Accreting neutron stars: the potential third MeV astrophysical neutrino source







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2/27 Neutrino astronomy



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"Neutrino astronomy is interesting for the same reason it is difficult."

John Bahcall, 1989

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Neutrino astronomy

Neutrino detection is *difficult*. But, precisely because of that, neutrinos

- Provide us with the conditions deep inside sources.
- Carry out *real-time* tracking.
- Are not affected by propagation.

Neutrinos interact weakly. But, precisely because of that, neutrinos

- Are a clean probe of new, weak, interactions.
- Are a clean probe of accumulating propagation effects.
- Are a clean probe of large densities.

4/27 Neutrino astronomy







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Neutrino astronomy: successes in astrophysics

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Article | Published: 24 October 2018

Comprehensive measurement of *pp*-chain solar neutrinos

The Borexino Collaboration

Nature 562, 505-510 (2018) Cite this article

12k Accesses | 164 Citations | 270 Altmetric | Metrics

Abstract

About 99 per cent of solar energy is produced through sequences of nuclear reactions that convert hydrogen into helium, starting from the fusion of two protons (the pp chain). The neutrinos emitted by five of these reactions represent a unique probe of the Sun's internal working and, at the same time, offer an intense natural neutrino beam for fundamental physics. Here we report a complete study of the pp chain. We measure the neutrino-electron

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JOURNAL ARTICLE

Helioseismic and neutrino data-driven reconstruction of solar properties @ Ningiang Song @, (Gonzalez-Garcia, Francesco L.Villante, Nuria Vinyoles, Aldo Serenelli @

Monthly Notices of the Royal Astronomical Society, Volume 477, Issue 1, June 2018, Pages 1397–1413, https://doi.org/10.1093/mnras/sty600 Published: 06 March 2018 Article history +

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Abstract

In this work, we use Bayesian inference to quantitatively reconstruct the solar properties most relevant to the solar composition problem using as inputs the information provided by helioseismic and solar neutrino data. In particular, we

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PHYSICAL REVIEW D

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15 JULY 1988

Observation in the Kamiokande-II detector of the neutrino burst from supernova SN1987A

The properties of the Kamiokande-II detector and the method of measurement are described in detail. The data on the neutrino burst from the supernova SN1987A on 23 February 1987 at 7:35:35 UT \pm 1 min are presented, with records of earlier and later observation periods in which other neutrino events possibly associated with SN1987A might have occurred. There is no evidence in the data for any excess of neutrino-induced events, either in a burst of a few seconds duration or over a longer time interval, relative to the usual count rate, excepting only the neutrino burst at 7:35:35 UT. The nature of the single, observed neutrino burst coincides remarkably well with the elements of the current model of type-II supernovae and neutron-star formation. This is the first direct observation in neutrino astronomy.

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BLAZARS AS ULTRA-HIGH-ENERGY COSMIC-RAY SOURCES: IMPLICATIONS FOR TeV GAMMA-RAY OBSERVATIONS

KOHTA, MURASH¹, CHARLES D, DERMIR², HAIME TAKAM¹, AND GIULIA MGUDRI⁴ ¹ Department of Physics, Center for Comology and Autro Particle Physics. The Ohio Sate Livewsity, Columbus, OH 43210, USA ² Space Science Division, Naval Research Laboratory, Washington, DC 20175, USA ³ Max Pinaci, Iostinico Forphysics, Folinger Ring, Carl, 80085 Minich, Germany ⁴ Havards-mithkomian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138, USA *Received 2011 July 27: accepted 2012 January 292*, published 2012 March 23 C

ABSTRACT

The spectra of BL Lac objects and Fanaroff-Riley I radio galaxies are commonly explained by the one-zone leptonic synchrotron self-Compton (SSC) model. Spectral modeling of correlated multiwavelength data gives the

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Neutrino astronomy: successes in astrophysics

Clean, real-time probes of dense interiors of sources.

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Neutrino astronomy: successes in particle physics





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^{10/27} Neutrino astronomy: successes in particle physics

Clean probes of new interactions, dense matter, and/or propagation effects.

^{11/27} See Ivanova et al, 1209.4302, for a review



A process where two stars share a common envelope, due to expansion or orbital decrease.

Suggested by Paczynski in 1976. Sounds exotic, but most likely necessary for

- Type la-supernovae
- X-ray binaries
- Gravitational-wave sources

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Yet still **unobserved** unambiguously (though see Dong et al., 2021, Science), and **very challenging to simulate!** We need new signatures.

From Ivanova et al, 1209.4302



- For the rest of the talk, I will focus on common envelope with *neutron stars* (relevant for, e.g., X-ray binaries or gravitational wave sources).
- Roughly, the main processes that happen during common envelope are
 - The neutron star **inspirals** due to drag.
 - The neutron star **accretes** material.

14/27 Accretion

The neutron star accretion rate can be estimated [Hoyle, Littleton (1939); Bondi (1952)]

$$\dot{M} \sim rac{G^2 M_{
m NS}^2
ho}{(
u^2 + c_s^2)^{3/2}} \sim 10^2 \, M_{
m sun}/{
m year} \ \sim 10^{25} \, {
m kg/s}$$

[This is an approximation, I'll lift it in a few slides]



^{15/27} Super-Eddington accretion

- Usually, accretion is limited because [Eddington (1926)]
 - Inflow gains kinetic energy, heats up, and radiation pressure can compensate gravity.
 - **2** That kinetic energy must go somewhere.
- Eddington limit! $\sim 10^{-8} \ensuremath{M_{\rm sun}}/{\rm year}$ for a neutron star.

How can we violate it by orders of magnitude in common envelope?

^{16/27} Super-Eddington accretion



Trapped light due to large inflow velocities

Basics well-understood since the 90s [Chevalier (1989); Houck & Chevalier (1991); Chevalier (1993)].

^{16/27} Super-Eddington accretion



Trapped light due to large inflow velocities

Neutrino cooling

Basics well-understood since the 90s [Chevalier (1989); Houck & Chevalier (1991); Chevalier (1993)].

^{17/27} Super-Eddington accretion

The main uncertainty is the accretion rate. Simulations are hard (10 km neutron star vs 100 $R_{
m sun} \sim 10^8 \, {
m km}$ giant)

- Angular momentum?
- How much of the energy is dissipated by neutrinos? Are jets formed?
- What are the details of the onset of super-Eddington accretion?

Fragos et al, 2019; Ivanova et al, 2012; Ricker et al, 2011; Houck & Chevalier, 1991; Brown, 1995; Ricket & Taam, 2012; Macleod & Ramirez-Ruiz, 2014; Macleod et al, 2017; Brown et al, 2000; ...

Yet recent 3D simulations still find super-Eddington accretion [Macleod & Ramirez-Ruiz, 2014; Hutchinson-Smith et al, 2023; Everson et al, 2023], with $\dot{M} \sim 0.1 M_{\rm sun}/{\rm year}$.

^{17/27} Super-Eddington accretion





re hard

Are jets formed?

tion [Macleod & on et al, 2023],

^{18/27} Other properties

Duration? The neutron star momentum loss (due to drag) can be estimated by the linear momentum gained by the inflow, Macleod & Ramirez-Ruiz, 2014...

$$t \sim {f month} imes \left({0.1 \, M_{
m sun}/
m year \over \dot M}
ight)$$

Rate in the Milky Way? From X-ray binary catalogs, estimates are Ginat et al, 2019; Hutilukejiang et al, 2018

$$au \sim 10^{-2}$$
 – $1\,{
m century}^{-1}$

^{19/27} Wrapping up

- Understanding common envelope is **key** to understand X-ray binaries, gravitational-wave sources. . .
- But direct, unambiguous observational windows lack.

Super-Eddington accretion, hypothesized since the 90s, would involve

- Accretion rates $\sim 0.1 M_{\rm sun}/{\rm year}$.
- For about a month.
- $\blacksquare \sim 10^{-2} 1 \, \rm century^{-1}$ in our galaxy.
- With neutrinos playing a key role.
- How can we look for this?

^{19/27} Wrapping up

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0/27 Average neutrino energy

Esteban, Beacom, Kopp; 2310.19868

- What are the properties of the neutrino signal?
- Neutrinos dissipate the kinetic energy gained by the inflow. Simple energy conservation determines the signal properties.

Inflow heating
$$\Rightarrow \frac{\mathrm{d}E}{\mathrm{d}t\mathrm{d}V} \sim \frac{GM_{\mathrm{NS}}\dot{M}}{r_{\mathrm{NS}}} \frac{1}{4\pi r_{\mathrm{NS}}^2(r_{\mathrm{NS}}/2)}$$

Neutrino-cooling $\Rightarrow \frac{\mathrm{d}E}{\mathrm{d}t\mathrm{d}V} \sim G_F^2 T^9$. So $T \sim \left(\frac{\mathrm{d}E}{\mathrm{d}t\mathrm{d}V}\right)^{1/9}$



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^{20/27} Average neutrino energy

Esteban, Beacom, Kopp; 2310.19868

- What are the properties of the neutrino signal?
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^{21/27} Neutrino flux

Esteban, Beacom, Kopp; 2310.19868

Neutrinos dissipate the kinetic energy gained by the inflow. Simple energy conservation determines the signal properties. $\langle E_{\nu} \rangle \sim 4 \text{ MeV}.$

$$rac{\mathrm{d}N_{
u}}{\mathrm{d}t}\sim rac{\mathrm{d}E/\mathrm{d}t}{\langle E_{
u}
angle}\sim rac{GM_{\mathrm{NS}}\dot{M}/r_{\mathrm{NS}}}{\langle E_{
u}
angle}\sim 10^{50}\,\mathrm{neutrinos/s}$$

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- ^{22/27} Order-of-magnitude estimations
 - Esteban, Beacom, Kopp; 2310.19868
 - Neutrinos dissipate the kinetic energy gained by the inflow. Simple energy conservation determines the signal properties. $\langle E_{\nu} \rangle \sim 4 \text{ MeV}, \ 10^{50} \text{ neutrinos/s}.$
 - At 10 kpc, $\phi_{\nu} \sim 10^4 \, {\rm cm}^{-2} \, {\rm s}^{-1}$. Is it observable? At \sim MeV energies, the most efficient detection channel is

$$\bar{\nu}_e + p \Rightarrow n + e^+$$

with $\sigma \sim 10^{-42} \,\mathrm{cm}^2$. $\phi \times \sigma \times N_{\mathrm{targets}} \sim 100 \,\mathrm{events/few months}$ at Super-Kamiokande! (Of course, beware of backgrounds!)

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^{13/27} Sharpening up the predictions

Up to now, I've only given (robust) order-of-magnitude estimations.

We sharpened them via simulation. Main input is accretion rate $[\langle E_{\nu} \rangle \propto \dot{M}^{1/9}; \phi_{\nu} \propto \dot{M}]$, that we took from the 3D simulation Macleod & Ramirez-Ruiz, 2014 $[\dot{M} \sim 0.1 M_{\rm sun}/{\rm year}]$. Our code is publicly available! github.com/ivan-esteban-phys/common-envelope-thermal \bigcirc



- 3-months integrated
- 10⁵⁰ $\nu/\mathrm{s} \times 3 \,\mathrm{months} \sim 10^{56} \nu$
- Very well approximated by FD, $T \sim 1.6 \text{ MeV}$ ($\langle E_{\nu} \rangle \sim 5 \text{ MeV}$).
- No oscillations in this plot, although in our results we included adiabatic oscillations (factor of few impact).

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^{4/27} Detector signals

Archival Super-Kamiokande data: e^+ and γ from neutron capture on H.



Collaboration-estimated background from DSNB search.

Detector signals

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Super-Kamiokande with low Gd: e^+ and γ from neutron capture on Gd.



Similar, collaboration-estimated, background sources, but higher efficiency.

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^{4/27} Detector signals

JUNO: e^+ and γ from neutron capture.



Collaboration-estimated background. Large efficiency due to scintillator.

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^{4/27} Detector signals

DUNE: ν on e^- , directional.



Background estimates from solar neutrino studies (Capozzi et al, 2018; Zhu, Li & Beacom, 2019).

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5/27 Distance reach

Improvable! (More Gd, dedicated background reduction strategies ...)





For Inverted Ordering, currently disfavored by \sim 2–3 sigma, distance reach worsens by a factor \sim 2.

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^{25/27} Distance reach

Improvable! (More Gd, dedicated background reduction strategies ...)





For Inverted Ordering, currently disfavored by \sim 2–3 sigma, distance reach worsens by a factor \sim 2.



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Neutron-star common-envelope is key to understand many binary systems. But it has never been observed unambiguously.

If accretion is super-Eddington, cooling proceeds emitting neutrinos.

 A Milky Way event (~ 10⁻²-1 century⁻¹) would produce detectable, months-long, MeV neutrino signals. The *third* MeV astrophysical neutrino source, after the Sun and supernovae. It is rare, but we won't see it unless we look for it!

Outlook

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In case of a detection,

- Common-envelope would be observationally established.
- Super-Eddington accretion would be observationally established. Even a non-detection with an electromagnetic/gravitational wave counterpart would be useful!

What's next?

- Astrophysics: more simulation. Accretion on white dwarfs?
- *Experiments*: dedicated searches, *including archival data*.
- *Theory*: neutrino properties? Particle physics?
- We extensively detail how to predict the neutrino flux. O