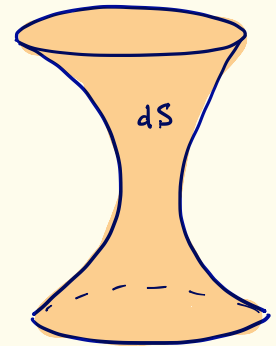
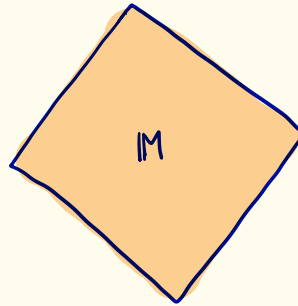
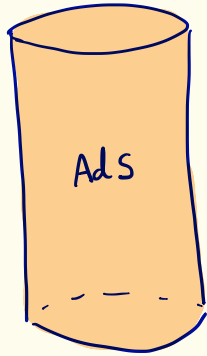


Towards the nonperturbative cosmological bootstrap



EPFL

João Penedones

HoloTube, 9/11/2021

Based on arXiv : 2107.13871 with :



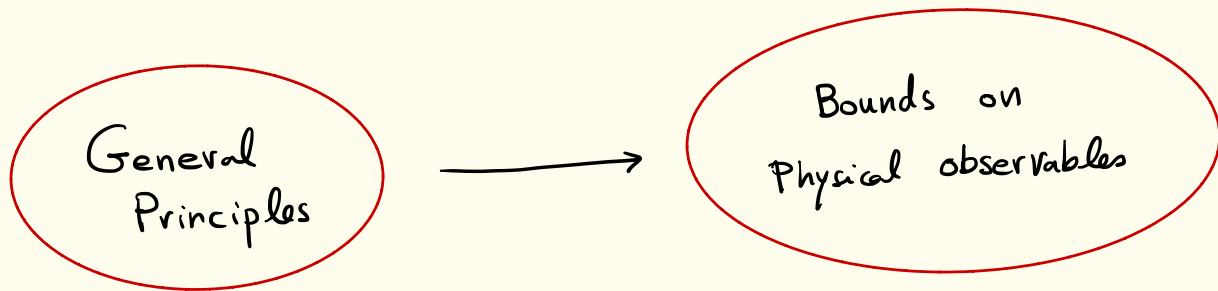
Matthijs Hogervorst



Kamran Vaziri

Introduction

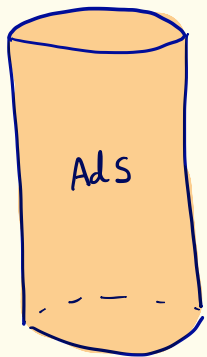
The bootstrap is a nonperturbative approach to QFT and Quantum Gravity.



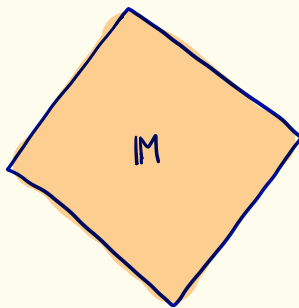
It works in different backgrounds: AdS, IM, dS.

Introduction

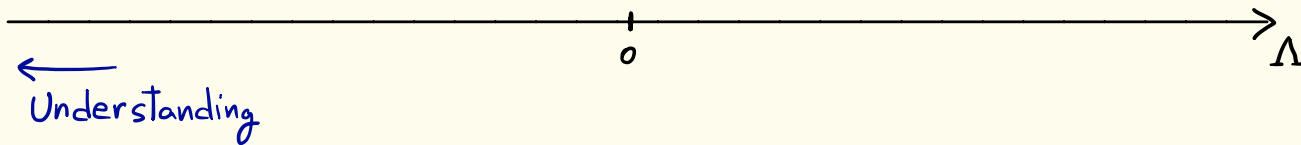
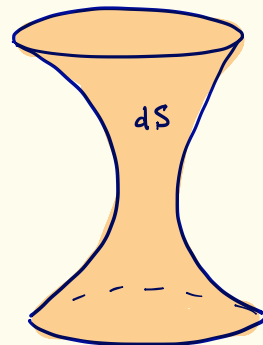
Conformal
Bootstrap



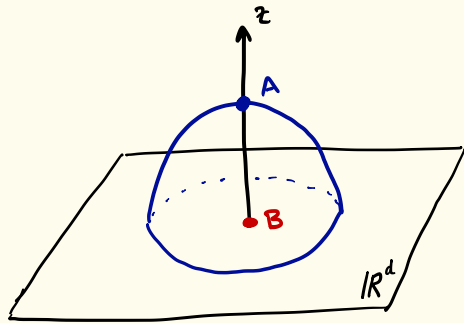
S-matrix
Bootstrap



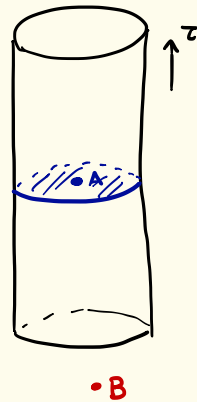
Cosmological
Bootstrap



QFT in AdS_{d+1}

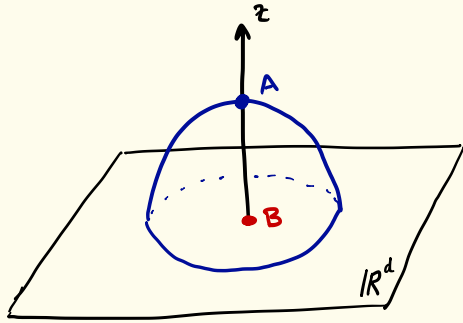


\Leftrightarrow

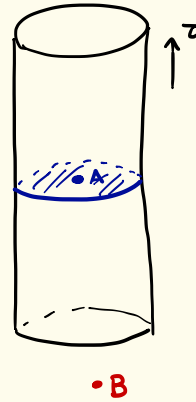


Boundary operator / bulk state map

QFT in AdS_{d+1}



\Leftrightarrow



Boundary operator / bulk state map

$$\Rightarrow \begin{cases} \phi(z, x) = \sum_k b_{\phi k} z^{\Delta_k} [\mathcal{O}_k(x) + \text{desc.}] \\ \mathcal{O}_i(x) \mathcal{O}_j(0) = \sum_k c_{ijk} |x|^{\Delta_k - \Delta_i - \Delta_j} [\mathcal{O}_k(0) + \text{desc.}] \end{cases}$$

bulk op. \downarrow

bound. op. \downarrow

OPE coefficients

Convergent

OPE

Boundary correlation functions obey the CFT_d axioms :

- Unitarity : $C_{ijk} \in \mathbb{R}$, $\Delta \geq \Delta_{\min} = \begin{cases} \frac{d-2}{2} & \text{scalars} \\ d-2+l & \text{spin } l \end{cases}$

- OPE convergence and associativity

$$\sum_k C_{12k} C_{34k} \begin{array}{c} 2 \\ \diagdown \\ \text{---} k \text{---} \\ \diagup \\ 1 \end{array} \begin{array}{c} 3 \\ \diagup \\ \text{---} \\ \diagdown \\ 4 \end{array} = \sum_q C_{14q} C_{23q} \begin{array}{c} 2 \\ \diagdown \\ \text{---} q \text{---} \\ \diagup \\ 1 \end{array} \begin{array}{c} 3 \\ \diagup \\ \text{---} \\ \diagdown \\ 4 \end{array} \leftarrow \text{conformal blocks}$$

Boundary correlation functions obey the CFT_d axioms:

• Unitarity: $C_{ijk} \in \mathbb{R}$, $\Delta \geq \Delta_{\min} = \begin{cases} \frac{d-2}{2} & \text{scalars} \\ d-2+l & \text{spin } l \end{cases}$

• OPE convergence and associativity

$$\sum_k \text{diagram}_k = \sum_q \text{diagram}_q$$

The diagram on the left shows two lines meeting at a vertex, with a horizontal line connecting them, and two lines meeting at a second vertex. The horizontal line is labeled with a 'k'. The diagram on the right shows two lines meeting at a vertex, with a vertical line connecting them, and two lines meeting at a second vertex. The vertical line is labeled with a 'q'.

• Stress tensor $T_{\mu\nu}$: $\Delta = d$, Ward identities

With dynamical gravity in AdS

QFT in Minkowski space

Focus on scattering amplitudes (boundary observables) :

- Unitarity (partial waves)
- Lorentz invariance
- Causality (some analyticity)
- Crossing symmetry

[- \exists massless spin 2 particle (graviton) : soft theorems
(need $d > 3$ for \mathbb{M}^{d+1})
(due to IR divergences)]

With dynamical gravity

What is the non-perturbative
bootstrap formulation for
QFT in de Sitter spacetime?

QFT on:	AdS_{d+1}	IM^{d+1}	dS_{d+1}
<u>Observable</u>	$\langle \mathcal{O}(x_1) \mathcal{O}(x_2) \mathcal{O}(x_3) \mathcal{O}(x_4) \rangle$ <p style="margin-left: 20px;">↙ boundary φ.</p> $= \frac{1}{(x_{13}^2 x_{24}^2)^{\Delta_0}}$ <p style="margin-left: 100px;">↑ ↑ Cross-ratios</p>	$\langle P_3, P_4 P_1, P_2 \rangle_{in}^{Conn} =$ $= i \delta(\sum P_i) T(s, t)$ <p style="margin-left: 100px;">↑ ↑ Mandelstam inv. $s+t+u=4m^2$</p>	?
<u>Crossing</u>	$F(u, v) = F(v, u)$ $= \frac{F(\frac{u}{v}, \frac{1}{v})}{u^{\Delta_0}}$	$T(s, t) = T(t, s)$ $= T(\frac{4m^2 - s - t}{u}, t)$?
<u>Decomposition in irreps</u>	$F(u, v) = \sum_{\Delta, \ell} \lambda_{\Delta, \ell} g_{\Delta, \ell}(u, v)$ <p style="margin-left: 100px;">↑ conf. blocks</p>	$T \propto \sum_{\ell} f_{\ell}(s) P_{\ell}(\underbrace{\cos \theta}_{1 - \frac{2t}{4m^2 - s}})$?
<u>Unitarity</u>	$\lambda_{\Delta, \ell} \geq 0$	$ 1 + i f_{\ell}(s) ^2 \leq 1$?

Related works

- Cosmological Bootstrap : perturbative
 - [Maldacena] [Arkani-Hamed , Maldacena]
 - [Arkani-Hamed , Baumann , Chun , Duaso Pueyo , Joyce , Lee , Pimentel]
 - [Goodnow , Jazayeri , Lee , Melville , Pajer , Stefanyshyn , Supel]
 - [Hollands] [Marolf , Morrison] ...
- Dictionary between dS and EAdS
 - [Sleight , Taronna]
 - [Di Pietro , Gorbenko , Komatsu]
 - ...
- dS / CFT
 - [Strominger] [Witten] [Anninos , Hartman , Strominger] [Susskind]
 - ...

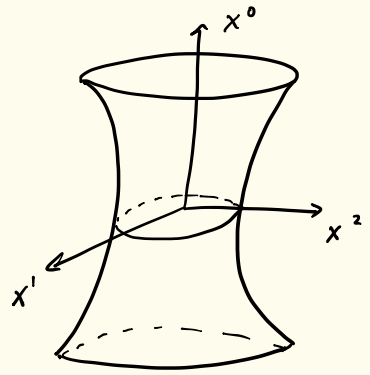
QFT in dS

de Sitter spacetime

$$dS_{d+1} \subset \mathbb{M}^{d+2} \ni x^A \quad A = 0, 1, \dots, d+1$$

$$-(x^0)^2 + (x^1)^2 + \dots + (x^{d+1})^2 = R^2$$

Isometry group = $SO(d+1, 1)$ = conformal group in \mathbb{R}^d

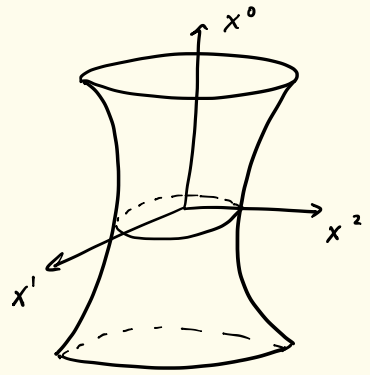


de Sitter spacetime

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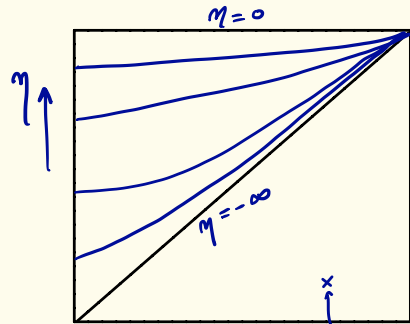
Isometry group = $SO(d+1, 1) =$ conformal group in \mathbb{R}^d



Conformal coordinates

$$ds^2 = R^2 \frac{d\eta^2 + dx^2}{\eta^2}$$

$$\eta < 0, \quad x \in \mathbb{R}^d$$

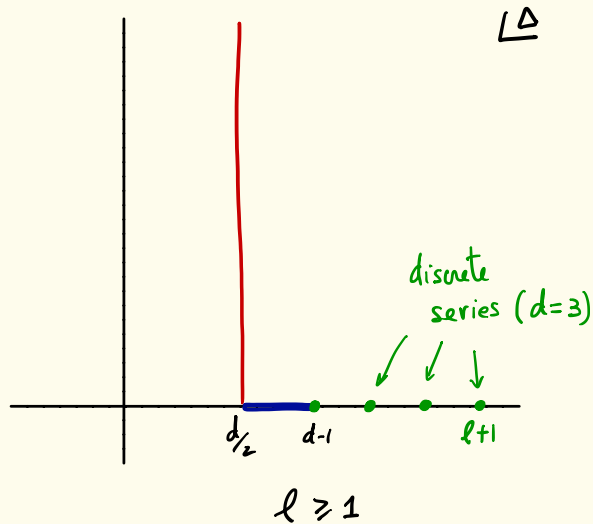
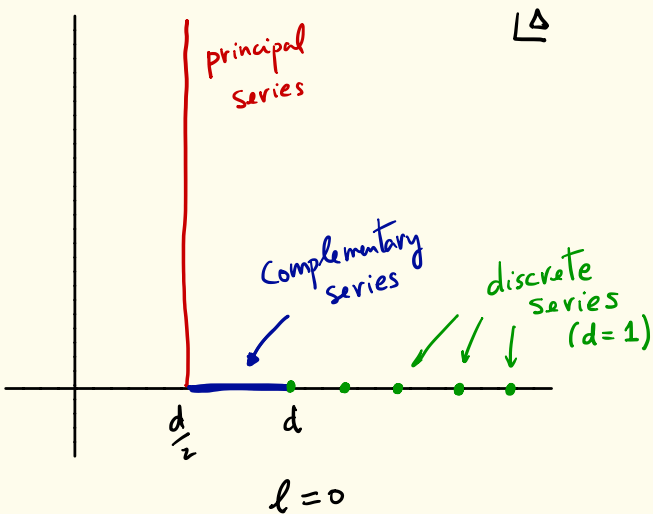


Penrose diagram S^{d-1}

Hilbert space

$$\mathcal{H} = \bigoplus [\text{Unitarity irreps of } SO(d+1, 1)]$$

$$SO(1, 1) \times SO(d)$$

 Δ
 $l \leftarrow \text{for simplicity } \boxed{\dots}$


Hilbert space

$$\mathcal{H} = \bigoplus \left[\text{Unitarity irreps of } \underset{\cup}{\text{SO}(d+1,1)} \right]$$

$\text{SO}(1,1) \times \text{SO}(d)$
 $\Delta \quad \ell \leftarrow \text{for simplicity } \boxed{\quad \dots \quad}$

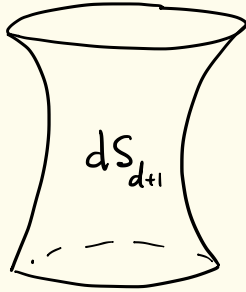
Basis for irrep (Δ, ℓ) : $|\Delta, x\rangle^{(\mu_1, \dots, \mu_d)}$, $x \in \mathbb{R}^d$, $\mu_i \in \{1, \dots, d\}$
 \uparrow transforms like primary local op $\mathcal{O}_{(\alpha)}^{\mu_1, \dots, \mu_d}(x)$

Projector
$$P_{\Delta, \ell} = \int d^d x |\Delta, x\rangle^{\mu_1, \dots, \mu_d} \langle \Delta, x |_{\mu_1, \dots, \mu_d}$$

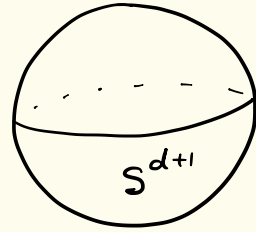
$$\mathbb{1} = |0\rangle\langle 0| + \sum_{\ell} \int_{\frac{d}{2}-i\infty}^{\frac{d}{2}+i\infty} d\Delta P_{\Delta, \ell} + \text{other irreps}$$

Correlation functions

Correlators in the Bunch-Davies vacuum can be obtained by analytic continuation from the sphere.



$$ds^2 = R^2(-dt^2 + \cosh^2 t d\Omega_d^2)$$



$$ds^2 = R^2(d\theta^2 + \cos^2 \theta d\Omega_d^2)$$

$$\langle 0 | \phi_1(t_1, \Omega_1) \dots \phi_n(t_n, \Omega_n) | 0 \rangle = \langle \phi(\theta_1 = \epsilon_1 + i t_1, \Omega_1) \dots \phi(\theta_n = \epsilon_n + i t_n, \Omega_n) \rangle_{\text{sphere}}$$

\uparrow
BD vacuum

$0 < \epsilon_n < \dots < \epsilon_1 \rightarrow 0$

Bulk Two-point function

$$\langle 0 | \phi(\eta, x) \phi(\eta', x') | 0 \rangle = G(\xi)$$

$$\xi = \frac{4\eta\eta'}{-(\eta-\eta')^2 + (x-x')^2} = \frac{4R^2}{\underbrace{(X-X')^2}_{\mathbb{M}^{d+2}}}$$

Bulk Two-point function

$$\langle 0 | \phi(\eta, x) \phi(\eta', x') | 0 \rangle = G(\xi)$$

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Källén-Lehmann decomposition

$$\langle 0 | \phi(\eta, x) \phi(\eta', x') | 0 \rangle = \langle 0 | \phi | 0 \rangle^2 + \int_{\frac{d}{2}-i\infty}^{d/2+i\infty} d\Delta \rho_\phi(\Delta) G_\Delta^{\text{free}}(\xi)$$

$\mathbb{1} = |0\rangle\langle 0| + \sum_i \int d\Delta \rho_{\Delta, i} + \dots$

positive spectral density
 $\rho_\phi(\Delta) = \rho_\phi(d-\Delta)$

Free propagator
 $m^2 R^2 = \Delta(\Delta-d)$

Bulk Two-point function

Continuation from the sphere:

$$\langle \phi(x) \phi(x') \rangle = G(x) = \sum_{J=0}^{\infty} a_J C_J^{d/2}(x)$$

$S^{d+1} \subset \mathbb{R}^{d+2}$
 $x^2 = 1 = x'^2$

↑ Gegenbauer polynomials

$$x = \cos \theta = X \cdot X'$$

$-1 \leq x \leq 1$

Bulk Two-point function

Continuation from the sphere:

$$\langle \phi(X) \phi(X') \rangle = G(x) = \sum_{J=0}^{\infty} a_J C_J^{d/2}(x)$$

$\overset{\text{in}}{\underbrace{S^{d+1} \subset \mathbb{R}^{d+2}}}$
 $X^2 = 1 = X'^2$

\uparrow Gegenbauer polynomials

$$x = \cos \theta = X \cdot X'$$

$-1 \leq x \leq 1$

$$a_J = \begin{cases} k_J \int_{-1}^1 dx (1-x^2)^{\frac{d-1}{2}} C_J^{d/2}(x) G(x) & \leftarrow \text{integer } J \\ \tilde{k}_J \int_1^{\infty} dx {}_2F_1\left(J+d, J+\frac{d}{2}+\frac{1}{2}, 2J+d+1, \frac{2}{1-x}\right) \frac{(x+1)^{\frac{d-1}{2}}}{(x-1)^{J+\frac{d}{2}+\frac{1}{2}}} \text{Disc}[G(x)] & \leftarrow \text{complex } J \end{cases}$$

$$\rho_{\phi}(\Delta) = \frac{2\pi^{d/2+1}}{\Gamma(\frac{d}{2})} \left[a_{J=-\Delta} + a_{J=\Delta-d} \right]$$

Bulk Two-point function - Examples

$$\langle 0 | \phi(\eta, x) \phi(\eta', x') | 0 \rangle = \langle \phi \rangle + \int_{\frac{d}{2}-i\infty}^{\frac{d}{2}+i\infty} d\Delta \rho_{\phi}(\Delta) G_{\Delta}^{\text{free}}(\xi)$$

Massive free boson

$$m^2 R^2 = \frac{d^2}{4} + \mu^2 \quad \Rightarrow \quad \rho_{\phi}(\Delta = \frac{d}{2} + i\nu) = \frac{1}{2} [\delta(\mu - \nu) + \delta(\mu + \nu)]$$

Bulk Two-point function - Examples

$$\langle 0 | \phi(\eta, x) \phi(\eta', x') | 0 \rangle = \langle \phi \rangle + \int_{\frac{d}{2}-i\infty}^{\frac{d}{2}+i\infty} d\Delta \rho_\phi(\Delta) G_\Delta^{\text{free}}(\xi)$$

Massive free boson

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Bulk CFT

$$\langle \phi \phi \rangle = \zeta^\delta \quad \Rightarrow \quad \rho_\phi(\Delta = \frac{d}{2} + i\nu) \propto \nu \sinh(\pi\nu) \underbrace{\Gamma(\delta - \frac{d}{2} - i\nu) \Gamma(\delta - \frac{d}{2} + i\nu)}_{\text{poles at } \Delta = \delta + iN_6}$$

bulk scaling dimension

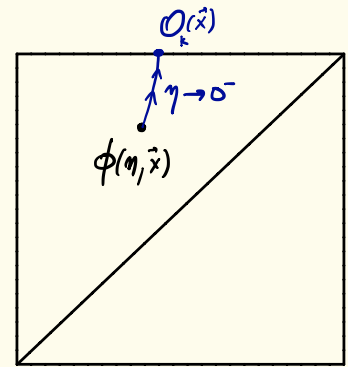
→ Need complementary series if $\frac{d-1}{2} < \delta < \frac{d}{2}$.

Boundary Operators

$$\phi(\eta, \vec{x}) = \sum_k b_{\phi k} (-\eta)^{\Delta_k} \left[\mathcal{O}_k(\vec{x}) + \text{desc.} \right]$$

\uparrow primary

$$\phi^{\dagger} = \phi \Rightarrow \begin{cases} \mathcal{O}_k^{\dagger} = \mathcal{O}_k, & \Delta_k \in \mathbb{R}, & b_{\phi k} \in \mathbb{R} \\ \mathcal{O}_k^{\dagger} = \mathcal{O}_{k'}, & \Delta_k^* = \Delta_{k'}, & b_{\phi k}^* = b_{\phi k'} \end{cases}$$

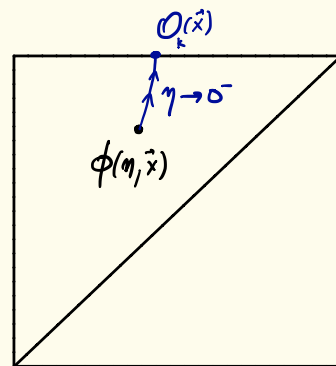


Boundary Operators

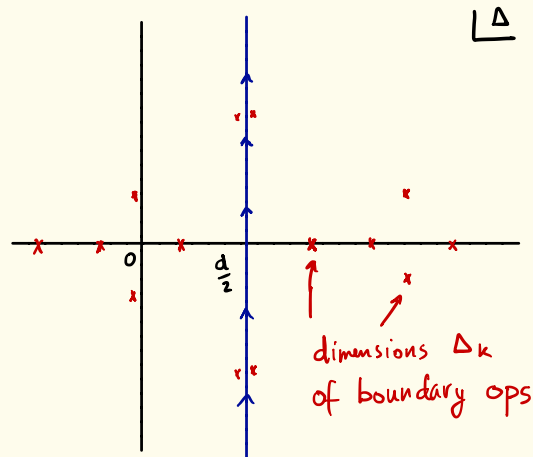
$$\phi(\eta, \vec{x}) = \sum_k b_{\phi k} (-\eta)^{\Delta_k} \left[\mathcal{O}_k(\vec{x}) + \text{desc.} \right]$$

\uparrow primary

$$\phi^\dagger = \phi \Rightarrow \begin{cases} \mathcal{O}_k^\dagger = \mathcal{O}_k, & \Delta_k \in \mathbb{R}, & b_{\phi k} \in \mathbb{R} \\ \mathcal{O}_k^\dagger = \mathcal{O}_{k'}, & \Delta_k^* = \Delta_{k'}, & b_{\phi k}^* = b_{\phi k'} \end{cases}$$



$$\langle \phi(\eta_1, \vec{x}_1) \phi(\eta_2, \vec{x}_2) \rangle = \int_{\frac{d}{2}-i\infty}^{\frac{d}{2}+i\infty} d\Delta \rho_\phi(\Delta) \underbrace{G_\Delta^{\text{free}}(\xi)}_{\substack{\downarrow \xi \rightarrow 0 \Leftrightarrow \eta_i \rightarrow 0 \\ \xi^\Delta(\dots) + \xi^{d-\Delta}(\dots)}}$$



No state-operator map!

Boundary 4pt - function

$$\langle \mathcal{O}(x_1) \mathcal{O}(x_2) \mathcal{O}(x_3) \mathcal{O}(x_4) \rangle = \sum_{\Delta, \ell} \int d\Delta \, I_{\Delta, \ell}$$

$$\mathbb{1} = \sum_{\Delta, \ell} \int d\Delta \int d^d y \, |\Delta, \ell\rangle \langle \Delta, \ell| + \dots$$

Unitarity \Rightarrow

$$I_{\Delta, \ell} \geq 0$$

$$\underbrace{\Psi_{\Delta, \ell}^{12,34}(x_1, x_2, x_3, x_4)} + \langle \mathcal{O}_1 \mathcal{O}_2 \rangle \langle \mathcal{O}_3 \mathcal{O}_4 \rangle$$

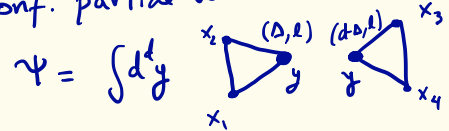
Conf. partial wave

$$\Psi = \int d^d y \, \begin{array}{c} x_2 \quad (\Delta, \ell) \\ \diagdown \quad \diagup \\ x_1 \quad y \end{array} \quad \begin{array}{c} (\Delta, \ell) \quad x_3 \\ \diagup \quad \diagdown \\ y \quad x_4 \end{array}$$

Boundary 4pt-function

$$\langle \mathcal{O}(x_1) \mathcal{O}(x_2) \mathcal{O}(x_3) \mathcal{O}(x_4) \rangle = \sum_{\ell} \int d\Delta \, I_{\Delta, \ell} \underbrace{\Psi_{\Delta, \ell}^{12,34}(x_1, x_2, x_3, x_4)}_{\text{Conf. partial wave}} + \langle \mathcal{O}_1 \mathcal{O}_2 \rangle \langle \mathcal{O}_3 \mathcal{O}_4 \rangle$$

Unitarity \Rightarrow $I_{\Delta, \ell} \geq 0$



Example: free scalar field in dS_{d+1} : $\Delta = \frac{d}{2} + i\mu$, $\mu \in \mathbb{R}$

$$\langle \mathcal{O}(x_1) \mathcal{O}^\dagger(x_2) \mathcal{O}^\dagger(x_3) \mathcal{O}(x_4) \rangle = \frac{1}{(x_{14}^2)^\Delta (x_{23}^2)^\Delta} + c \delta(x_{13}) \delta(x_{24}) + c \delta(x_{12}) \delta(x_{34})$$

$$\Rightarrow I_{\Delta, \ell} \geq 0 \quad \checkmark$$

$$\lambda \phi^4 \text{ X} \rightarrow I_{\Delta, \ell} = I_{\Delta, \ell}^{(0)} + \lambda I_{\Delta, \ell}^{(1)} + \mathcal{O}(\lambda^2)$$

Crossing + Unitarity → Bootstrap

$d=1$ QFT in dS_2

$$\langle \mathcal{O}(x_1) \mathcal{O}^\dagger(x_2) \mathcal{O}(x_3) \mathcal{O}^\dagger(x_4) \rangle = \frac{1}{(x_{12}^2 x_{34}^2)^{\Delta_{re}}} \left| \frac{x_{24}}{x_{13}} \right|^{2i\Delta_{im}}$$

\uparrow
 $\Delta_{\mathcal{O}} = \Delta_{re} + i\Delta_{im}$

unique cross ratio

$$\mathcal{G}(z) \equiv \frac{x_{12} x_{34}}{x_{13} x_{24}}$$

$d=1$ QFT in dS_2

$$\langle \mathcal{O}(x_1) \mathcal{O}^\dagger(x_2) \mathcal{O}(x_3) \mathcal{O}^\dagger(x_4) \rangle = \frac{1}{(x_{12}^2 x_{34}^2)^{\Delta_{re}}} \left| \frac{x_{24}}{x_{13}} \right|^{2i\Delta_{im}} \mathcal{G}(z \equiv \frac{x_{12} x_{34}}{x_{13} x_{24}})$$

\uparrow
 $\Delta_{\mathcal{O}} = \Delta_{re} + i\Delta_{im}$

Crossing :
 $x_1 \leftrightarrow x_3$

$$(1-z)^{2\Delta_{re}} \mathcal{G}(z) = z^{2\Delta_{re}} \mathcal{G}(1-z)$$

$$\mathcal{G}(z) = \underbrace{\sum_{\ell=0,1} \int_0^\infty \frac{dv}{2\pi} I_{\frac{1}{2}+iv,\ell} \Psi_{\frac{1}{2}+iv,\ell}(z)}_{\text{principal series}} + \underbrace{\sum_{\substack{n \in \mathbb{N} \\ \ell=0,1}} \tilde{I}_{n,\ell} \Psi_{n,\ell}(z)}_{\text{discrete series}}$$

positive

$d=1$ QFT in dS_2

$$\langle \mathcal{O}(x_1) \mathcal{O}^\dagger(x_2) \mathcal{O}(x_3) \mathcal{O}^\dagger(x_4) \rangle = \frac{1}{(x_{12}^2 x_{34}^2)^{\Delta_{re}}} \left| \frac{x_{24}}{x_{13}} \right|^{2i\Delta_{im}} \mathcal{G}(z \equiv \frac{x_{12} x_{34}}{x_{13} x_{24}})$$

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positive

Problem: Not absolutely convergent.

Solution: Act with linear functional: $\omega[f] = \int_0^1 dz z^\sigma (1-z)^\sigma f(z)$

Regularized Crossing Equations

$$\sum_{\ell=0,1} \int_0^{\infty} \frac{dv}{2\pi} I_{\frac{1}{2}+iv,\ell} \tilde{F}_{\frac{1}{2}+iv,\ell}(\tau,\sigma) + \sum_{\ell=0,1} \sum_{n=1}^{\infty} \tilde{I}_{n,\ell} \tilde{F}_{n,\ell}(\tau,\sigma) = 0 \quad \forall \tau,\sigma > -1$$

Example of bound: Take $\mathcal{O} = \mathcal{O}^+$ with $\Delta_{\mathcal{O}} = \frac{1}{2} + \frac{1}{8} \Rightarrow l=0$ & n even

$$\int_0^{\infty} \frac{dv}{2\pi} I_{\frac{1}{2}+iv,0} \tilde{F}_{\frac{1}{2}+iv,0}(\tau,\sigma) + \sum_{\substack{n=2 \\ \text{even}}}^{\infty} \tilde{I}_{n,0} \tilde{F}_{n,0}(\tau,\sigma) + D(\tau,\sigma) = 0$$

↙ disconnected term

If $I_{\frac{1}{2}+iv,0} = 0$ for $v < 8.53$ then $\tilde{I}_{2,0} < 5.7$

QFT on:	AdS_{d+1}	IM^{d+1}	dS_{d+1}
<u>Observable</u>	boundary φ . $\langle \mathcal{O}(x_1) \mathcal{O}(x_2) \mathcal{O}(x_3) \mathcal{O}(x_4) \rangle$ $= \frac{1}{(x_{13}^2 x_{24}^2)^{\Delta_0}}$ $F(u, v)$ <small>↑ ↑ Cross-ratios</small>	$\langle P_3, P_4 P_1, P_2 \rangle_{in}^{Conn} =$ $= i \delta(\sum P_i) T(s, t)$ <small>↑ ↑ Mandelstam inv. $s+t+u=4m^2$</small>	boundary φ . $\langle \mathcal{O}(x_1) \mathcal{O}(x_2) \mathcal{O}(x_3) \mathcal{O}(x_4) \rangle$ $= \frac{1}{(x_{13}^2 x_{24}^2)^{\Delta_0}}$ $F(u, v)$ <small>↑ ↑ Cross-ratios</small>
<u>Crossing</u>	$F(u, v) = F(v, u)$ $= \frac{F(\frac{u}{v}, \frac{1}{v})}{u^{\Delta_0}}$	$T(s, t) = T(t, s)$ $= T(\frac{4m^2 - s - t}{u}, t)$	$F(u, v) = F(v, u)$ $= \frac{F(\frac{u}{v}, \frac{1}{v})}{u^{\Delta_0}}$
<u>Decomposition in irreps</u>	$F(u, v) = \sum_{\Delta, \ell} \lambda_{\Delta, \ell} g_{\Delta, \ell}(u, v)$ <small>↑ conf. blocks</small>	$T \propto \sum_{\ell} f_{\ell}(s) P_{\ell}(\underbrace{\cos \theta}_{1 - \frac{2t}{4m^2 - s}})$	$F(u, v) = \sum_{\ell} \int d\Delta I_{\Delta, \ell} \Psi_{\Delta, \ell}(u, v)$ <small>↑ CPWs</small>
<u>Unitarity</u>	$\lambda_{\Delta, \ell} \geq 0$	$ 1 + i f_{\ell}(s) ^2 \leq 1$	$I_{\Delta, \ell} \geq 0$

Open Questions

Open Questions

For a generic QFT in dS , what is the:

- Hilbert space?

- Spectrum of boundary ops.?

Computable examples: $\left\{ \begin{array}{l} \text{free QFT} \\ \text{bulk CFT} \end{array} \right. + \text{small deformations}$

What are the relevant quantities to bound?

Better numerical bootstrap methods? " $I_{\Delta, \ell} = \sum_{\ell'} \int d\Delta' \mathcal{J}(\Delta, \Delta', \ell, \ell') I_{\Delta', \ell'}$ "

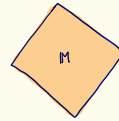
Flat space limit?

Conformal Bootstrap



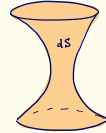
$R \rightarrow 0$

S-matrix Bootstrap



$R \rightarrow 0$

Cosmological Bootstrap



Can we include dynamical gravity?

Thank you !