Exploring the Phase Diagram of V-QCD with Neutron Star Mergers

Christian Ecker







HoloTube Seminar 27 September 2022

Based on 2112.12157 (PRX) and 2205.05691 (SciPost) with Tuna Demircik, Matti Järvinen, Luciano Rezzolla, Samuel Tootle and Konrad Topolski

DFG



























Density Profile



Demircik, CE, Järvinen 2009.10731 (ApJL)













Equation of State

Relates properties of dense matter to properties of neutron stars.

$$p = p(e, \ldots) \xrightarrow{\mathsf{Einstein Eq.}} M(R)$$

Constrained by theory and neutron star observations:

- Causality & thermodynamic stability: $0 < c_s^2 = rac{\mathrm{d} p}{\mathrm{d} e} < 1$
- Direct mass and radius measurements: $M_{\rm TOV}\gtrsim 2\,M_{\odot}$, 11 km $\lesssim R\lesssim$ 14 km
- \blacktriangleright First GW detection from a BNS merger GW170817: $\tilde{\Lambda} \lesssim 720$



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Model Agnostic Approach

 Impose boundary conditions from nuclear theory and pertubative QCD. Hebeler, Lattimer, Pethick, Schwenk 1303.4662 (ApJ);

Fraga, Kurkela, Vuorinen 1311.5154 (ApJL)

- Connect with random, but causal and thermodynamically consistent EOS.
- Constrain with pulsar and binary neutron star observations.



Lots of recent work:

 ...; Annala, Gorda, Kurkela, Nättilä, Vuorinen 1903.09121 (Nature Physics), 2105.05132 (PRX); Komoltsev, Kurkela 2111.05350 (PRL); Altiparmak, CE, Rezzolla 2203.14974; Gorda, Komoltsev, Kurkela 2204.11877; CE, Rezzolla 2207.04417 (ApJL); ...

Sampling the EOS

- Statistics of millions of EOS models that pass theory+astro constraints.
- Mass bound of generic (sub-conformal) models: $M_{\rm TOV} < 3(2.1) M_{\odot}$.
- Radius estimates (+ error bars) for typical stars:

 $R_{1.4} = 12.42^{+0.52}_{-0.99} \,\mathrm{km} \,, \, R_{2.0} = 12.11^{+1.11}_{-1.23} \,\mathrm{km} \left(95\% \,\mathrm{confidence}\right) \,. \tag{1}$

• Remarkable agreement with $R_{min}(M)$ from threshold mass calculations.



Altiparmak, CE, Rezzolla 2203.14974

Koeppel, Rezzolla 1901.09977 (ApJL)

Simple formulas for NS and BNS properties



$$c_{s}^{2}\left(\frac{r}{R}\right) = \left(\alpha e^{\beta\left(\frac{r}{R}\right)^{2}} + \gamma e^{\delta\left(\frac{r}{R} - \epsilon\right)^{2}}\right) \left[1 - \tanh\left(\zeta\frac{r}{R} - \eta\right)\right] \qquad \tilde{\Lambda}_{min(max)} = a + b \mathcal{M}_{chip}^{c}$$

CE, Rezzolla 2207.04417 (ApJL); Altiparmak, CE, Rezzolla 2203.14974

PSR J0952-0607

Recent estimate for the black widow binary pulsar PSR J0952-0607:

 $M = 2.35 \pm 0.17 \, M_{\odot}$.

Romani, Kandel, Filippenko, Brink, Zheng 2207.05124 (ApJL)



New estimates for neutron star radii and binary-tidal deformability:

 $R_{1.4} \in [12.48, 12.97] \, \mathrm{km} \,, \tilde{\Lambda}_{1.186}^{\mathrm{min}} \in [236, 301] \text{ for } M_{\mathrm{TOV}} \in [2.18, 2.52] \, M_{\odot}.$ CE. Rezzolla 2209.08101

A Holographic Approach to Dense QCD Matter



Applications of Holography to Neutron Stars

• Early attempts with D3-D7 hybrid EOS models.

Hoyos, Fernandez, Jokela, Vuorinen 1603.02943 (PRL); Annala, CE, Hoyos, Fernandez, Jokela, Vuorinen 1711.06244 (JHEP)

- First BNS merger simulations with input from holography.
 CE. Järvinen, Niis, van der Schee 1908.03213 (PRD)
- Rapidly rotating stars (GW190814) with V-QCD model.

Demircik, CE, Järvinen 2009.10731 (ApJL)

- Superconducting holographic quark matter in compact stars. Bitaghsir Fadafan, Cruz Rojas, Evans 2009.14079 (PRD)
- Neutron stars with superconducting QCD matter from 6D holography. Ghoroku, Kashiwa, Nakano, Tachibana, Toyoda 2107.14450 (PRD)
- Realistic neutron stars from Witten-Sakai-Sugimoto (WSS) model. Kovensky, Poole, Schmitt 2111.03374 (PRD)
- Merger simulations with holographic hard-wall model.

Bartolini, Gudnason, Leutgeb, Rebhan 2202.12845 (PRD)

 More references on holographic EOS modelling in two review articles. Järvinen 2205.05691 (EPJC), Hoyos, Jokela, Vuorinen 2112.08422 (Prog.Part.Nucl.Phys.)

Holographic Veneziano QCD

Holographic bottom-up model for QCD with many parameters that are tuned to mimic QCD and that are constrained by lattice QCD data.

Järvinen, Kiritsis 1112.1261 (JHEP)

Gluons: Improved holographic QCD (Einstein-dilaton gravity)

$$S_{\rm g} = N_c^2 M^3 \int d^5 x \sqrt{-g} \left[R - rac{4}{3} rac{(\partial \lambda)^2}{\lambda^2} + V_g(\lambda)
ight]$$

 $\lambda \equiv e^{\phi} \leftrightarrow \text{Tr}F^2$ sources the 't Hooft coupling in YM theory Gürsoy, Kiritsis 0707.1324 (JHEP); Gürsoy, Kiritsis, Nitti 0707.1349 (JHEP)

Quarks: Tachyonic Dirac-Born-Infeld (DBI) action

$$\begin{split} S_{\rm f} &= -N_f N_c M^3 \int d^5 x \frac{V_{f0}(\lambda) e^{-\tau^2} \sqrt{-\det\left[g_{ab} + \kappa(\lambda) \partial_a \tau \partial_b \tau + w(\lambda) F_{ab}\right]}}{F_{rt} = \Phi'(r) , \quad \Phi(0) = \mu \,, \end{split}$$

tachyon $\tau \leftrightarrow \bar{q}q$ controls chiral symmetry breaking.

Bigazzi et al. 0505140 (JHEP); Casero et al. 0702155 (Nucl.Phys.B)

Baryons: homogeneous solution of non-Abelian DBI + Cern–Simons action. Ishii, Järvinen, Nijs 1903.06169 (JHEP)

Veneziano limit: $N_c \to \infty$ and $N_f \to \infty$ with $N_f/N_c = \mathcal{O}(1)$ fixed Järvinen, Kiritsis 1112.1261 (JHEP)

Cold V-QCD Hybrid Equation of State

- Computing neutron star properties requires EOS at low and high densities.
- Homogeneous approximation for V-QCD baryons not reliable at low densities: use results from nuclear theory to model the low density part.
- Hybrid EOSs: nuclear theory model at low densities + V-QCD model for dense baryonic and quark matter at large densities.



CE, Järvinen, Nijs, van der Schee 1908.03213 (PRD); Jokela, Järvinen, Nijs, Remes 2006.01141 (JHEP); Demircik, CE, Järvinen 2112.12157 (PRX)

Hot V-QCD Hybrid Equation of State

- Low density part: Hemple-Schaffner-Bielich model (HS)DD2 + Particle Data Group (PDG) mesons. Hempel, Schaffner-Bielich 0911.4073 (Nucl.Phys.A); Zyla et al. (Particle Data Group 2020)
- (HS)DD2 too stiff around n_s, replace with APR model.

Akmal, Pandharipande, Ravenhall 9804027 (PRC)

- Van der Waals construction to extend cold V-QCD baryons to finite-T. Vovchenko, Motornenko, Alba, Gorenstein, Satarov, Stoecker 1707.09215 (PRC)
- Mixed baryon and quark matter phase from Gibbs construction



V-QCD Critical Point



Estimates for the critical point:

 $\begin{array}{l} 110 \ \mathrm{MeV} \lesssim \ T_c \lesssim 130 \ \mathrm{MeV} \\ 0.3 \ n_s \lesssim \ n_c \lesssim 0.6 \ n_s \end{array}$

model	$\frac{n_{\rm bc}}{n_{\rm s}}$	$\frac{\mu_{bc}}{\text{MeV}}$	$\frac{T_{c}}{MeV}$
soft	0.46	485	128
interm.	0.62	575	118
stiff	0.32	565	112



Demircik, CE, Järvinen 2112.12157 (PRX)

Binary neutron star mergers

- Significant sources of gravitational radiation.
- Microscopic properties of dense matter encoded in GW and EM signal.
- Likely the origin of short gamma-ray bursts
- ▶ and of heavy elements like gold, platinum, uranium, ...



Picture: Baiotti, Rezzola 1607.03540 (Rept.Prog.Phys.)

GW170817

• Best event seen so far: GW, optical, X-ray, γ , UV, IR

 $d \approx 40 \, {
m Mpc} \,, \quad {\cal M}_{
m chirp} = rac{(M_1 M_2)^{3/5}}{(M1+M2)^{1/5}} = 1.188^{+0.004}_{-0.002} \, M_{\odot} \,, \quad q = M_2/M_1 > 0.73 \,.$

• Inspiral (but no post-merger) waveform: $\tilde{\Lambda} \lesssim 720$

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LIGO/Virgo 1805.11579 (PRX)
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• EM-counterpart: estimates for post-merger HMNS lifetime $\approx 1 \text{ ms.}$



Gill, Nathanail, Rezzolla 1901.04138 (ApJ)

LIGO/Virgo 1710.05832 (PRL), 1710.05834 (ApJL) 20/30

Simulating Binary Neutron Star Mergers

Have to solve the 3+1D General Relativistic hydrodynamic equations:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi G_N T_{\mu\nu}, \quad \nabla_\mu T^{\mu\nu} = 0, \quad \nabla_\mu J^\mu = 0$$

with initial spacetime and fluid distribution modelling a NS binary system.

- Equation of State $p = p(n_b, T, Y_e)$ as input required.
- Spectral code Frankfurt University/Kadath (FUKA) for initial data.
 Papenfort, Tootle, Grandclement, Most, Rezzolla 2103.09911 (PRD)
- Frankfurt/Illinois (FIL) code for binary evolution with tabulated EOS.

Most, Papenfort, Rezzolla 1907.10328 (MNRAS)

Implemented in the Einstein Toolkit.

http://einsteintoolkit.org

High-Performance Computing Center Stuttgart

▶ Project BNSMIC: 100+ million core-hours on HAWK.



Picture Copyright: Ben Derzian for HLRS



Hot, Warm and Cold Quarks



Tootle, CE, Topolski, Demircik, Järvinen, Rezzolla 2205.05691 (SciPost)

Phase Diagram

 $t - t_{\rm mer} \, [{\rm ms}]$



Tootle, CE, Topolski, Demircik, Järvinen, Rezzolla 2205.05691 (SciPost)

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Quark Abundance





Tootle, CE, Topolski, Demircik, Järvinen, Rezzolla 2205.05691 (SciPost)

Extracting Waveforms

• Newman-Penrose scalar ψ_4 from simulation

 $\psi_4 := C_{\mu\nu\rho\sigma} k^\mu \bar{m}^\nu k^\rho \bar{m}^\sigma \,, \quad \text{null tetrade: } \{ I^\mu, k^\mu, m^\mu, \bar{m}^\mu \} \,.$

• GW polarization amplitudes (h_+, h_{\times}) are related to ψ_4 via

see e.g. book by M. Alcubierre (2005)

$$\ddot{h}_{+} - i\ddot{h}_{\times} = \psi_4 = \sum_{\ell=2}^{\infty} \sum_{m=-\ell}^{\ell} \psi_4^{\ell m} {}_{-2} \boldsymbol{Y}_{\ell m}(\theta, \varphi) \,.$$

• Consider for example the dominant $\ell = m = 2 \mod \ell$

$$h_{+,\times} = \sum_{\ell=2}^{\infty} \sum_{m=-\ell}^{\ell} h_{+,\times-2}^{\ell m} \boldsymbol{Y}_{\ell m}(\boldsymbol{\theta},\boldsymbol{\varphi}) \approx h_{+,\times-2}^{22} \boldsymbol{Y}_{22}(\boldsymbol{\theta},\boldsymbol{\varphi}) \,,$$

- Optimal orientation $({}_{-2}Y_{22}(0,0) = \frac{1}{2}\sqrt{5/\pi})$ with respect to the detector.
- Extrapolate signal to same luminosity distance (40 Mpc) as GW170817.

Power Spectral Density

Post-merger power spectral density (PSD) has typical three peak structure

$$\tilde{h}(f) \equiv \sqrt{\frac{|\tilde{h}_+(f)|^2 + |\tilde{h}_\times(f)|^2}{2}}, \quad \tilde{h}_{+,\times}(f) \equiv \int h_{+,\times}(t) e^{-i2\pi f t} dt.$$

Characteristic frequencies f_1 , f_2 , f_3 encode information about EOS.



Waveforms and Frequency Spectra

Small impact of quark matter on post-merger frequencies.



Tootle, CE, Topolski, Demircik, Järvinen, Rezzolla 2205.05691 (SciPost) 29/30

Waveforms and Frequency Spectra

- Small impact of quark matter on post-merger frequencies.
- Phase transition triggered collapse (PTTC) leads to shorter lifetime of the hyper massive neutron star (HMNS).



Tootle, CE, Topolski, Demircik, Järvinen, Rezzolla 2205.05691 (SciPost) 29/30

Summary

 Hot V-QCD hybrid EOS with deconfinement phase transition, three tabulated versions soon on CompOSE database (stay tuned).

https://compose.obspm.fr/

Predictions for the QCD critical point:

 $110~{\rm MeV} \lesssim T_c \lesssim 130~{\rm MeV}\,, \quad 0.3~n_s \lesssim n_c \lesssim 0.6~n_s\,.$

Demircik, CE, Järvinen 2112.12157 (PRX)

- Neutron star merger simulations with V-QCD EOS.
- Identified 3 different stages in HMNSs: hot, warm, cold quark matter.
- Small impact on frequencies, but PTTC shortens HMNS lifetime.

Tootle, CE, Topolski, Demircik, Järvinen, Rezzolla 2205.05691 (SciPost)

Ongoing work:

▶ ...

- Prompt collapse: Imprint of quark formation on threshold mass?
- Long-time stability (lifetime) of HMNS: Need simulations up to 1-2 sec. (expensive, simplified spacetime evolution, axial-symmetry, ...)

V-QCD without baryons (I)

Consider first the non-baryonic V-QCD action, whose solutions will serve as background for the probe baryons

$$S_{
m V-QCD}^{(0)} = S_{
m glue} + S_{
m DBI}^{(0)}$$
 .

The gluon part is given by the IHQCD (dilaton gravity) action

$$S_{
m glue} = N_c^2 M^3 \int d^5 x \sqrt{-g} \left[R - rac{4}{3} rac{(\partial \lambda)^2}{\lambda^2} + V_g(\lambda)
ight],$$

where $\lambda \equiv e^{\phi} \leftrightarrow \text{Tr}F^2$ ($\approx g^2 N_c$ near the boundary) sources the 't Hooft coupling in YM theory, the dilaton potential is chosen¹ to mimic QCD

$$V_g(\lambda) = 12 \left[1 + V_1 \lambda + \frac{V_2 \lambda^2}{1 + \lambda/\lambda_0} + V_{\rm IR} e^{-\lambda_0/\lambda} (\lambda/\lambda_0)^{4/3} \sqrt{\log(1 + \lambda/\lambda_0)} \right] \,.$$

Finite T is implemented by homogeneous+isotropic black brane metric

$$ds^{2} = e^{2A(r)}(-f(r)dt^{2} + d\vec{x}^{2} + f^{-1}(r)dr^{2}).$$

¹E.g. V_1 and V_2 are fixed by requiring the UV RG flow of the 't Hooft coupling to be the same as in QCD up to two-loop order.

V-QCD without baryons (II)

The flavor part is modelled by the tachyonic DBI-action²

$$\begin{split} S_{\rm DBI}^{(0)} &= -N_f N_c M^3 \int d^5 x \frac{V_{f0}(\lambda) e^{-\tau^2} \sqrt{-\det \left[g_{ab} + \kappa(\lambda) \partial_a \tau \partial_b \tau + w(\lambda) F_{ab}\right]}}{F_{rt} = \Phi'(r) \,, \quad \Phi(0) = \mu \,, \end{split}$$

where the tachyon $\tau \leftrightarrow \bar{q}q$ controls chiral symmetry breaking.

Several potentials: { $V_g(\lambda), V_{f0}(\lambda), w(\lambda), \kappa(\lambda)$ }, chosen to match pQCD in UV ($\lambda \rightarrow 0$), qualitative agreement with QCD in IR ($\lambda \rightarrow \infty$) and tuned to lattice QCD in the middle ($\lambda \sim O(1)$).

For details see Appendix B of Ishii, Järvinen, Nijs arXiv:1903.06169

Different solutions:

without/with horizon ↔ confined/deconfined phase without/with tachyon ↔ chirally symmetric/chirally broken phase

²Without baryons we have a vectorial flavor singlet gauge field $A^{(L/R)} = \mathbb{I}_f \Phi(r) dt$ giving nonzero charge density and chemical potential.

Probe baryons in V-QCD

Each baryon maps to a solitonic "instanton" configuration of non-Abelian gauge fields in the bulk.

Witten; Gross, Ooguri; ...

Consider the full non-Abelian brane action $S = S_{\text{DBI}} + S_{\text{CS}}$ where

Bigazzi, Casero, Cotrone, Kiritsis, Paredes; Casero, Kiritsis, Paredes

$$\begin{split} S_{\text{DBI}} &= -\frac{1}{2} M^3 N_c \, \mathbb{T}r \int d^5 x \, V_{f0}(\lambda) e^{-\tau^2} \left(\sqrt{-\det \mathbf{A}^{(L)}} + \sqrt{-\det \mathbf{A}^{(R)}} \right) \,, \\ \mathbf{A}_{MN}^{(L/R)} &= g_{MN} + \delta_M^r \delta_N^r \kappa(\lambda) \tau'(r)^2 + \delta_{MN}^{rt} w(\lambda) \Phi'(r) + w(\lambda) F_{MN}^{(L/R)} \end{split}$$

gives the dynamics of the solitons.

The Cern-Simons term sources the baryon number for the solitions

$$S_{\rm CS} = \frac{N_c}{8\pi^2} \int \Phi(r) e^{-b\tau^2} dt \wedge \left(F^{(L)} \wedge F^{(L)} - F^{(R)} \wedge F^{(R)} + \cdots\right) \, .$$

Non-Abelian DBI action only known to first few orders in $F^{(L/R)}$: expand to second order on top of solution $(g_{MN}, \Phi, \lambda, \tau)$ obtained from $S_{V-QCD}^{(0)}$.

Homogeneous Baryon Ansatz

Set $N_f = 2$ and consider the SU(2) Ansatz

Rozali, Shieh, Van Raamsdonk, Wu

$$A_L^i = -A_R^i = h(r)\sigma^i$$

Immediate consequence: baryon charge integrates to zero?

$$N_b \propto \int dr \frac{d}{dr} \left[e^{-b\tau^2} h^3 (1 - 2b\tau^2) \right] \stackrel{?}{=} 0$$

However finite baryon number may can be realized by discontinuity of $h \leftrightarrow$ smeared solitons in singular gauge.

Ishii, Järvinen, Nijs, arXiv:1903.06169

The free parameter b of the model is used to tune the baryon onset to its physical value in QCD.

Binary tidal deformability and chirp mass

Chirp mass can be extracted accurately from inspiral GW strain signal

$$\mathcal{M}_{chirp} = \frac{(M_1 M_2)^{3/5}}{(M_1 + M_2)^{1/5}} = \left(\frac{5}{96}\pi^{-8/3}f^{-11/3}\dot{f}\right)^{3/5}.$$
 (2)

Binary tidal deformability encodes EOS properties via Λ_i

$$\tilde{\Lambda} = \frac{16}{13} \frac{(12M_2 + M_1) M_1^4 \Lambda_1 + (12M_1 + M_2) M_2^4 \Lambda_2}{(M_1 + M_2)^5} \,. \tag{3}$$

GW170817 only "good" event so far:

$$\mathcal{M}_{chirp} = 1.188^{+0.004}_{-0.002} M_{\odot}, \quad q = M_2/M_1 > 0.7, \quad \tilde{\Lambda} \lesssim 720.$$
 (4)

LIGO/Virgo: arXiv:1805.11579

Mechanical Toy Model



Takami, Rezzolla, Baiotti arXiv:1412.3240





Phase Transition Signatures in Wavefrom Spectra

- NPT: No Phase Transition
- PTTC: Phase Transition Triggered Collapse
- DPT: Delayed Phase Transition



Weih, Hanauske, Rezzolla arXiv:1912.09340

Finite Temperature Extension

- Temperature dependence is essential when studying BNS-mergers.
- T-dependence of V-QCD baryons is trivial: confining background has no horizon, artefact of large N_c limit.
- ▶ Van der Waals construction to extend cold V-QCD baryons to finite-*T*.
- Ideal gas of protons, neutrons and electrons with excluded volume correction for nucleons.

 $p_{\mathrm{ex}}(T, \{\mu_i\}) = p_{\mathrm{id}}(T, \{\tilde{\mu}_i\}), \quad \tilde{\mu}_i = \mu_i - v_0 p_{\mathrm{ex}}(T, \{\mu_i\}) \quad (i = p, n)$

• Add potential term to match with V-QCD at T = 0

$$p_{\rm vdW}(T, \{\mu_i\}) = p_{\rm ex}(T, \{\mu_i\}) + \Delta p(\{\mu_i\})$$
$$\Delta p(\{\mu_i\}) = p_{\rm V-QCD}(T = 0, \{\mu_i\}) - p_{\rm ex}(T = 0, \{\mu_i\})$$

Rischke, Gorenstein, Stoecker, Greiner, Z Phys. C 51, 485 (1991)

Vovchenko, Gorenstein, Stoecker, arXiv:1609.03975

Vovchenko, Motornenko, Alba, Gorenstein, Satarov, Stoecker, arXiv:1707.09215

Building blocks of the EOSs

• Lowest densities $(0 - 0.5 n_s)$: Baym-Pethick-Sutherland (BPS) EOS

Baym, Pethick, Sutherland (1971)

- ► Low densities $(0.5 1.2 n_s)$: $p(n) = K n^{\Gamma}$, $\Gamma \in [1.77, 3.23]$ Hebeler, Lattimer, Pethick, Schwenk arXiv:1303.4662
- High densities (1.2 40 n_s): speed of sound parametrization Annala, Gorda, Kurkela, Nättilä, Vuorinen arXiv:1903.09121
- Highest densities (\gtrsim 40 n_s): perturbative QCD results

$$p(\mu) = \frac{3}{4\pi^2} \left(\frac{\mu}{3}\right)^4 \left(c_1 - \frac{d_1 X^{-\nu_1}}{\mu/GeV - d_2 X^{-\nu_2}}\right), \quad X \in [1, 4], \quad (5)$$

 $c_1 = 0.9008$ $d_1 = 0.5034$ $d_2 = 1.452$ $\nu_1 = 0.3553$ $\nu_2 = 0.9101$.

Kurkela, Romatschke, Vuorinen arXiv:0912.1856; Fraga, Kurkela, Vuorinen arXiv:1311.5154



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