

Post-inflationary Axions:

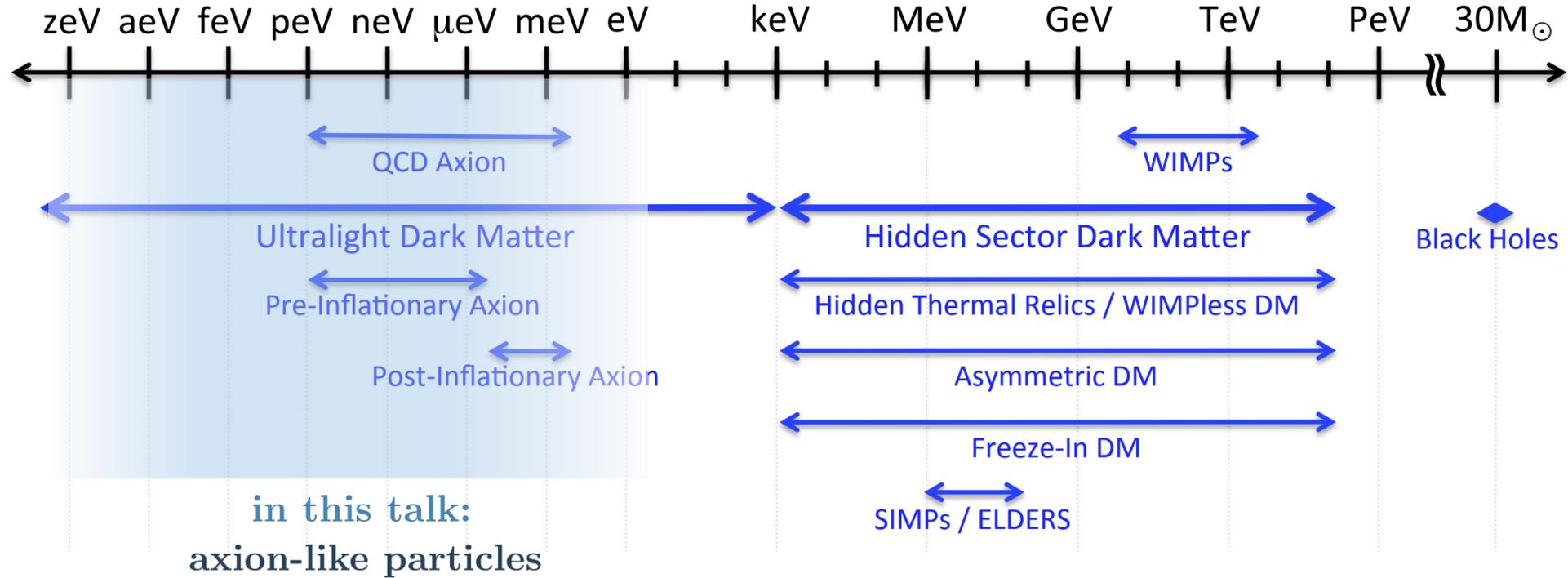
Gravitational waves and targets for haloscopes

Marco Gorghetto

WEIZMANN INSTITUTE OF SCIENCE

Mostly based on: [2212.13263] **MG, E.Hardy**
[2101.11007] **MG, E.Hardy, H.Nicoleascu; JCAP**
[2007.04990] **MG, E.Hardy, G.Villadoro; SciPost**

Dark Matter Candidates



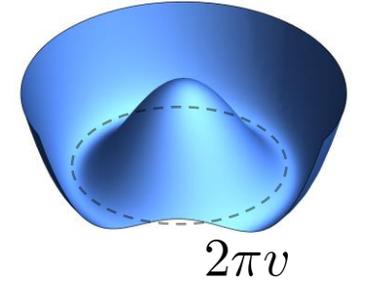
Axion (= ALP)

- PNGB of a new U(1) symmetry, spontaneously broken at a high scale v

$$\phi = \frac{r+v}{\sqrt{2}} e^{i\frac{a}{v}}$$

$$\mathcal{L}_\phi = |\partial_\mu \phi|^2 - \frac{m_r^2}{2v^2} \left(|\phi|^2 - \frac{v^2}{2} \right)^2$$

$$m_r \simeq v$$

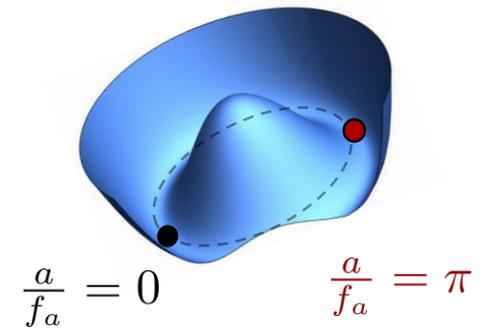
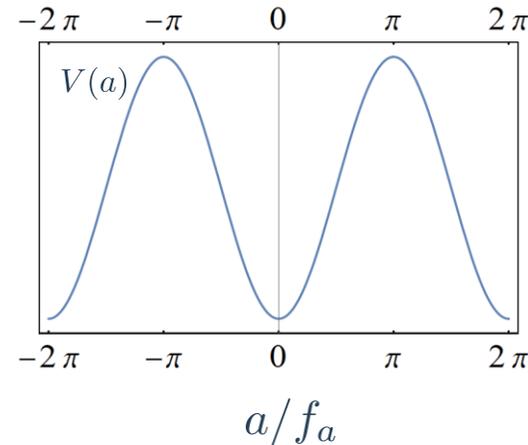


- U(1) broken by $V(a) \implies$ axion mass m_a

$\rightarrow V(a)$ invariant under $a \rightarrow a + 2\pi f_a$

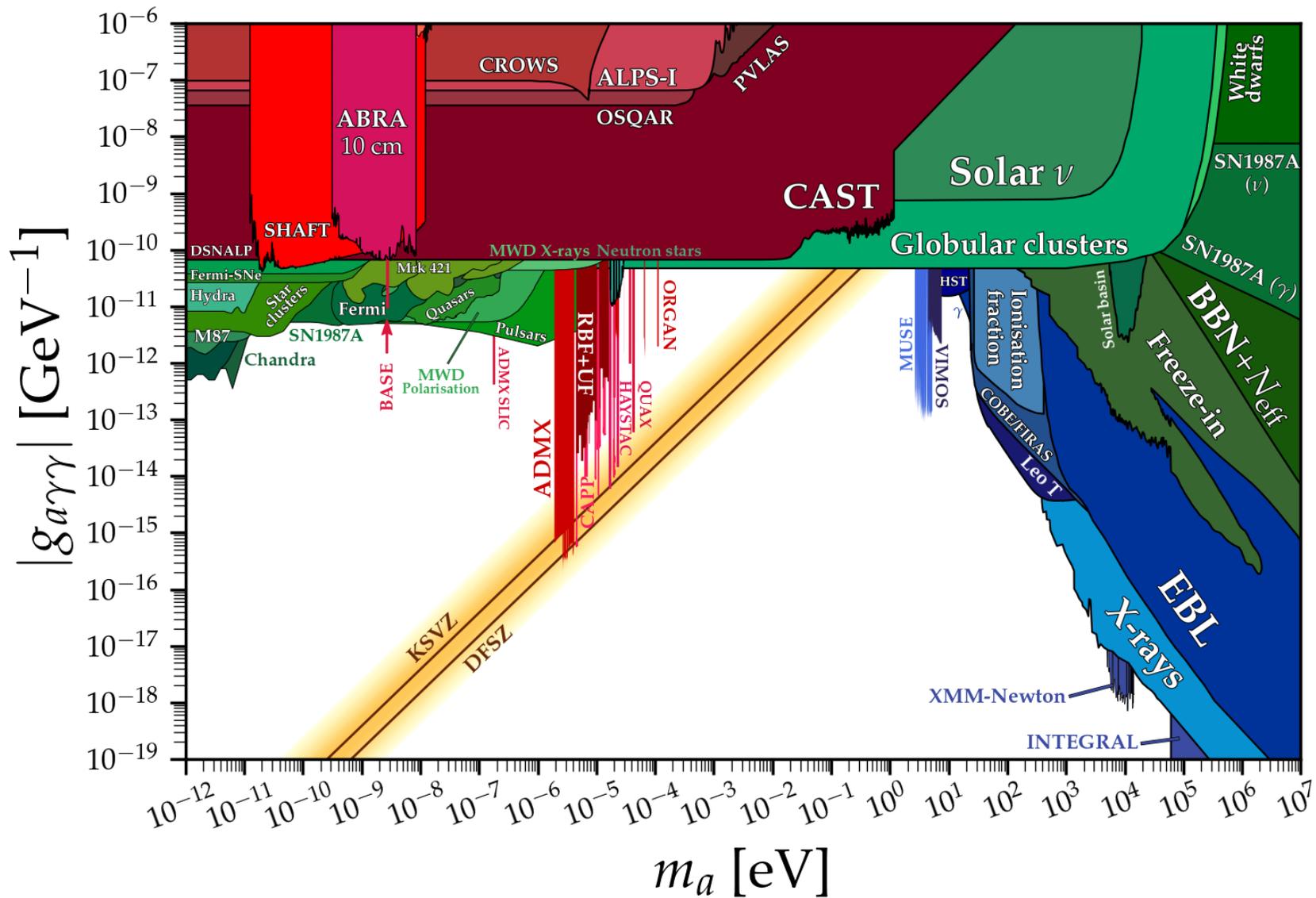
\rightarrow in general $v = N f_a$, N integer

\rightarrow for now assume $N = 1$, i.e. $v = f_a$



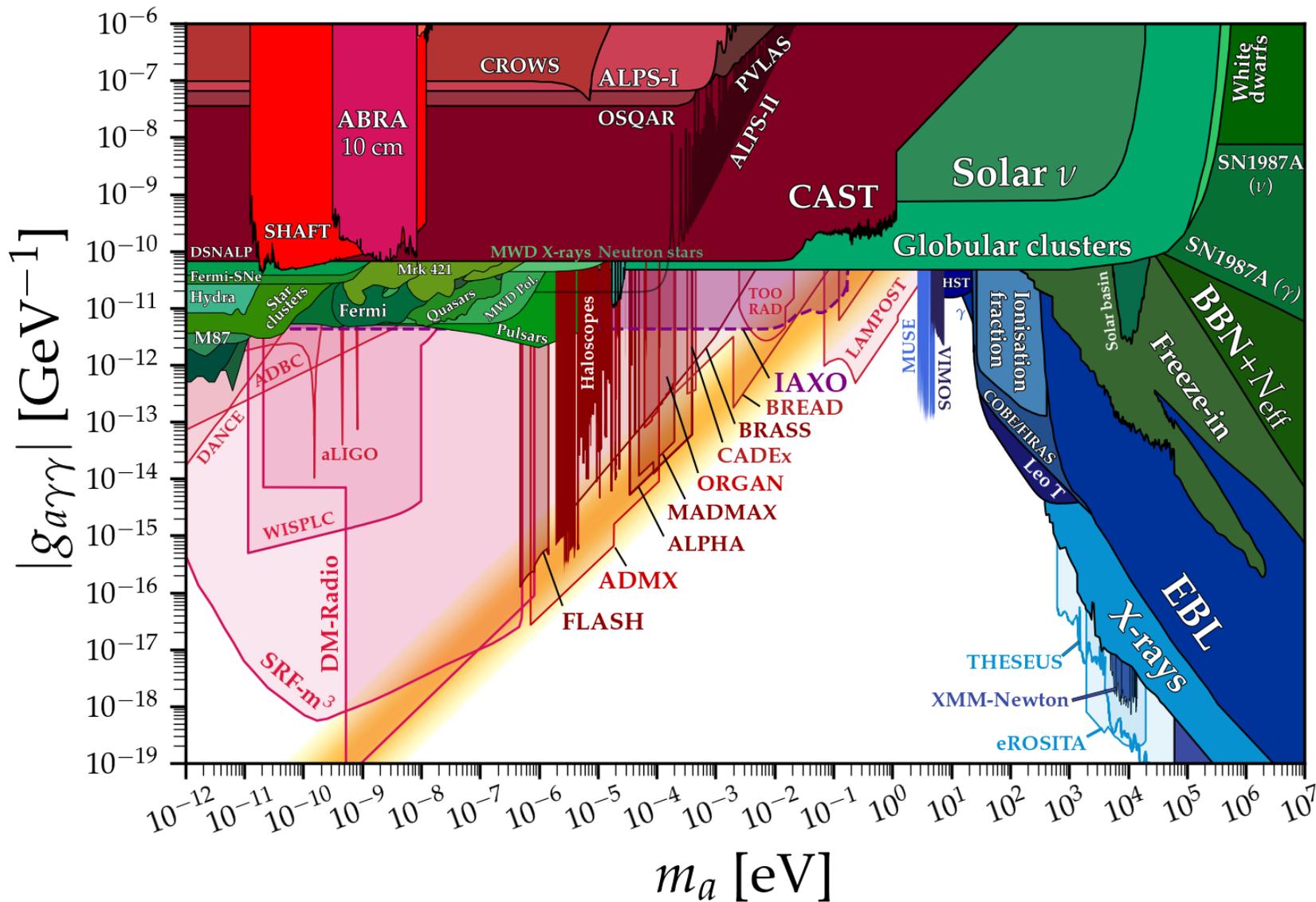
Axion searches

$$\mathcal{L} \supset -\frac{1}{4} g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$



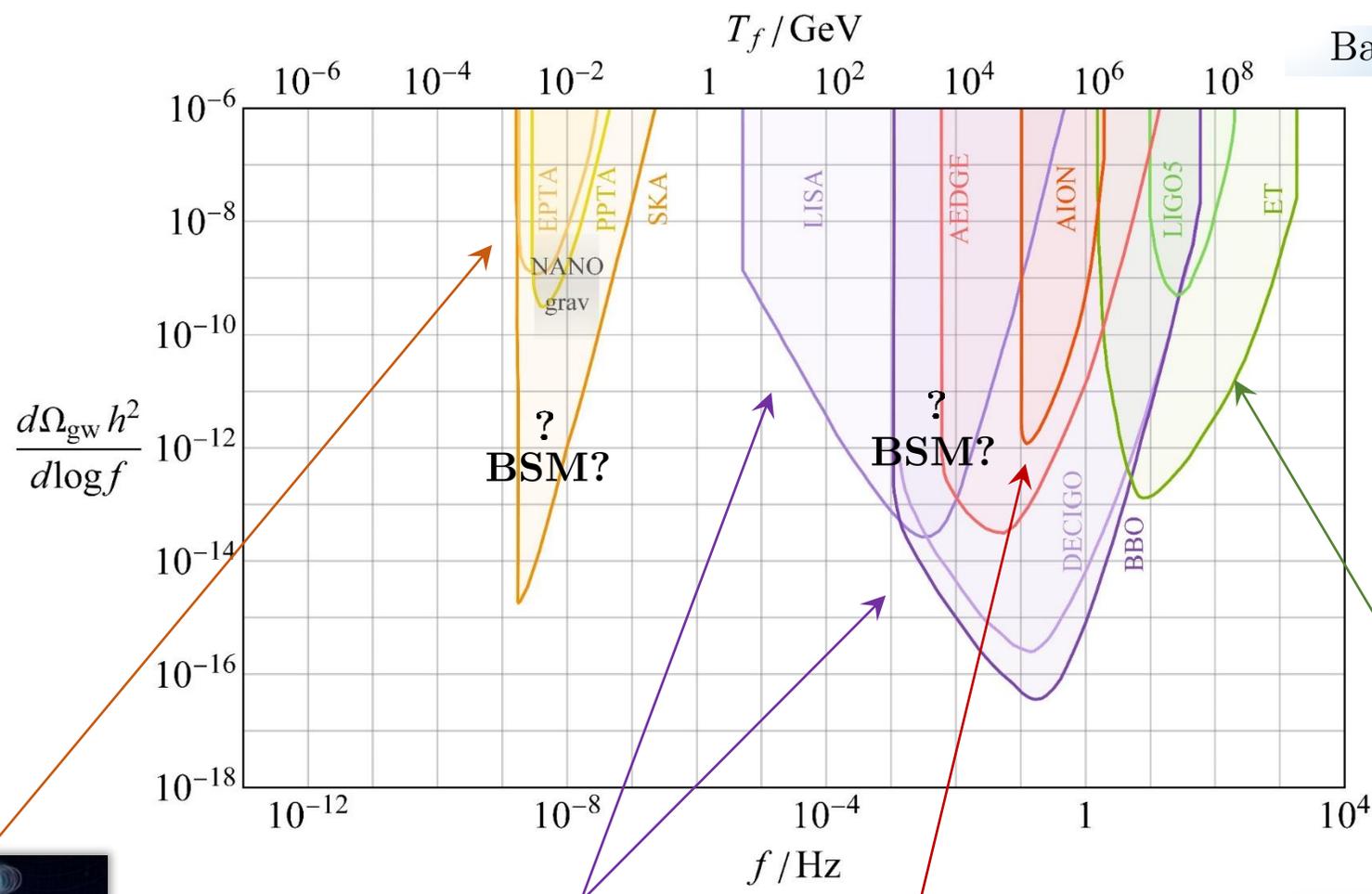
Axion searches

$$\mathcal{L} \supset -\frac{1}{4} g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$



GW searches

Back to the past 

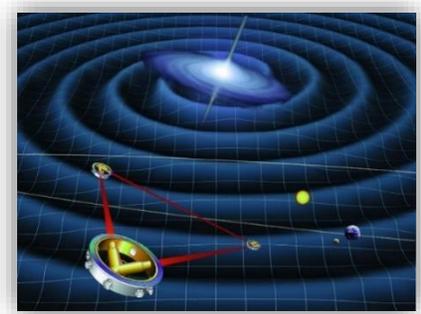


Caution:
Many experimental challenges to overcome

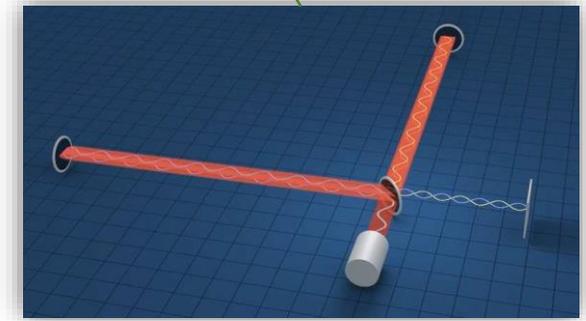
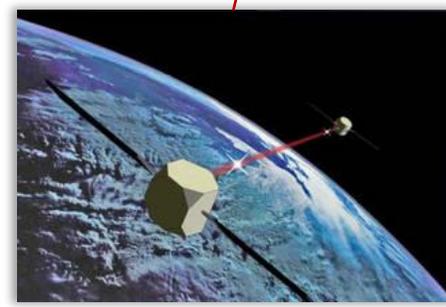
Pulsar



Pulsar Timing Arrays



Space-based Light/Atom Interferometers



Ground-based Interferometers

Outline

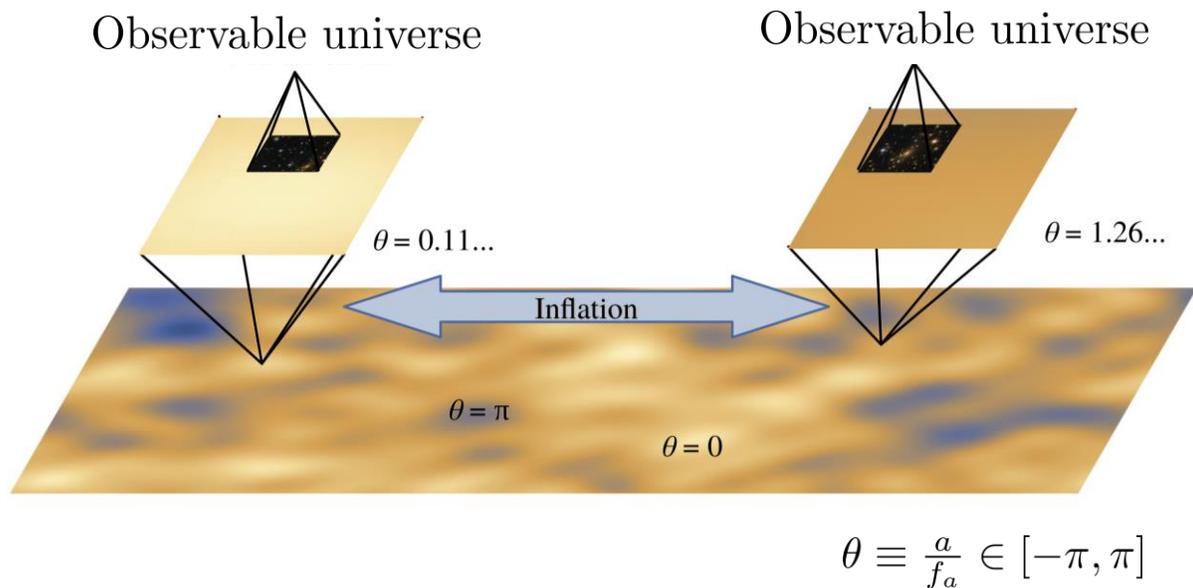
- Axion post-inflationary scenario
- Gravitational waves from axion strings
- $N > 1$: minimal targets for haloscopes & substructure

Cosmological Initial Conditions

Pre-inflationary scenario

$$T_R \lesssim f_a \text{ and } H_I \lesssim f_a$$

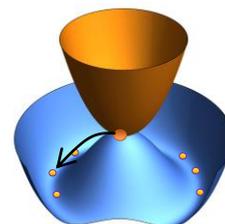
↑ Reheating temperature ↑ Hubble parameter during inflation



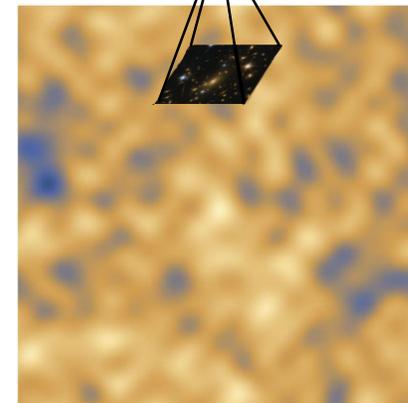
Post-inflationary scenario

$$T_R \gtrsim f_a \text{ or } H_I \gtrsim f_a$$

$$T \gtrsim f_a$$
$$T \lesssim f_a$$

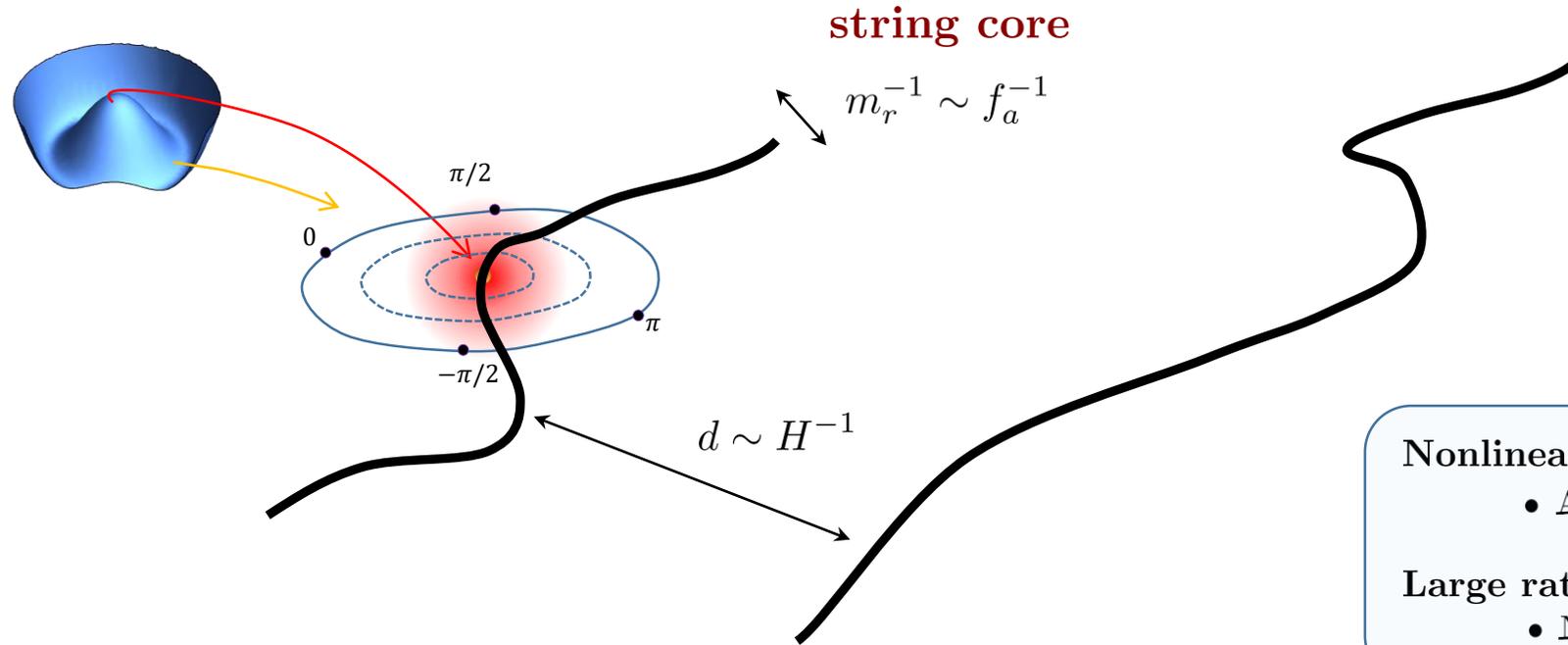


Observable universe



focus of this talk

Strings



Nonlinear dynamics:

- ~~Analytical approach~~ 😞

Large ratio of scales:

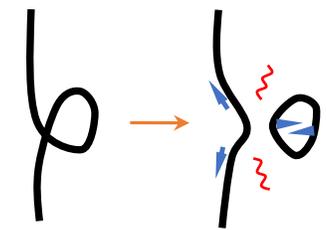
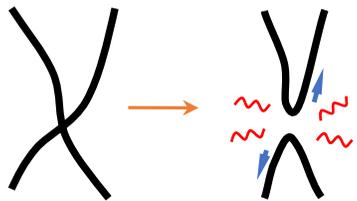
- ~~Numerical approach~~ 😞

string tension

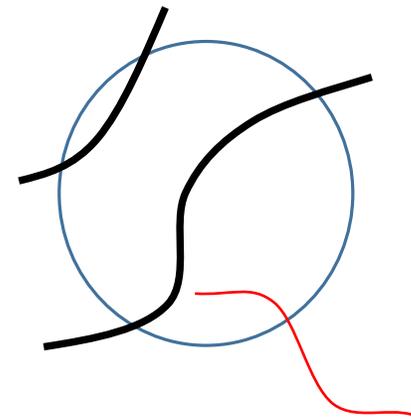
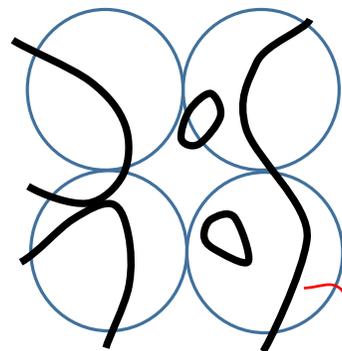
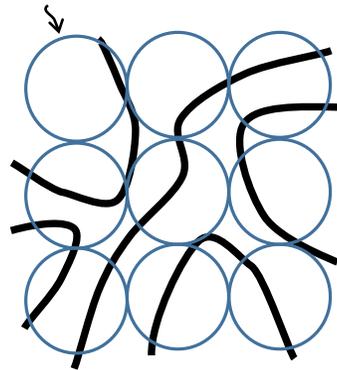
$$\mu = \frac{E}{L} \sim \underbrace{\pi f_a^2}_{\text{core}} \underbrace{\log \frac{d}{m_r^{-1}}}_{\text{axion gradient}} \sim \pi f_a^2 \log \frac{m_r}{H}$$

\uparrow
 grows logarithmically in time
 \downarrow
 T^2/M_p

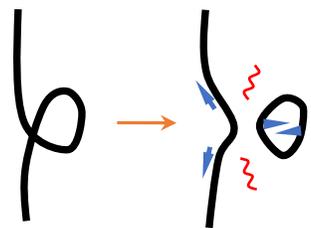
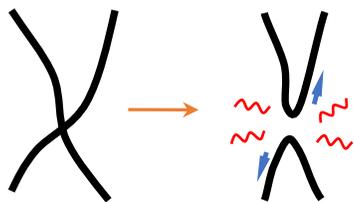
The Scaling Regime



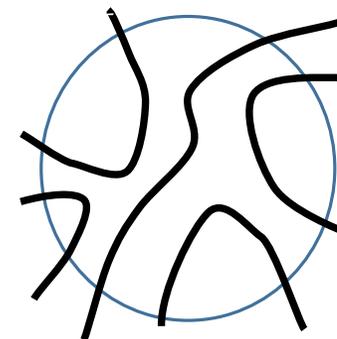
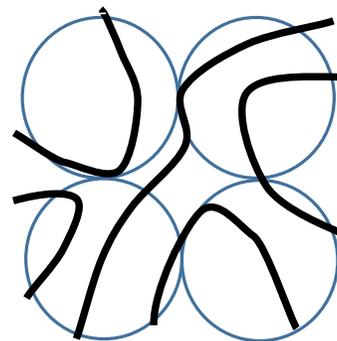
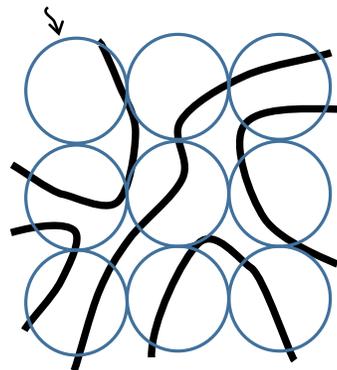
causal patch $\propto 1/H = 2t$



The Scaling Regime

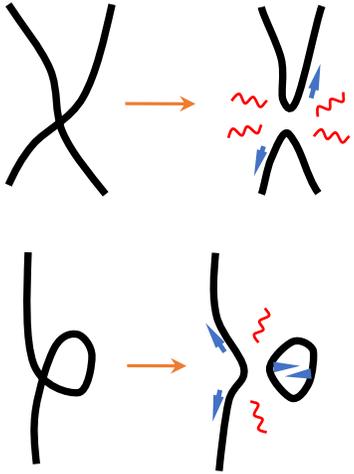


causal patch $\propto 1/H = 2t$

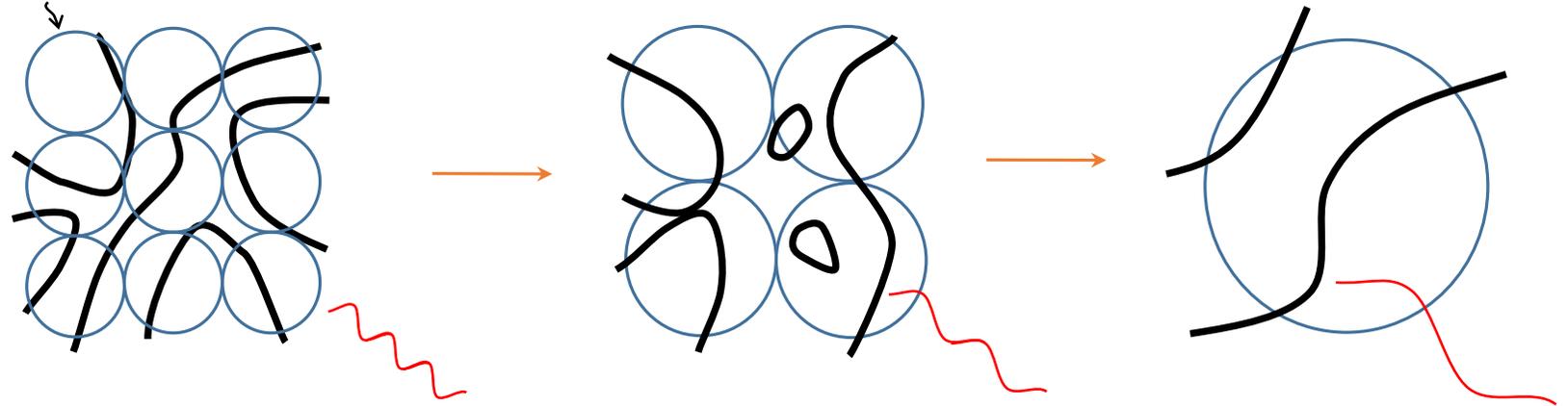


free strings: $\rho^{\text{free}} \propto \frac{1}{R^2} \propto \frac{1}{t}$

The Scaling Regime



causal patch $\propto 1/H = 2t$



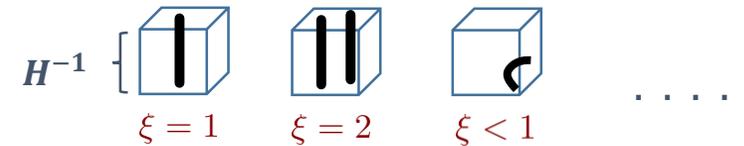
rate of energy loss:

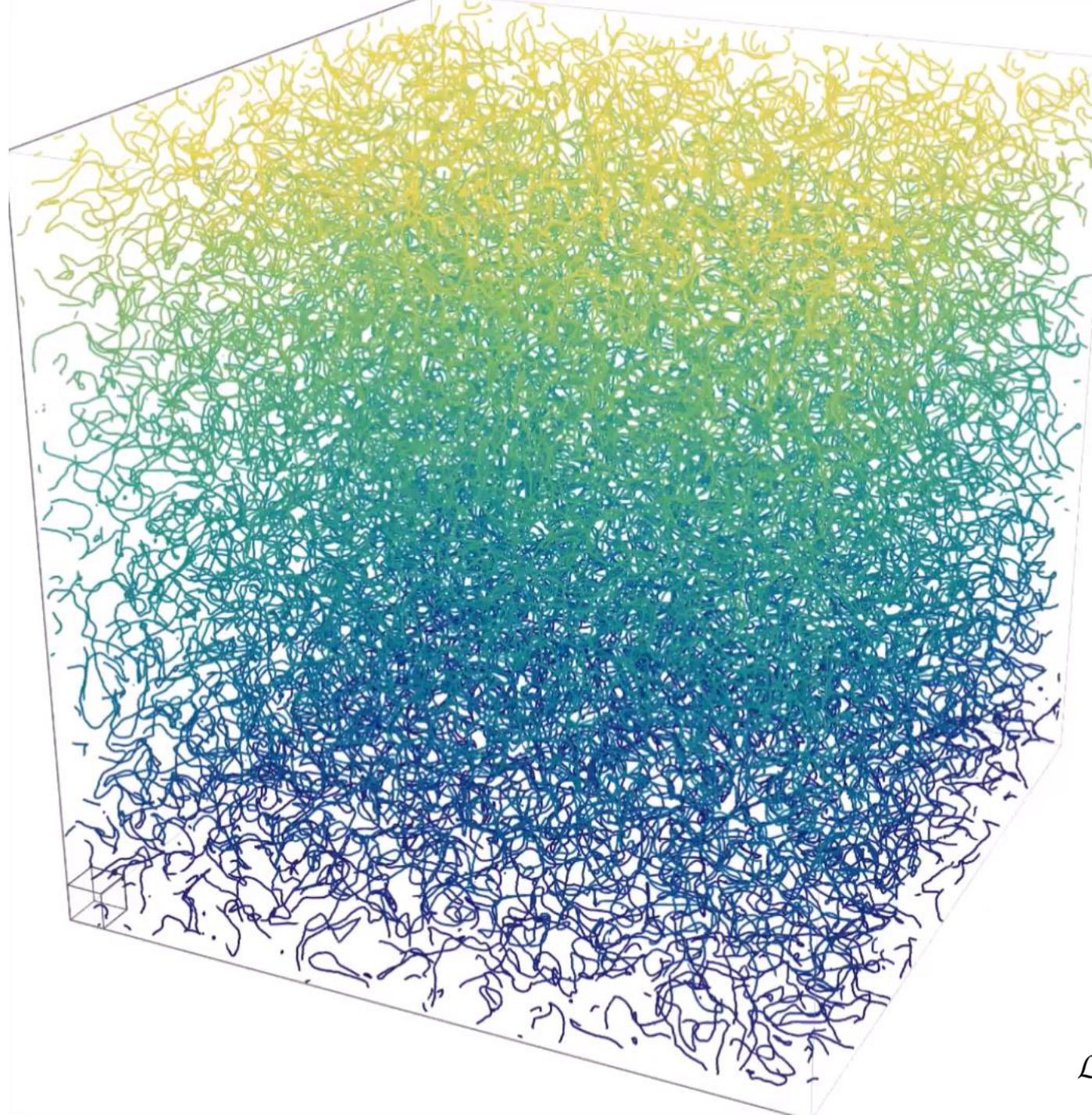
$$\Gamma \equiv \frac{d}{dt} [\rho^{\text{free}} - \rho^{\text{scal}}] \simeq \frac{\xi \mu}{t^3}$$

$$\propto \frac{1}{R^2} \propto \frac{1}{t}$$

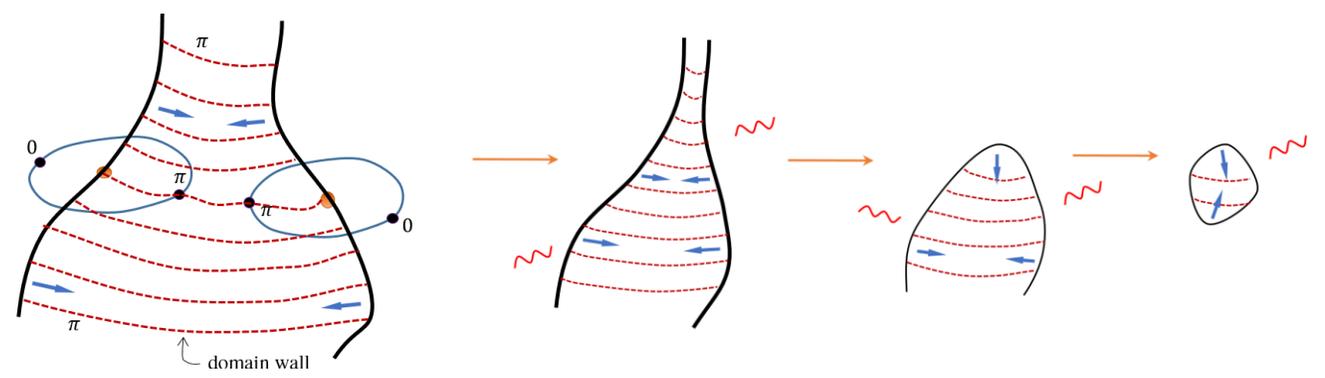
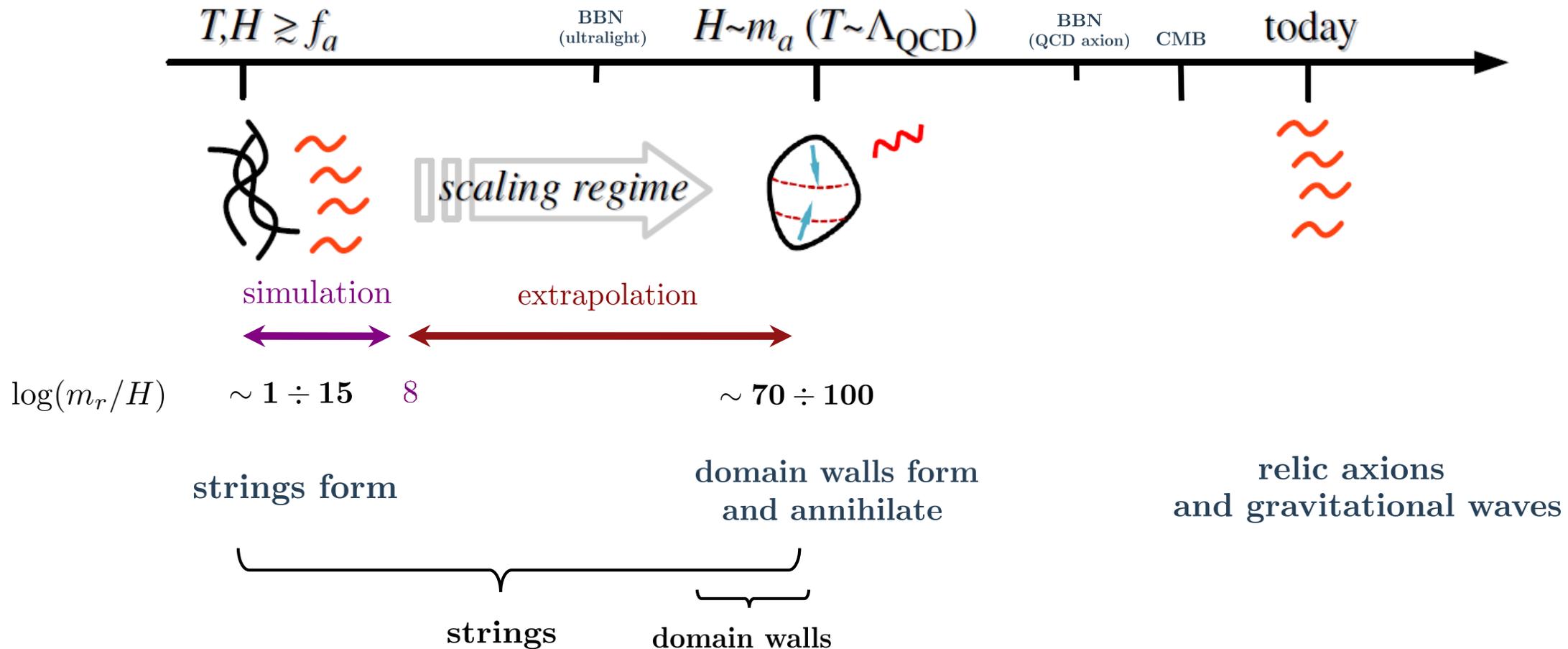
$$\boxed{\xi} \frac{\mu}{t^2}$$

number of strings
per Hubble patch

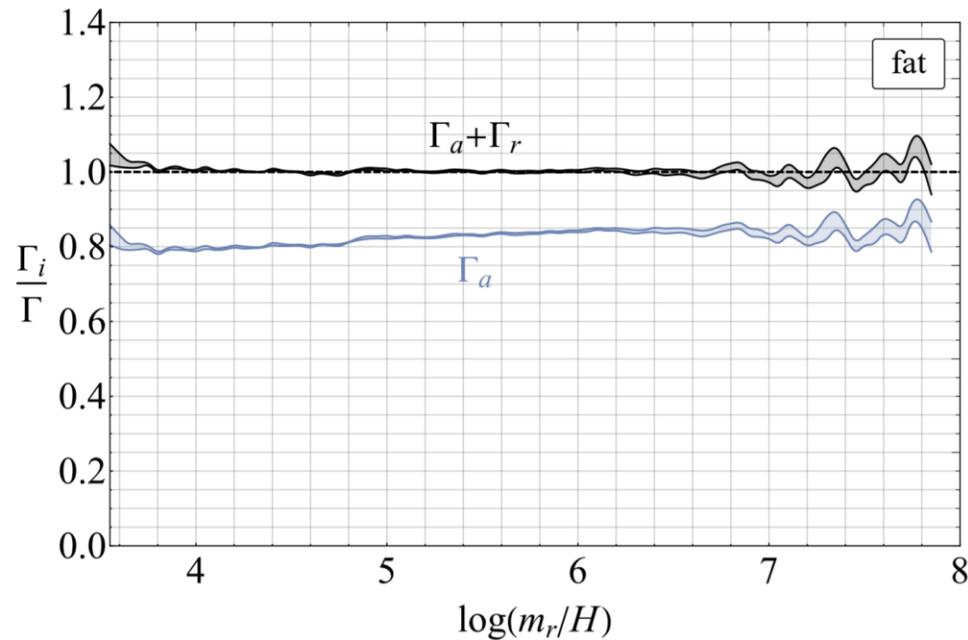
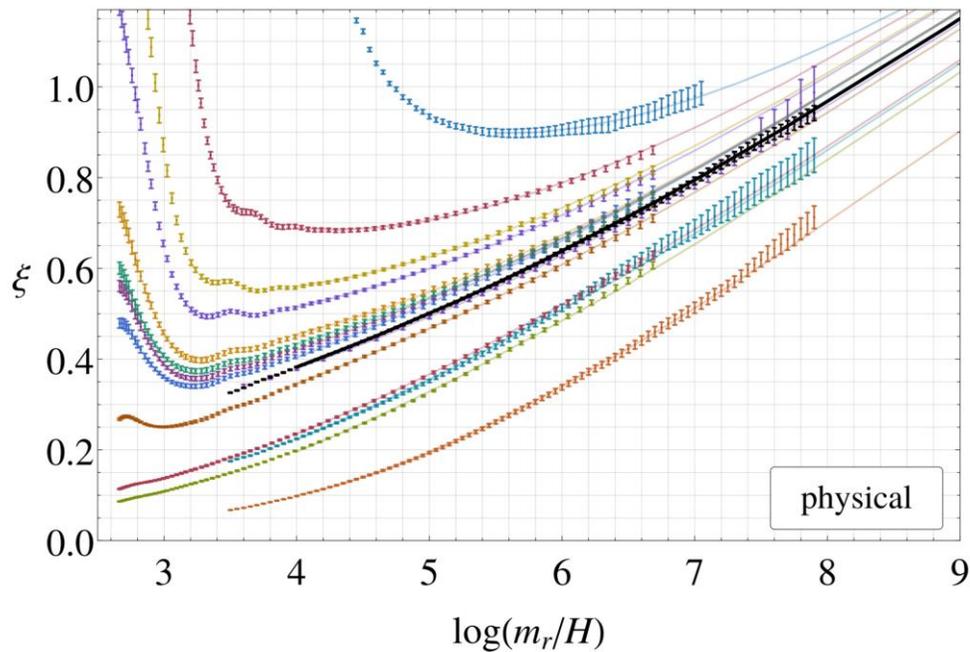
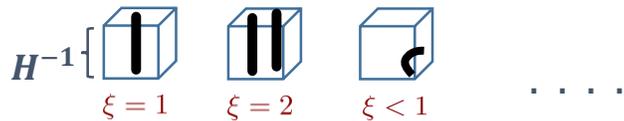




$$\mathcal{L}_\phi = |\partial_\mu \phi|^2 - \frac{m_r^2}{2v^2} \left(|\phi|^2 - \frac{v^2}{2} \right)^2$$



Attractor and Energy Emission



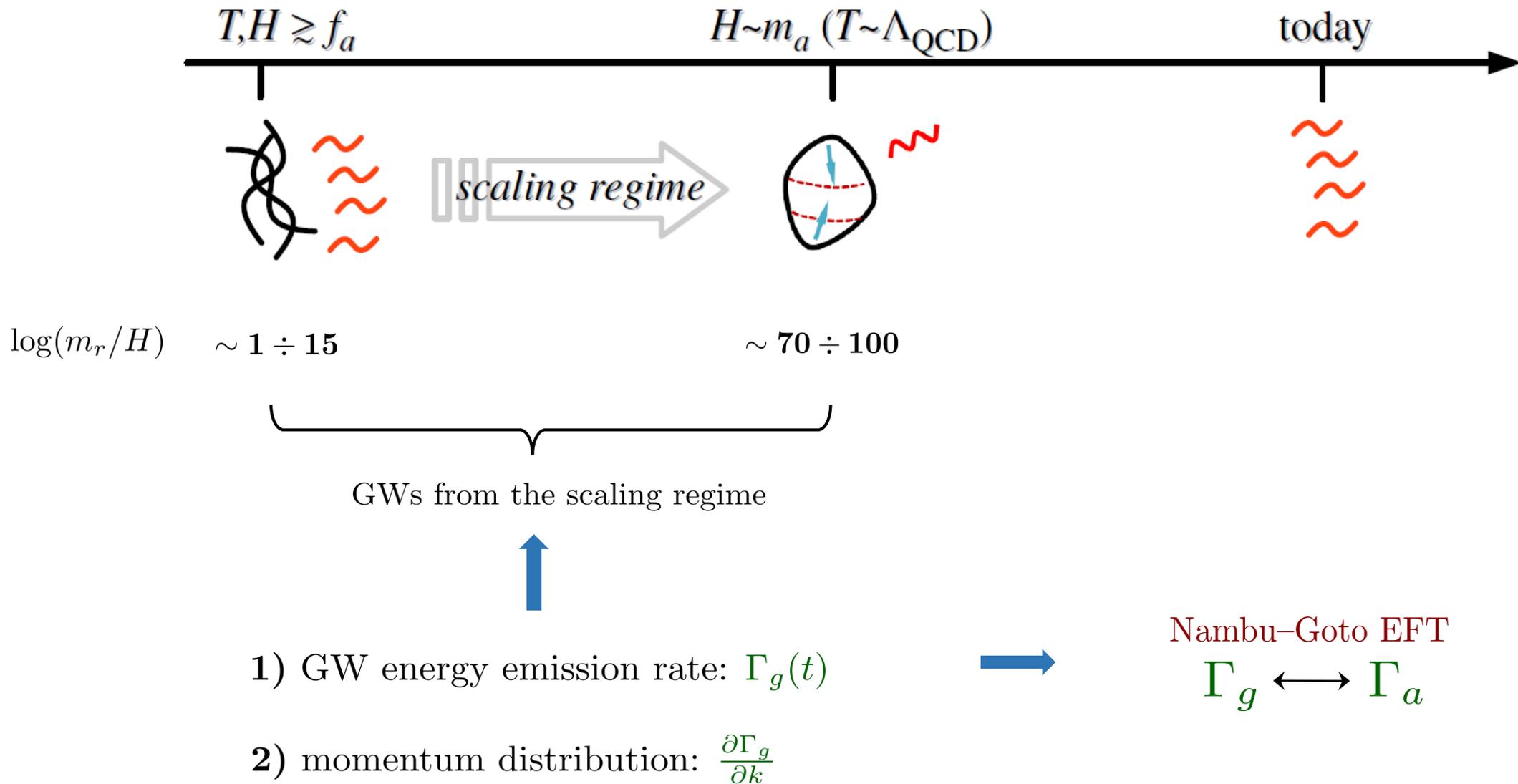
Scaling Violation

$$\xi = c_1 \log + c_0 + \frac{c_{-1}}{\log} + \frac{c_{-2}}{\log^2}$$



$$\Gamma_a = \frac{\xi \mu}{t^3} \propto \frac{f_a^2 \log^2}{t^3}$$

Gravitational Waves



String Effective Theory

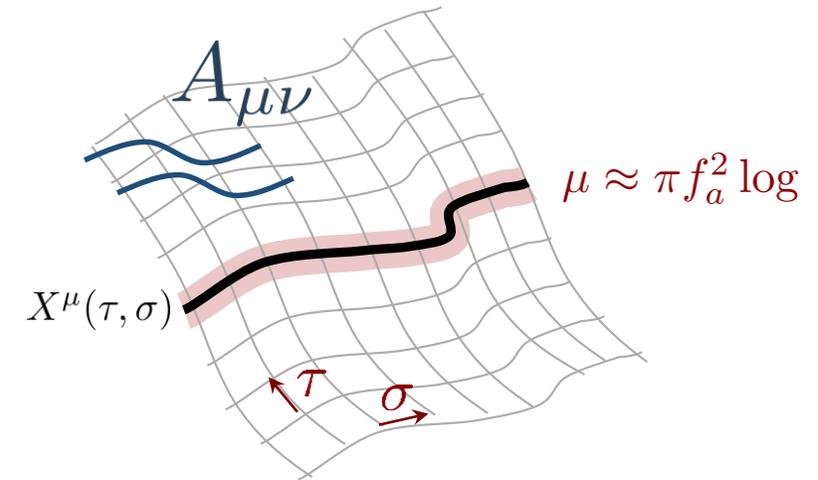
- 1) GW energy emission rate: $\Gamma_g(t)$
- 2) momentum distribution: $\frac{\partial \Gamma_g}{\partial k}$

Degrees of freedom:

- $a \longleftrightarrow A_{\mu\nu}$

$$\partial A \sim F^{\mu\nu\rho} = \epsilon^{\mu\nu\rho\sigma} \partial_\sigma a$$

- $X^\mu(\tau, \sigma)$



$$S_\phi[\phi]$$



$$S_{\text{EFT}}[X, A] =$$

$$-\mu \int d\tau d\sigma \sqrt{-\gamma}$$

Nambu-Goto action

$$\gamma_{ab} = \partial_a X^\mu \partial_b X_\mu$$

$$-\frac{1}{6} \int d^4x (\partial A)^2$$

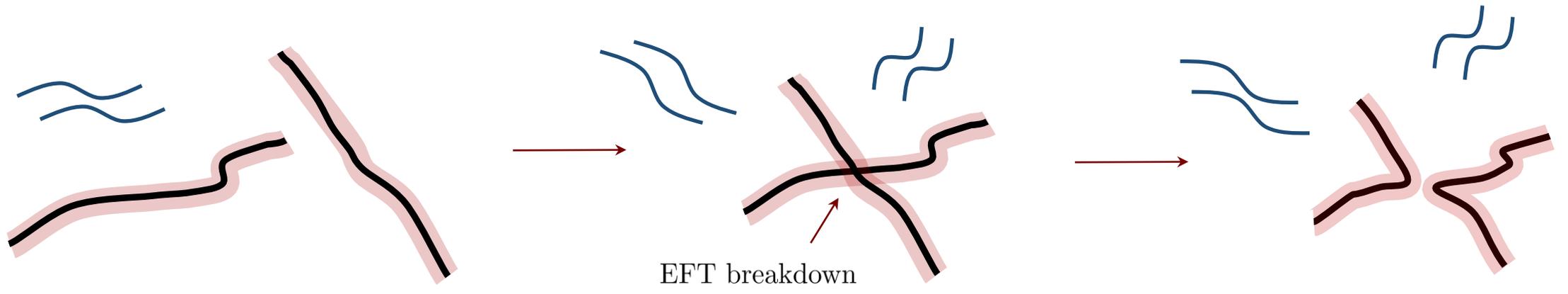
Axion kinetic term

axion-to-string coupling

$$+ 2\pi f_a \int d\tau d\sigma \partial_\tau X^\mu \partial_\sigma X^\nu A_{\mu\nu}$$

Axion-string interaction
(Kalb-Ramond action)

- 1) GW energy emission rate: $\Gamma_g(t)$
- 2) momentum distribution: $\frac{\partial \Gamma_g}{\partial k}$



Gravitational Wave Emission

- 1) GW energy emission rate: $\Gamma_g(t)$
- 2) momentum distribution: $\frac{\partial \Gamma_g}{\partial k}$

EoM: $\square_x A^{\mu\nu} = 2\pi \boxed{f_a} \int d\sigma \dot{X}^{[\mu} X'^{\nu]} \delta^3(\vec{x} - \vec{X})$

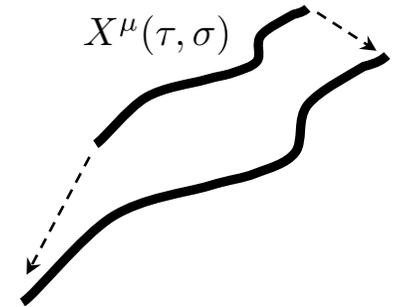
Einstein Eq: $\square_x h^{\mu\nu} = 16\pi G (T_s^{\mu\nu} - \frac{1}{2}\eta^{\mu\nu} T_s^\lambda{}_\lambda)$

$$T_s^{\mu\nu} = \boxed{\mu} \int d\sigma (\dot{X}^\mu \dot{X}^\nu - X'^\mu X'^\nu) \delta^3(\vec{x} - \vec{X})$$

$$\frac{dE_a}{dt} = \underbrace{r_a[X]}_{\uparrow} \boxed{f_a^2}$$

$$\frac{dE_g}{dt} = \underbrace{r_g[X]}_{\uparrow} G \boxed{\mu^2}$$

dimensionless functionals of the shape of the string trajectory X^μ



$$\frac{\Gamma_g}{\Gamma_a} = \underbrace{\frac{r_g[X]}{r_a[X]}}_{\equiv r} \frac{G\mu^2}{f_a^2}$$

$$\equiv r = \text{const} = \mathcal{O}(1)$$

$$\Gamma_g = r \frac{G\mu^2}{f_a^2} \Gamma_a$$

$\frac{\xi_\mu}{t^3}$

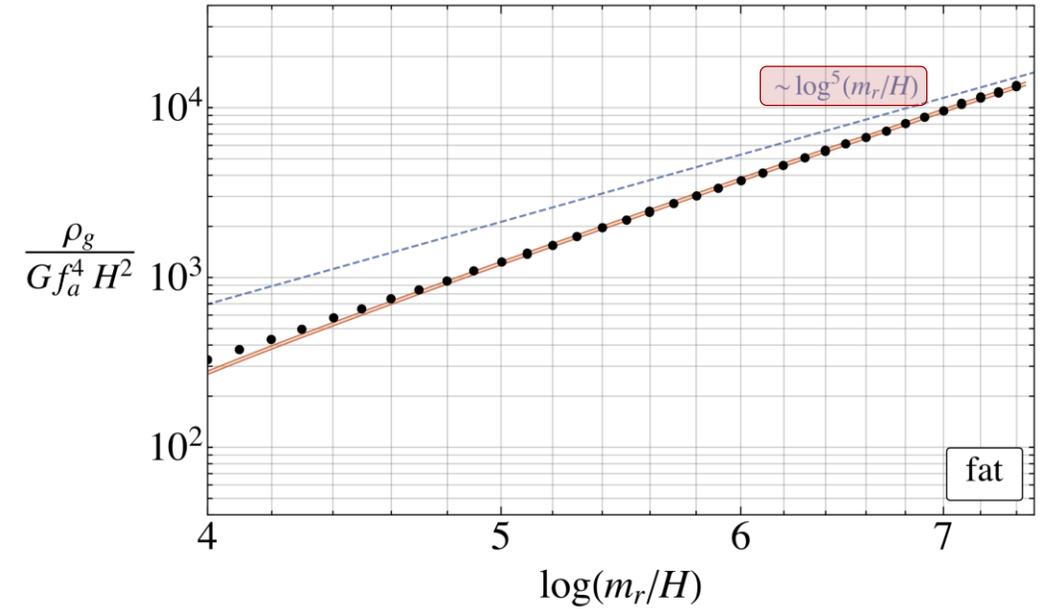
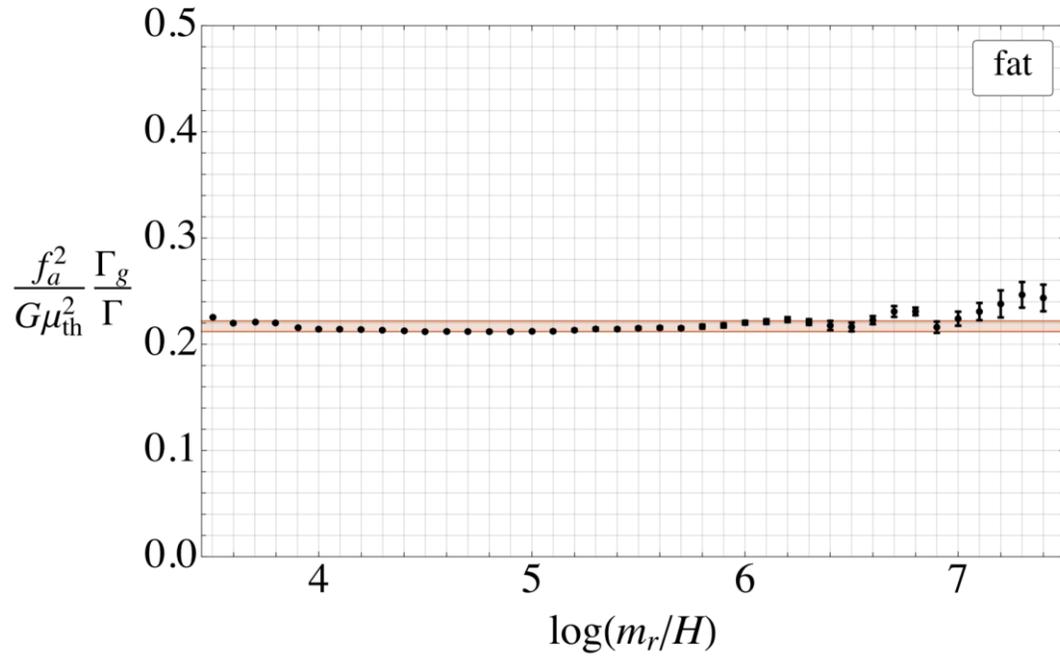
$$\propto \frac{\log^4}{t^3}$$

Comparison with field theory simulations

- 1) GW energy emission rate: $\Gamma_g(t)$
- 2) momentum distribution: $\frac{\partial \Gamma_g}{\partial k}$

$$\Gamma_g = r \frac{G\mu^2}{f_a^2} \Gamma_a$$

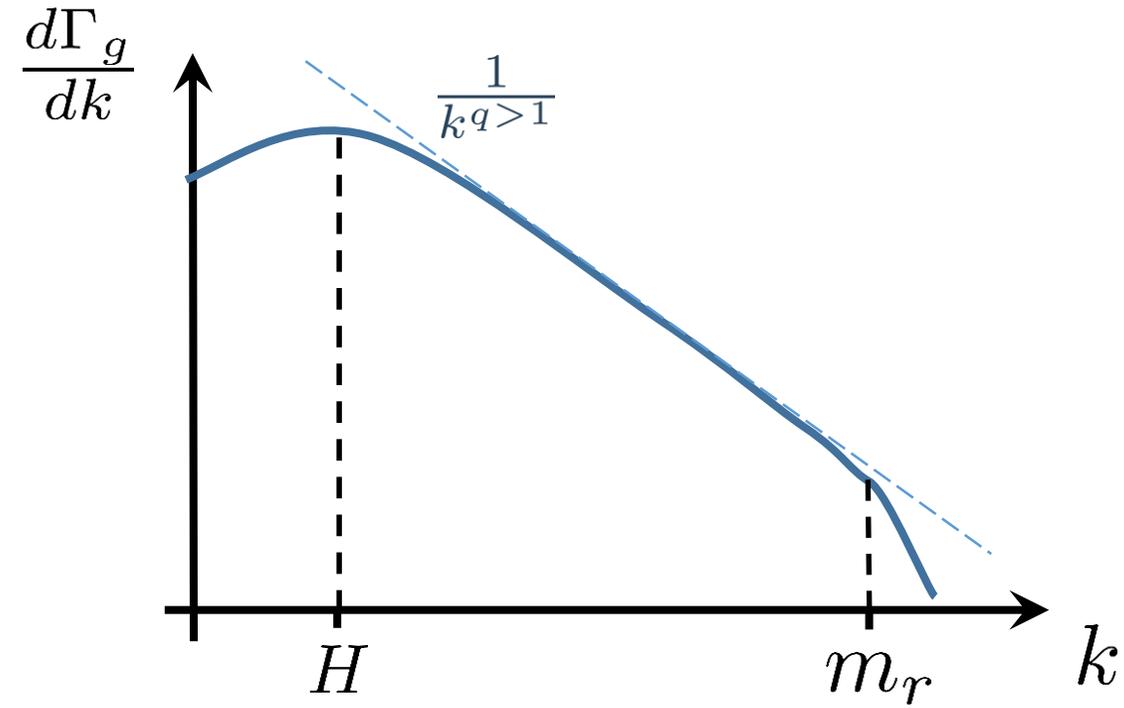
$\swarrow \frac{\xi\mu}{t^3}$



total GW energy: $\rho_g = \int dt' \left(\frac{R'}{R}\right)^4 \Gamma'_g \propto \frac{\log^5}{t^2}$

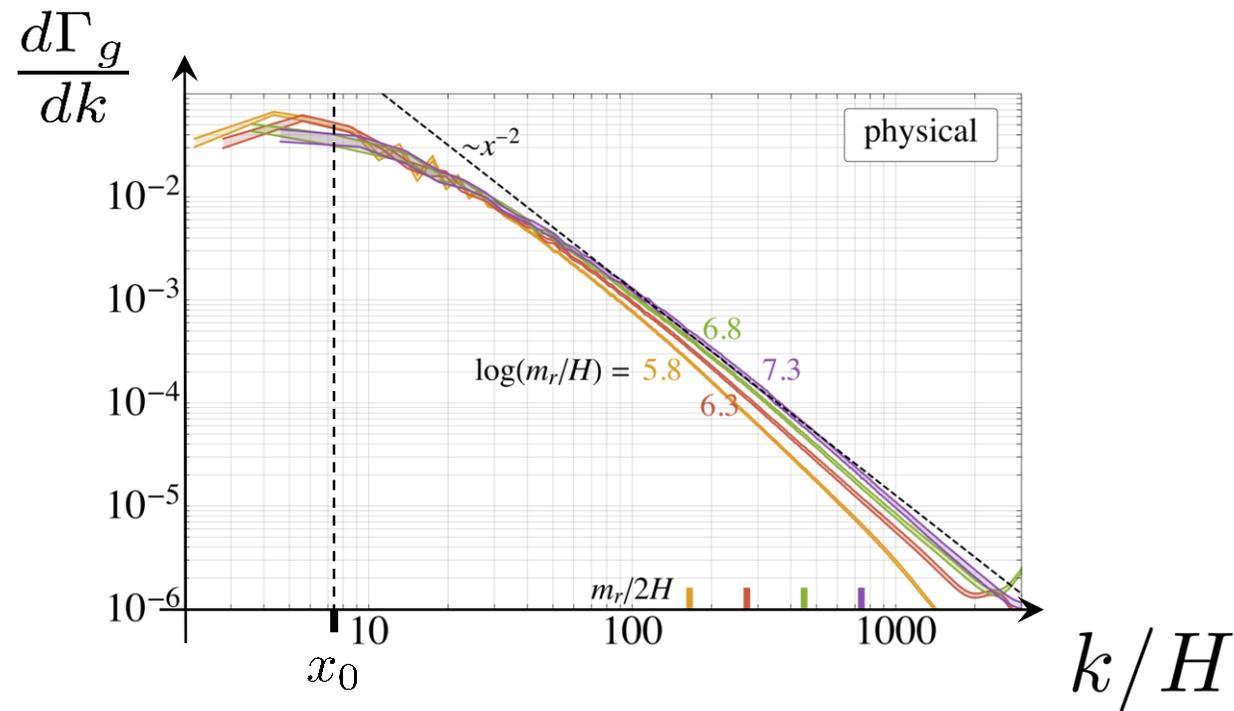
The Gravitational Wave Spectrum

- 1) GW energy emission rate: $\Gamma_g(t)$
- 2) momentum distribution: $\frac{\partial \Gamma_g}{\partial k}$

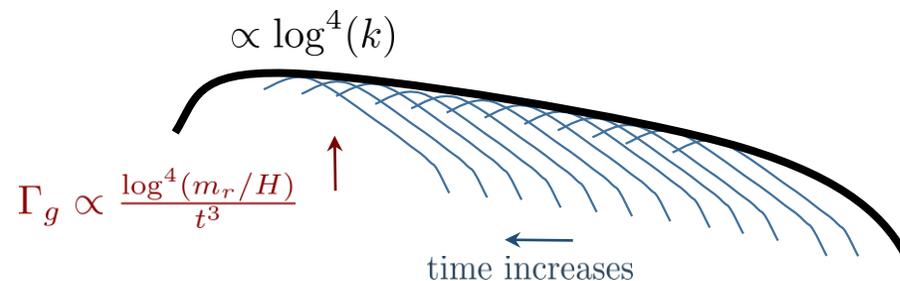


The Gravitational Wave Spectrum

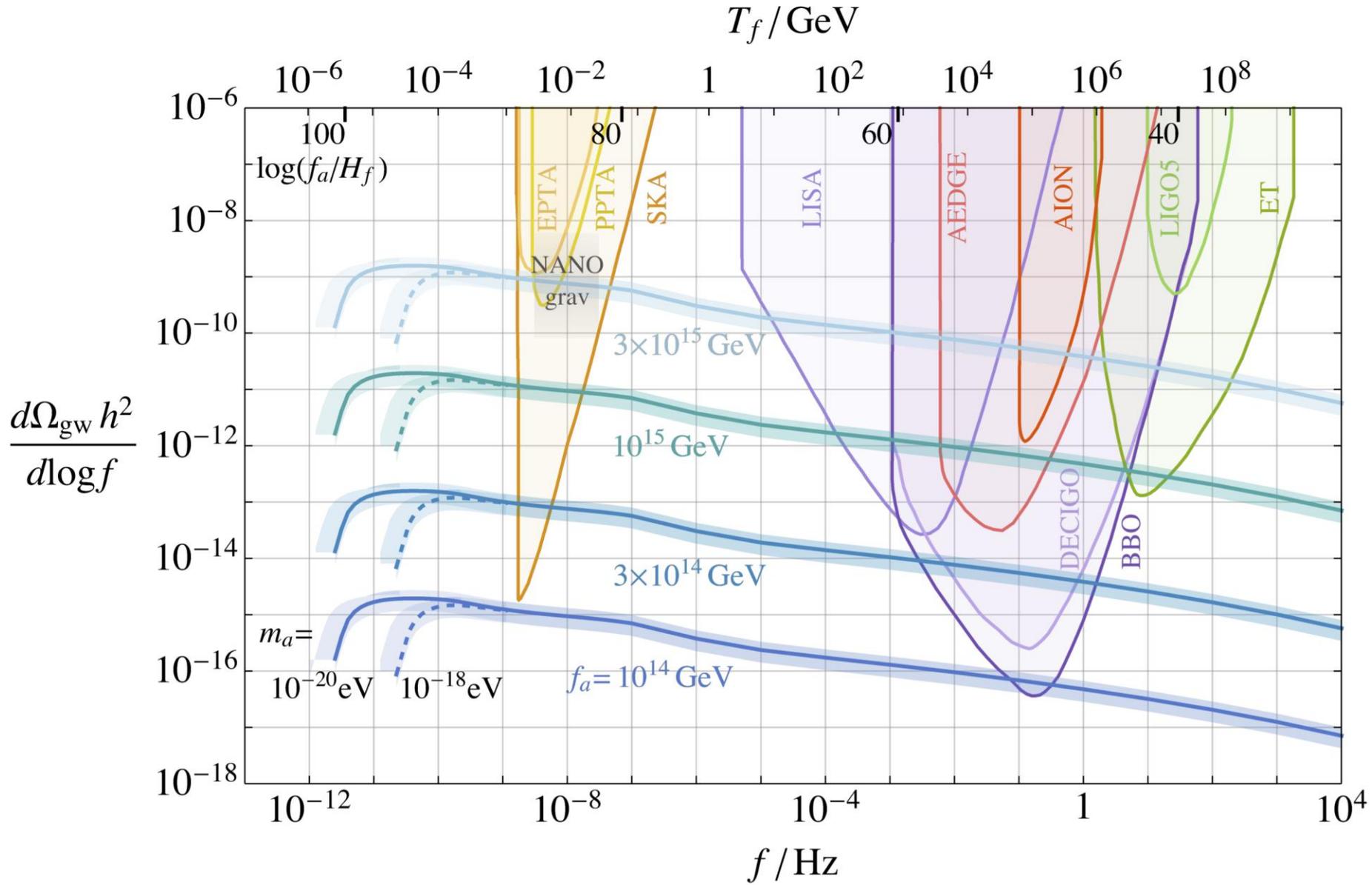
- 1) GW energy emission rate: $\Gamma_g(t)$
- 2) momentum distribution: $\frac{\partial \Gamma_g}{\partial k}$



$$\frac{\partial \rho_g}{\partial \log k} \equiv \int dt' \frac{d\Gamma'_g}{d \log k} \left(\frac{R'}{R}\right)^4 =$$



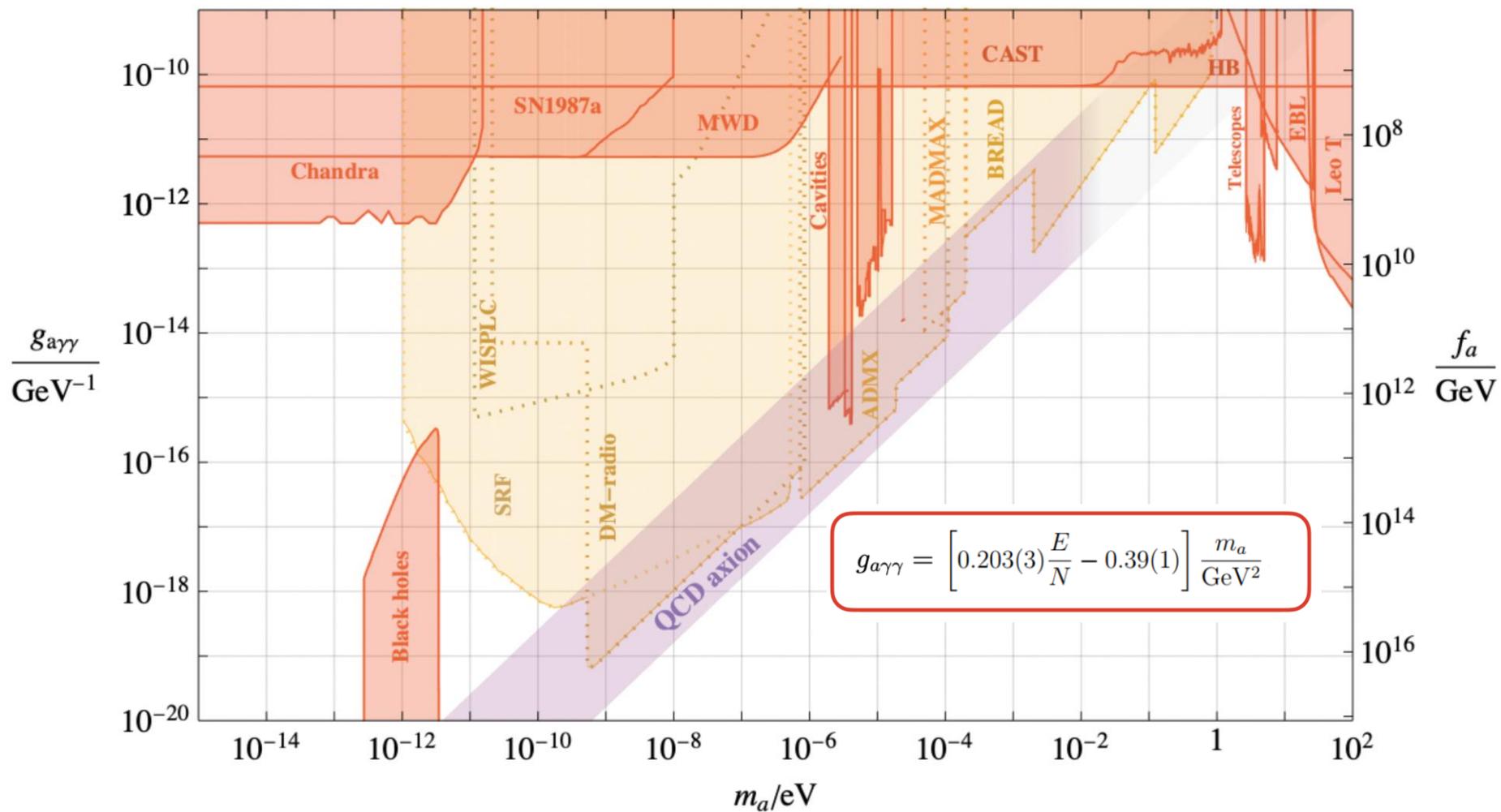
- approximately scale invariant
- \log^4 enhancement



$$\left[\begin{array}{l} f_a \lesssim 10^{15} \text{ GeV} \\ m_a \lesssim 10^{-18} \text{ eV} \end{array} \right.$$

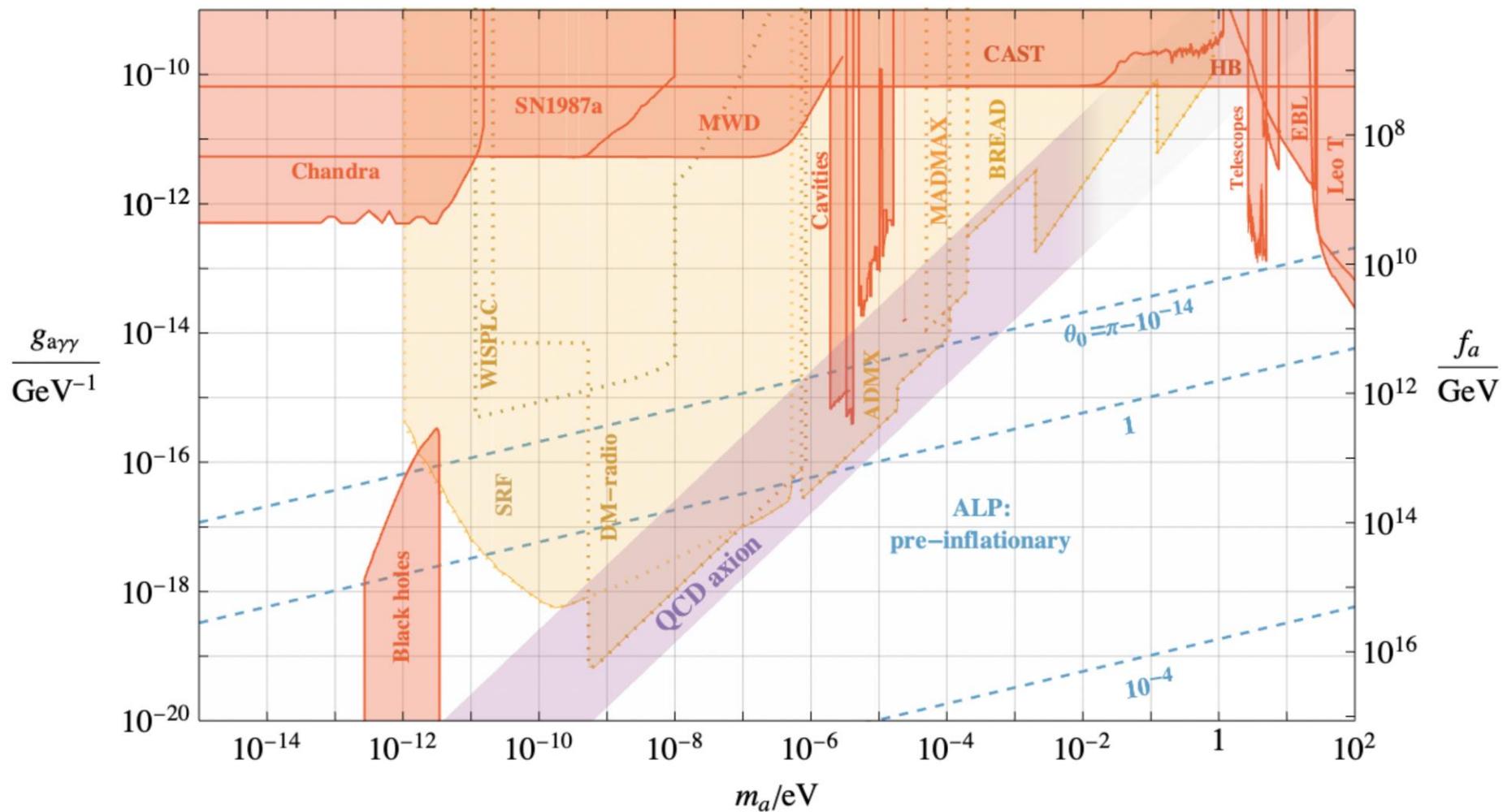
$$\frac{d\Omega_{\text{gw}} h^2}{d\log f} \simeq 10^{-15} \left(\frac{r}{0.26} \right) \left(\frac{f_a}{10^{14} \text{ GeV}} \right)^4 \left(\frac{10}{g_f} \right)^{\frac{1}{3}} \left\{ 1 + 0.12 \log \left[\left(\frac{m_r}{10^{14} \text{ GeV}} \right) \left(\frac{10^{-8} \text{ Hz}}{f} \right)^2 \right] \right\}^4$$

Targets for haloscopes



$$g_{a\gamma\gamma} \simeq \frac{\alpha_{em}}{2\pi f_a}$$

Targets for haloscopes

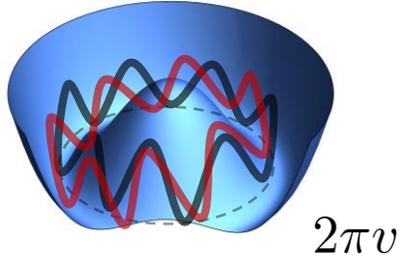


$$\frac{\Omega_a^{\text{mis}}}{\Omega_{\text{DM}}} \simeq 2.2 \cdot 10^{-3} \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^2 \left(\frac{m_a}{10^{-6} \text{ eV}} \right)^{1/2} h(\theta_0) \theta_0^2$$

$$g_{a\gamma\gamma} \simeq \frac{\alpha_{em}}{2\pi f_a}$$

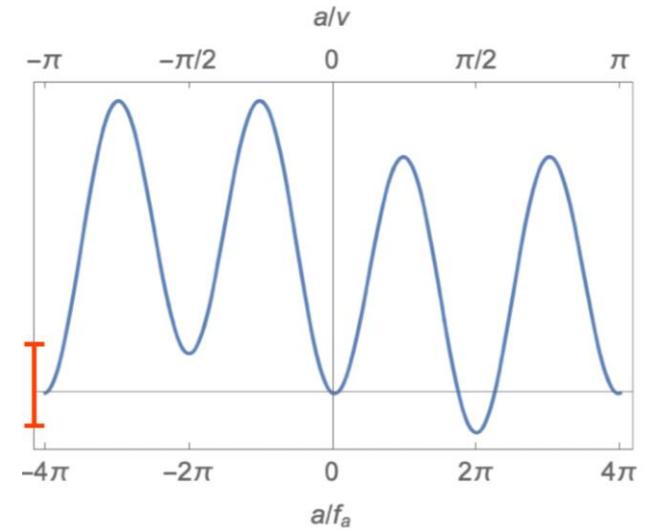
Domain wall number $N > 1$

$$\phi = \frac{r+v}{\sqrt{2}} e^{i\frac{a}{v}}$$

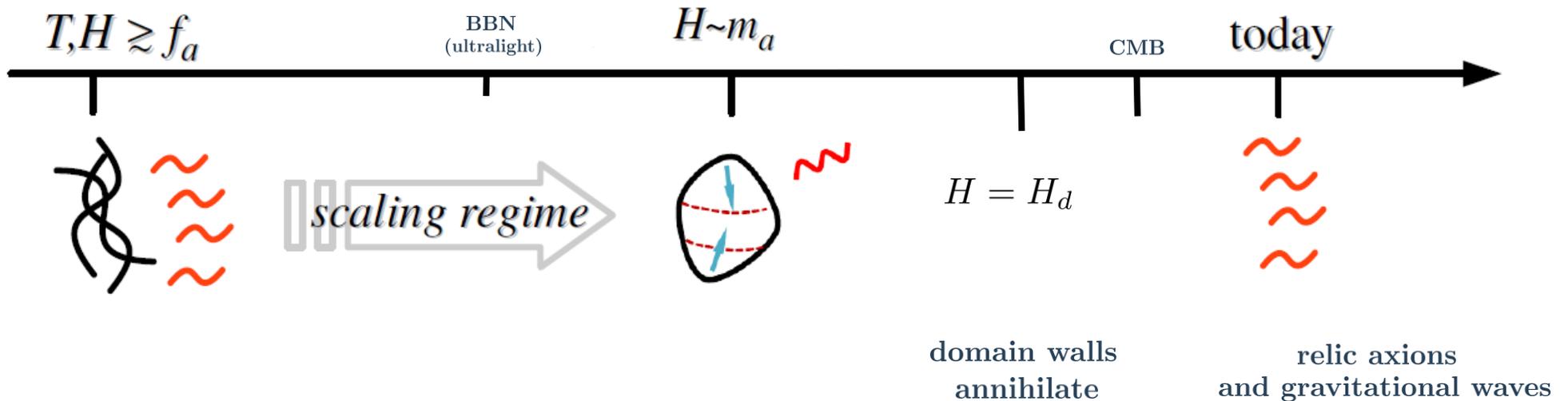


$$v = N f_a$$

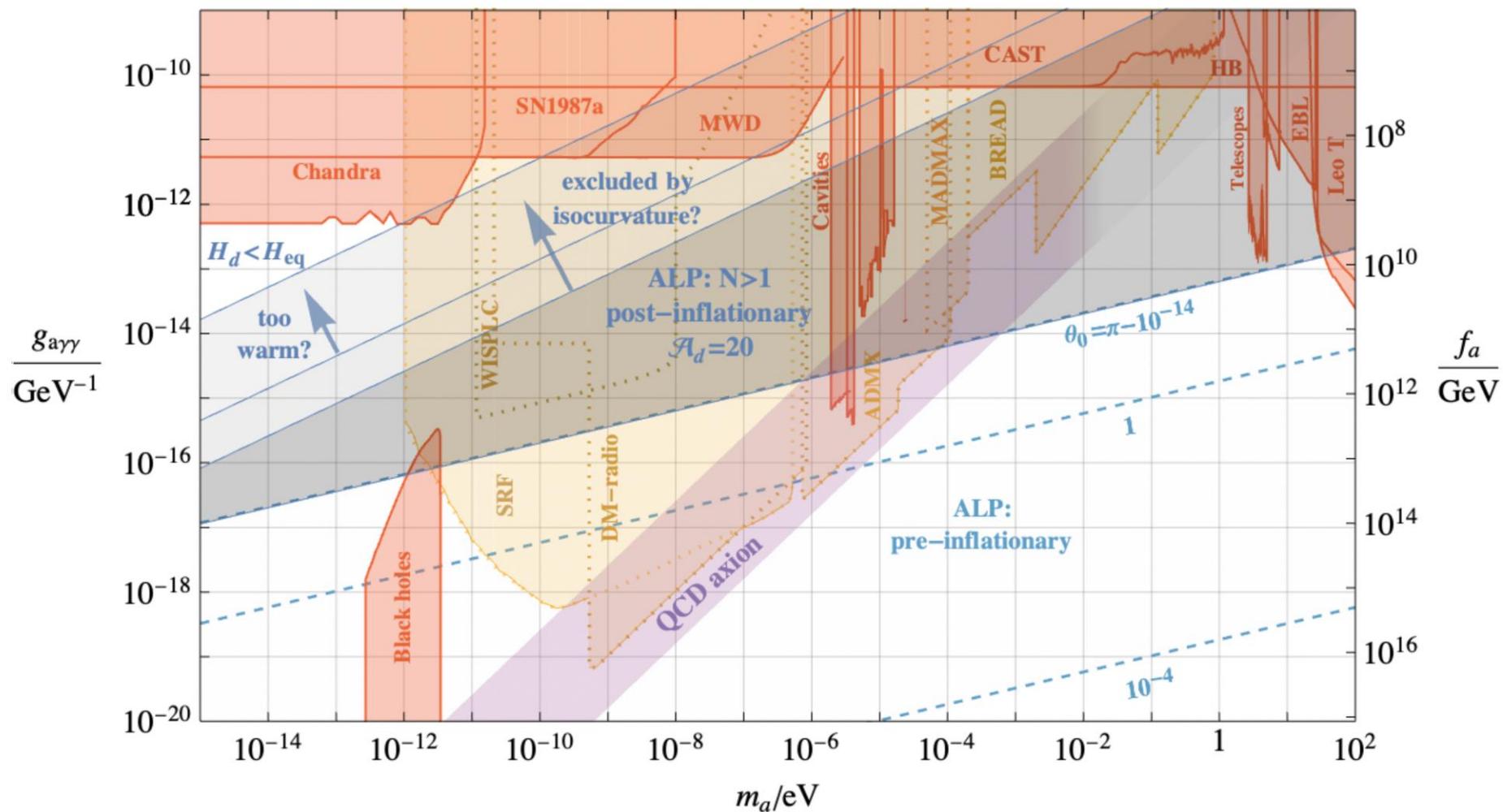
$$V(a) = \frac{m_a^2 v^2}{N^2} \left(1 - \cos \left(N \frac{a}{v} \right) \right) + \delta V_{PQV}$$



$N = 2$

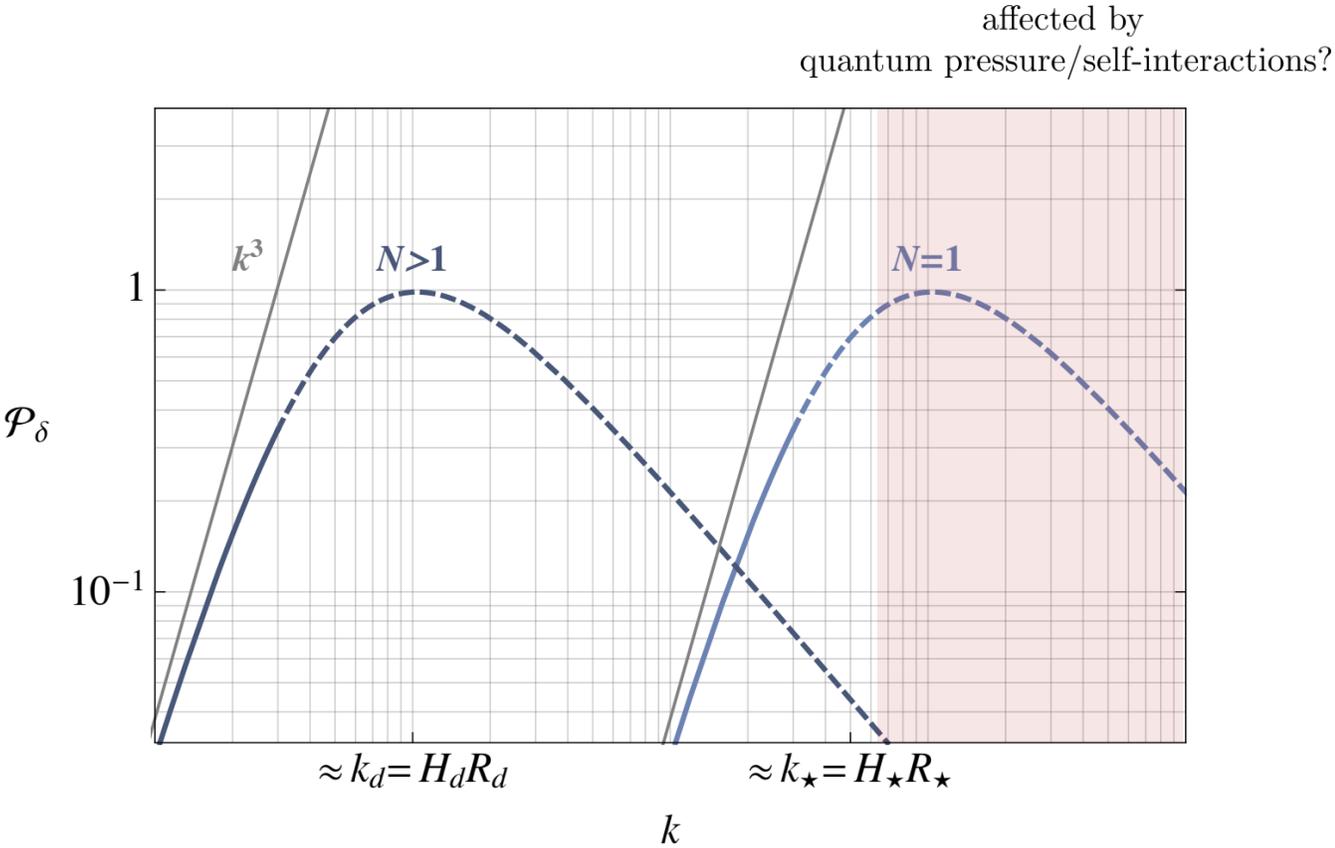
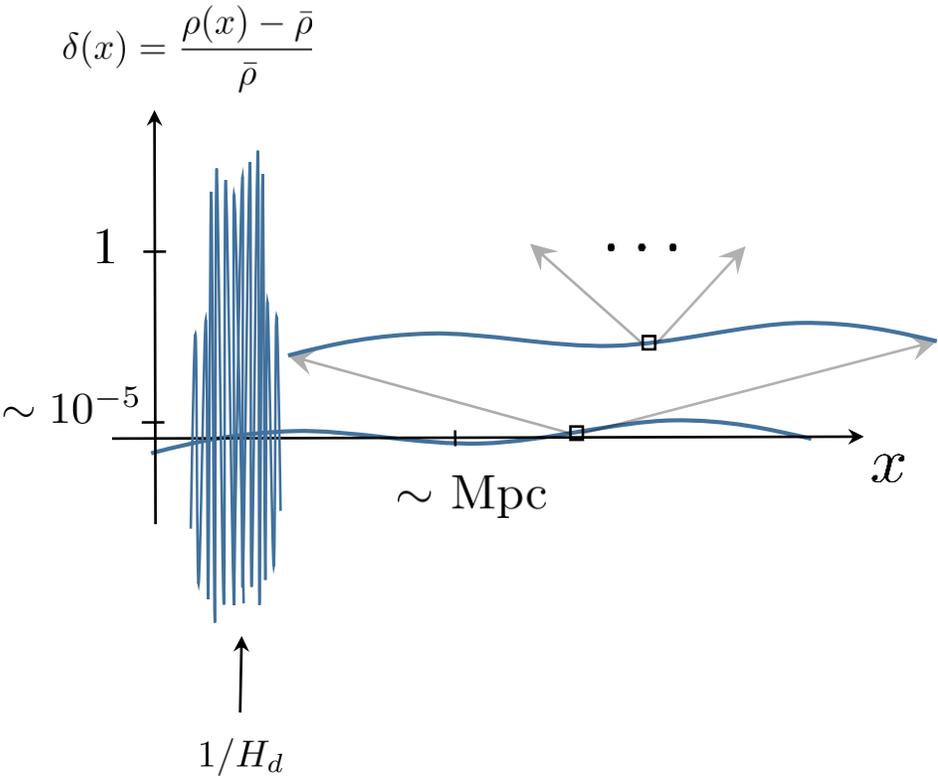


Targets for haloscopes



$$\frac{\Omega_a}{\Omega_{\text{DM}}} \simeq 2 \left(\frac{\mathcal{A}_d}{20} \right) \left(\frac{m_a}{H_d} \right)^{1/2} \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^2 \left(\frac{m_a}{10^{-6} \text{ eV}} \right)^{1/2}$$

Overdensities at small scales



work in progress

standard lore: collapse at around matter-radiation equality into miniclusters

Conclusions

- **The scaling regime produces an approximately scale invariant GW spectrum**
 - $\Gamma_g \propto \log^4$ (from the increase in $\mu \propto \log$ and $\xi \propto \log$)
 - enhances the spectrum at low frequencies
- **The spectrum is visible by multiple experiments for $f_a > 10^{14}$ GeV**
 - best prospects in PTAs and LISA
- **ALPs with $N > 1$ in the post-inflationary scenario have a potentially large coupling to photons**

Much more to do

- Local strings?
- QCD axion dark matter substructure
- Similar dynamics in other theories of light dark matter (e.g. dark photons)

Thank You