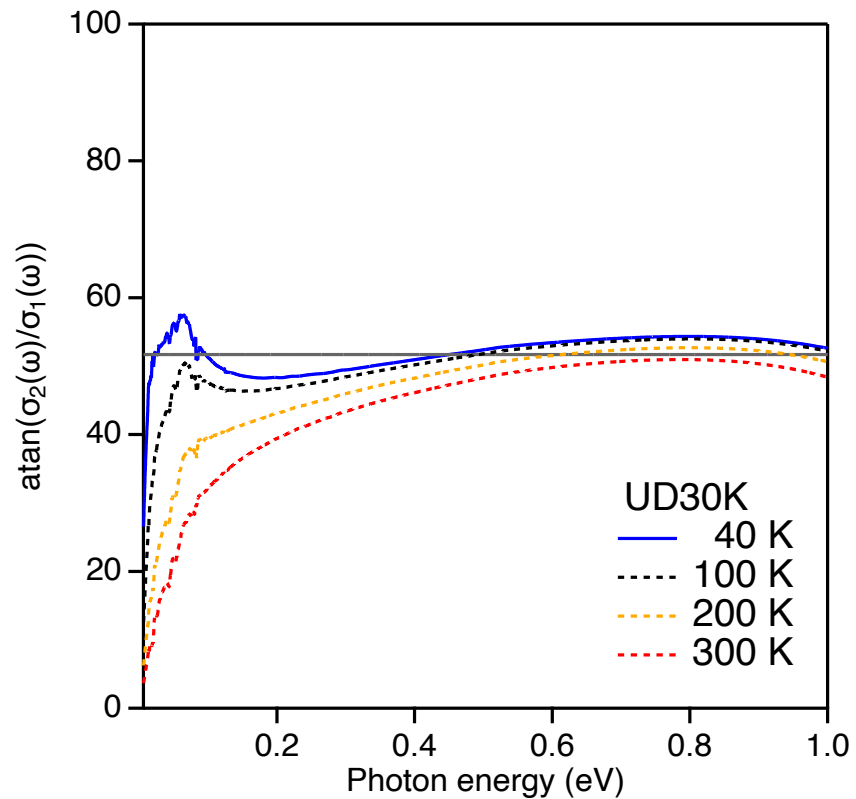
- At high energy  $|\sigma|$  follows a powerlaw behaviour.

$$\sigma(\omega) \sim 1/(i\omega)^\alpha$$

- Should be accompanied by a constant phase angle.

Marel, D. van der *et al. Nature* **425**, 271 (2003)

## Component 2: doping dependence of the phase angle



# A two-component model

- To get ahead we assume:

$$\hat{\sigma}(\omega) = \hat{\sigma}^D(\omega) + \hat{\sigma}^{inc.}(\omega)$$

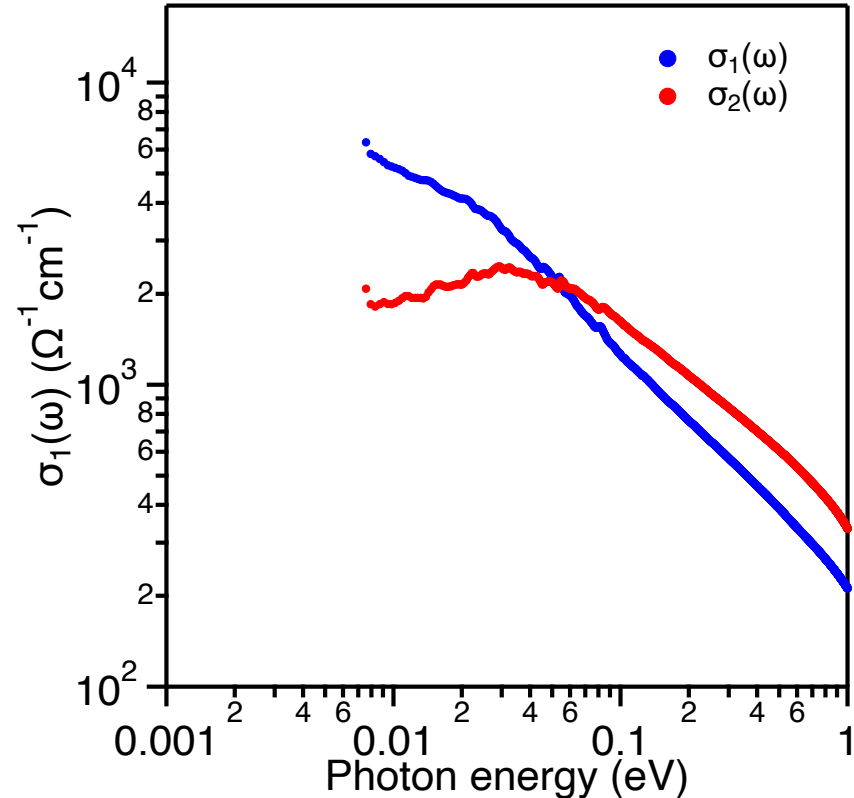
The Drude is easy:

$$\hat{\sigma}^D(\omega) = \frac{D_{Dr}}{4\pi} \frac{1}{\Gamma_{Dr} + i\omega}$$

However,

$$\sigma(\omega) \sim 1/(i\omega)^\alpha$$

both  $\sigma_1$  and  $\sigma_2$  diverge



# A two-component model

- To get ahead we assume:

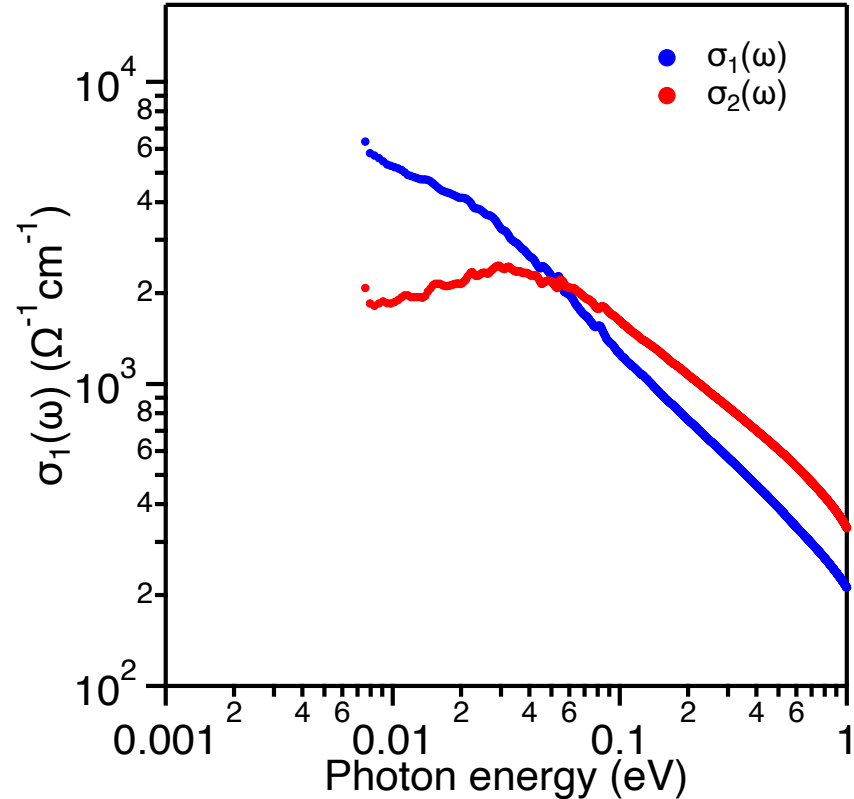
$$\hat{\sigma}(\omega) = \hat{\sigma}^D(\omega) + \hat{\sigma}^{inc.}(\omega)$$

The Drude is easy:

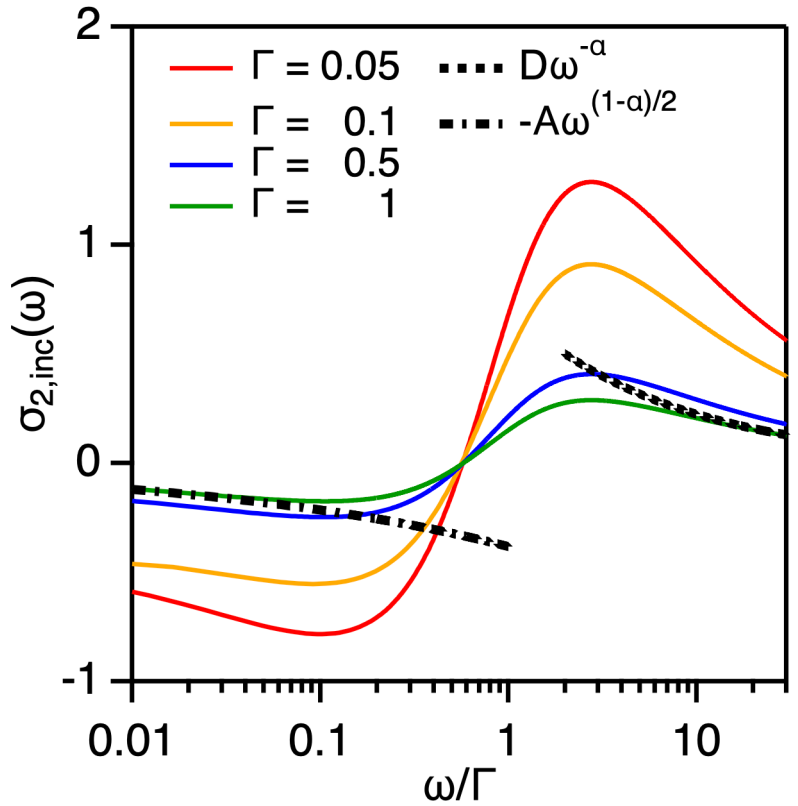
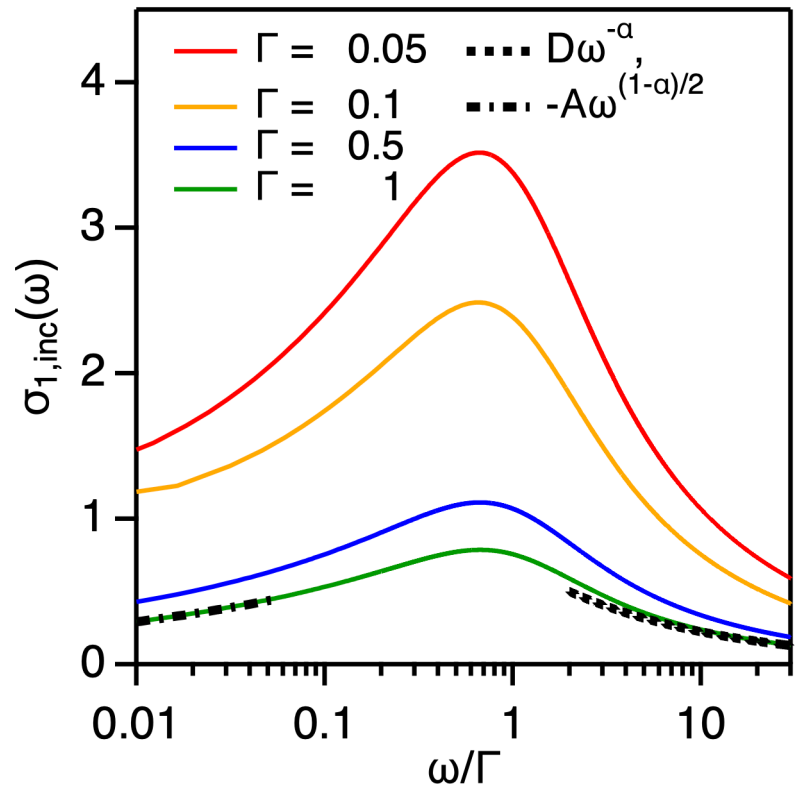
$$\hat{\sigma}^D(\omega) = \frac{D_{Dr}}{4\pi} \frac{1}{\Gamma_{Dr} + i\omega}$$

The incoherent part:

$$\hat{\sigma}^{inc.}(\omega) = \frac{-iD_{inc}\omega}{\sqrt{\Delta^2 - \omega^2 - i\Gamma_{inc}\omega}^\beta}$$



# Generalized interband transitions



# A two-component model

- To get ahead we assume:

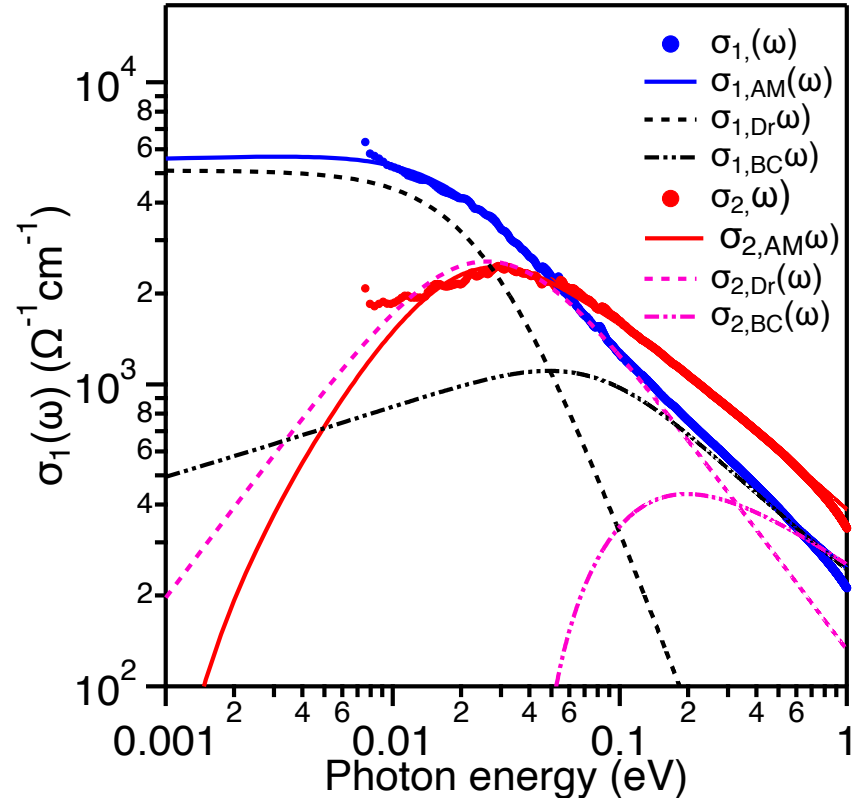
$$\hat{\sigma}(\omega) = \hat{\sigma}^D(\omega) + \hat{\sigma}^{inc.}(\omega)$$

The Drude is easy:

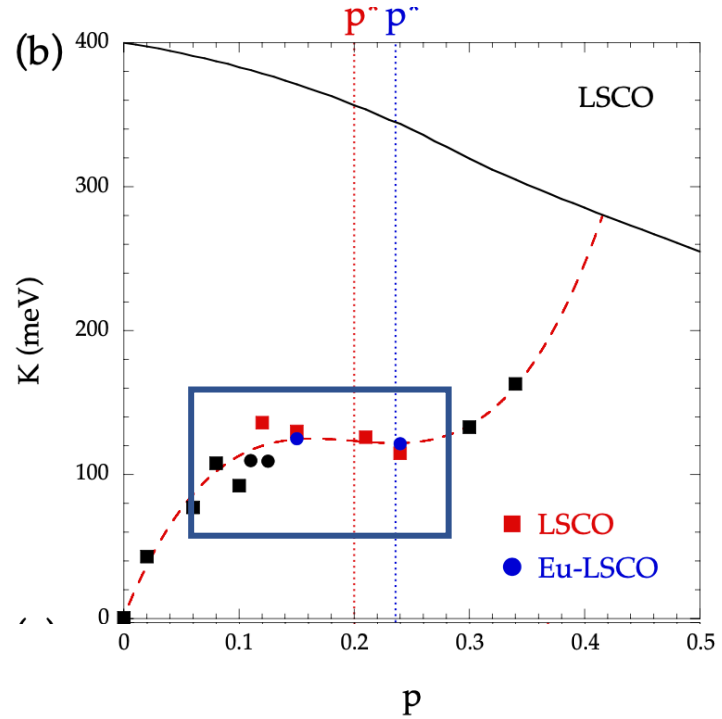
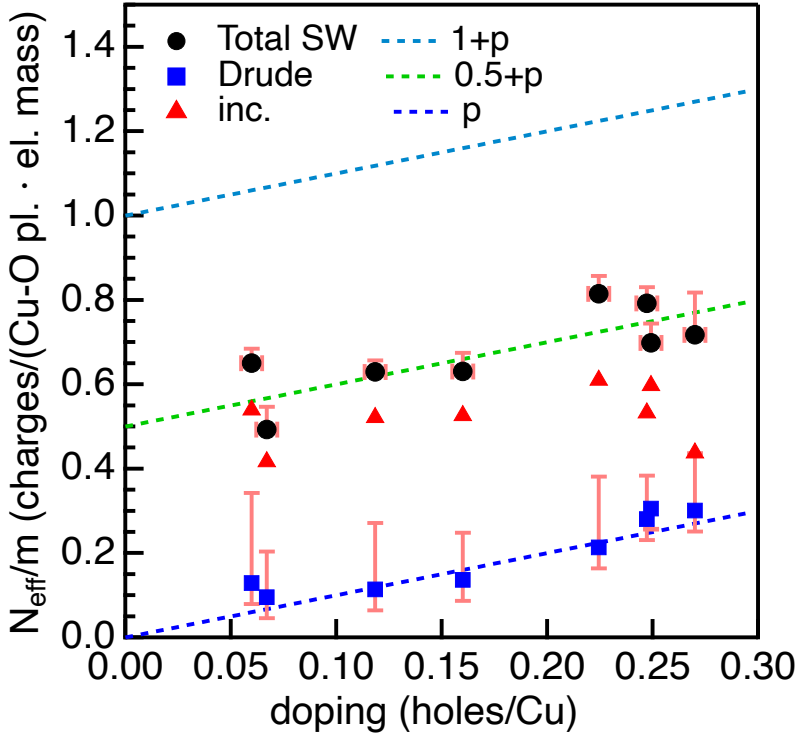
$$\hat{\sigma}^D(\omega) = \frac{D_{Dr}}{4\pi} \frac{1}{\Gamma_{Dr} + i\omega}$$

The incoherent part:

$$\hat{\sigma}^{inc.}(\omega) = \frac{-iD_{inc}\omega}{\sqrt{\Delta^2 - \omega^2 - i\Gamma_{inc}\omega}^\beta}$$



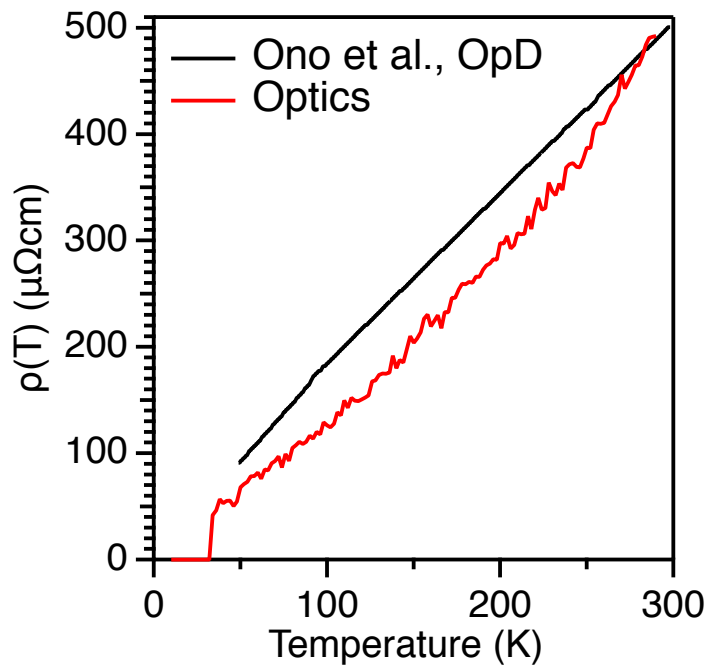
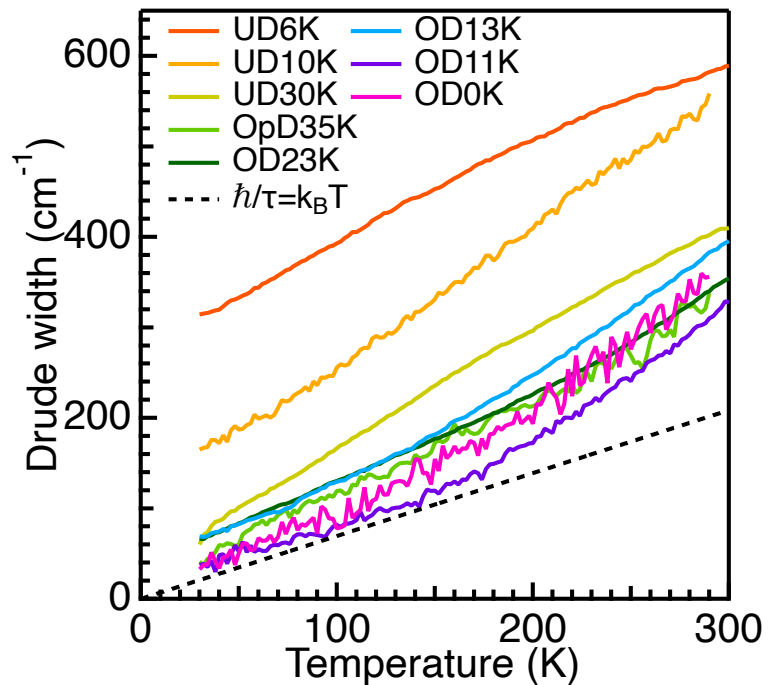
# Doping dependence of carrier density



# The current relaxation rate

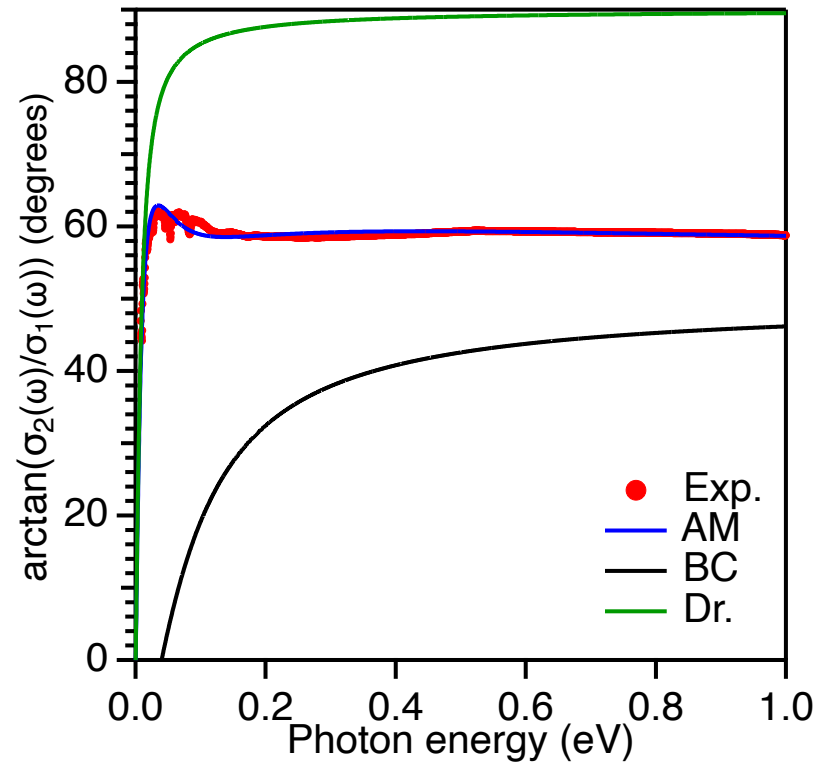
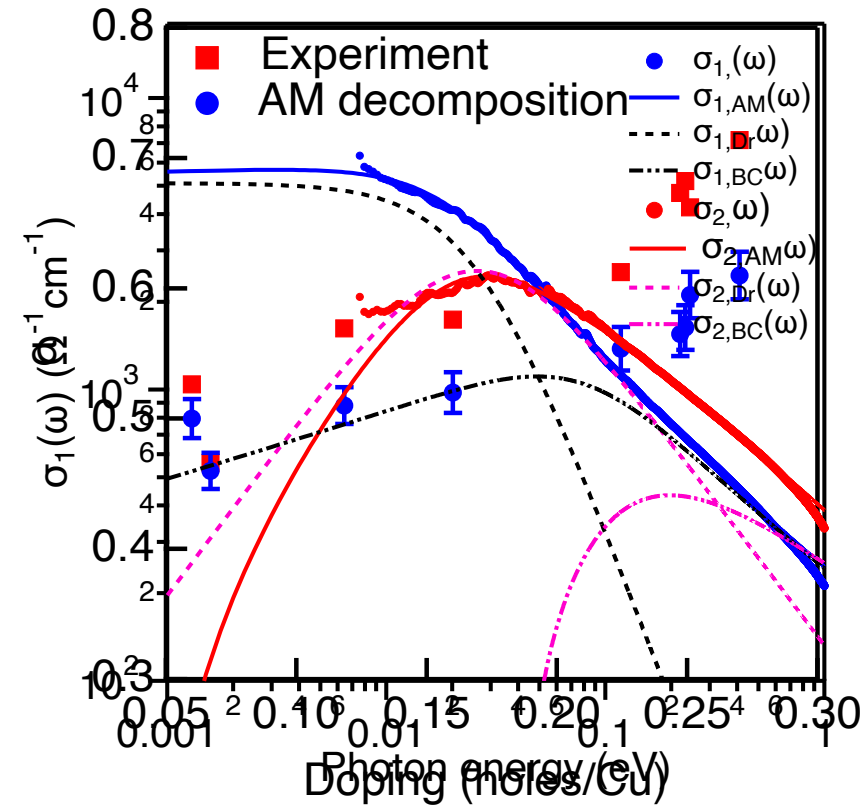
- Within our error bars,  $\Gamma$  is frequency independent.

$$\hat{\sigma}^D(\omega) = \frac{D_{Dr}}{4\pi} \frac{1}{\Gamma_{Dr} + i\omega}$$

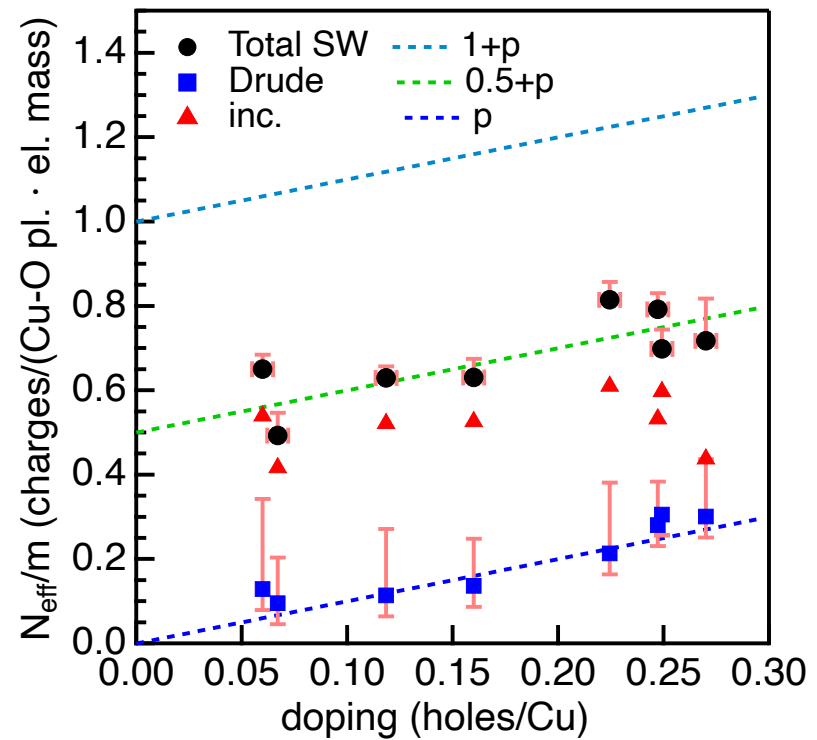
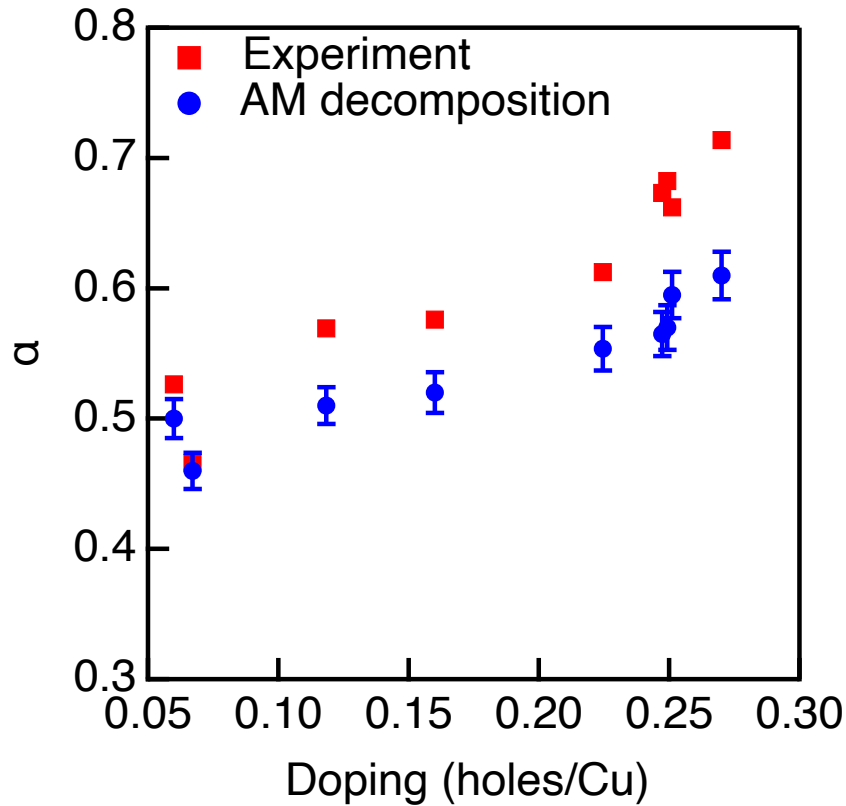




# The MIR component: a 'strange' continuum



# The MIR component: a 'strange' continuum



# Summary

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- Optical response has a Drude and mid-infrared response.
  - Low energy Drude response:
    - Spectral weight is doping but not temperature dependent.
    - Scattering rate is doping and temperature dependent.
    - No divergencies/anomalies in either.
  - The mid-infrared response
    - Spectral weight is not doping/temperature dependent.
    - Broadening is doping dependent.
- incoherent response persists to strongly overdoped regime!
- Optics and normal state transport agree. The puzzle is low T/high B!