

DW-genesis: inducing the baryon number with domain walls

Miguel Vanvlasselaer miguel.vanvlasselaer@vub.be

VUB, IIHE and COST Action

February 26, 2025





Miguel Vanvlasselaer (VUB, IIHE and COST Action)

DW-genesis: inducing the baryon number with domain walls

Mexican hat potential



• Mexican hat potential for a complex field ϕ with U(1) global symmetry

$$\mathsf{high} \ \mathsf{T} \gg f_a: \quad V(\phi) \approx |\phi|^2 T^2 \qquad \qquad \mathsf{low} \ \mathsf{T} \ll f_a: \quad V(\phi) = \lambda \Big(|\phi|^2 - f_a^2/2 \Big)$$

1

 $\sqrt{2}$

Mexican hat potential



• Mexican hat potential for a complex field ϕ with U(1) global symmetry

$$\text{high } \mathsf{T} \gg f_a: \quad V(\phi) \approx |\phi|^2 T^2 \qquad \qquad \text{low } \mathsf{T} \ll f_a: \quad V(\phi) = \lambda \bigg(|\phi|^2 - f_a^2/2 \bigg)^2$$

• Axions are characterized by a spontaneously broken U(1) symmetry with periodicity f_a

$$\phi = \rho e^{ia/f_a}$$

\$ 2

Mexican hat potential



 \bullet Mexican hat potential for a complex field ϕ with U(1) global symmetry

$$\text{high }\mathsf{T}\gg f_a:\quad V(\phi)\approx |\phi|^2T^2\qquad\qquad \text{low }\mathsf{T}\ll f_a:\quad V(\phi)=\lambda \Big(\,|\phi|^2-f_a^2/2\,\Big)^{-1}$$

• Axions are characterized by a spontaneously broken U(1) symmetry with periodicity f_a

$$\phi = \rho e^{ia/f_a}$$

• Formation of cosmic strings at $T \sim f_a$ and spreads of value of a around strings

\$ 2

Universe high-T after inflation:cooling of primordial soup







Universe high-T after inflation:cooling of primordial soup



• $T \sim \Lambda \sim \sqrt{m_a f_a}$, dark sector confines:

$$\frac{a}{f_a}G_d\tilde{G}_d \quad \Rightarrow \quad V(a) = \Lambda^4 \left(1 - \cos\left(\frac{aN_{\rm DW}}{f_a}\right)\right)$$





Universe high-T after inflation:cooling of primordial soup



• $T \sim \Lambda \sim \sqrt{m_a f_a}$, dark sector confines:

$$\frac{a}{f_a}G_d\tilde{G}_d \quad \Rightarrow \quad V(a) = \Lambda^4 \left(1 - \cos\left(\frac{aN_{\rm DW}}{f_a}\right)\right)$$

• Breaking of the $U(1) \rightarrow Z_{N_{\rm DW}}$





Universe high-T after inflation:cooling of primordial soup



• $T \sim \Lambda \sim \sqrt{m_a f_a}$, dark sector confines:

$$\frac{a}{f_a}G_d\tilde{G}_d \quad \Rightarrow \quad V(a) = \Lambda^4 \left(1 - \cos\left(\frac{aN_{\rm DW}}{f_a}\right)\right)$$

- Breaking of the $U(1) \rightarrow Z_{N_{\rm DW}}$
- Formation of DW for $T < \Lambda$





Formation of domain walls

Blasi, Mariotti, Rase, MV



 $\eta = 0$



Miguel Vanvlasselaer (VUB, IIHE and COST Action)

DW-genesis: inducing the baryon number with domain walls





DW domination and collapse

• Domain wall precludes discrete symmetries ? Zel'dovich, Kobzarev, Okun (1974)

$$ho_{
m DW} \sim \sigma H \sim m_a f_a^2 H \qquad \qquad rac{\Omega_{
m DW}}{\Omega_{
m rad}} \propto rac{m_a f_a^2}{T^2 M_{
m pl}} \qquad \Rightarrow {\sf DW} \ {\sf domination} \ {\sf at} \ T_{
m dom} pprox \sqrt{\sigma/M_{
m pl}}$$

DW domination and collapse

• Domain wall precludes discrete symmetries ? Zel'dovich, Kobzarev, Okun (1974)

$$ho_{
m DW} \sim \sigma H \sim m_a f_a^2 H \qquad \qquad rac{\Omega_{
m DW}}{\Omega_{
m rad}} \propto rac{m_a f_a^2}{T^2 M_{
m pl}} \qquad \Rightarrow {\sf DW} \ {\sf domination} \ {\sf at} \ T_{
m dom} pprox \sqrt{\sigma/M_{
m pl}}$$

 \bullet Solution: DWs have to decay at $T_{\rm ann}>T_{\rm dom}$

DW domination and collapse

• Domain wall precludes discrete symmetries ? Zel'dovich, Kobzarev, Okun (1974)

$$\rho_{\rm DW} \sim \sigma H \sim m_a f_a^2 H \qquad \qquad \frac{\Omega_{\rm DW}}{\Omega_{\rm rad}} \propto \frac{m_a f_a^2}{T^2 M_{\rm pl}} \qquad \Rightarrow {\sf DW} \text{ domination at } T_{\rm dom} \approx \sqrt{\sigma/M_{\rm pl}}$$

 \bullet Solution: DWs have to decay at $T_{\rm ann} > T_{\rm dom}$

• Introduce a bias:



• Example: gravitational breaking

$$V_{\rm bias} \sim \frac{\phi^5}{M_{\rm pl}}$$

$$T_{\rm ann}^2 \approx \frac{\Delta V M_{\rm pl}}{24 \, m_a f_a^2}$$

• Amount of breaking:

$$\Delta \theta \equiv \frac{\Delta V}{V_0}$$

4

Baryogenesis with Domain Walls: DW-genesis

Baryogenesis with Domain Walls: DW-genesis [2411.13494]: Alberto Mariotti, Xander Nagels, Aaron Rase, MV

Problem of baryogenesis and Sakharov conditions

• Inflation: erase the initial conditions and produce $n_B = n_{\bar{B}}$ (likely...)



• $\frac{n_B-n_{\bar{B}}}{n_{\gamma}} \rightarrow 0(10^{-20}...)$ (naive expectation): empty universe

Problem of baryogenesis and Sakharov conditions

• Inflation: erase the initial conditions and produce $n_B = n_{\bar{B}}$ (likely...)



- $\frac{n_B n_{\bar{B}}}{n_{\gamma}} \rightarrow 0(10^{-20}...)$ (naive expectation): empty universe
- $\frac{n_B n_{\bar{B}}}{n_{\gamma}} \rightarrow 10^{-10}$ (observation): need a process that turns one out of 10^{10} anti-baryon \rightarrow baryon

Sakharov conditions

- Out-of-equilibrium situation; if CPT theorem
- **CP-violation**: $\Gamma \neq \overline{\Gamma}$
- B-number violation

Problem of baryogenesis and Sakharov conditions

• Inflation: erase the initial conditions and produce $n_B = n_{\bar{B}}$ (likely...)



- $\frac{n_B n_{\bar{B}}}{n_{\gamma}} \rightarrow 0(10^{-20}...)$ (naive expectation): empty universe
- $\frac{n_B n_{\bar{B}}}{n_{\gamma}} \rightarrow 10^{-10}$ (observation): need a process that turns one out of 10^{10} anti-baryon \rightarrow baryon

Sakharov conditions

- Out-of-equilibrium situation; if CPT theorem
- **CP-violation**: $\Gamma \neq \overline{\Gamma}$
- B-number violation
- Example: electroweak baryogenesis, leptogenesis, GUT baryogenesis, Affleck-Dine baryogenesis

Observation prospects of GW



GW from baryogenesis

Baryogenesis with gravitational waves signal

GW from baryogenesis

Baryogenesis with gravitational waves signal







Eung, Dutka, Jung, Nagels, MV, 23':

[2411.13494]: Mariotti, Nagels, Rase, MV

• See-saw model: breaks L number

$$\mathcal{L}_{\not L} = y_N (\tilde{H} \bar{L}) N_R + rac{1}{2} M_N \bar{N}_R^c N_R + {\sf h.c.}$$

• See-saw model: breaks L number

$$\mathcal{L}_{\not\!L} = y_N(\tilde{H}\bar{L})N_R + \frac{1}{2}M_N\bar{N}_R^cN_R + \text{h.c.}$$

• Coupling between an axion and the lepton number

$$\mathcal{L}_{a-L} = rac{c_L \partial_\mu a}{f_a} j_L^\mu, \qquad j_L^\mu \equiv \bar{L} \gamma^\mu L \qquad \mathcal{L} = \mathcal{L}_{a-j} + \mathcal{L}_L$$

• See-saw model: breaks L number

$$\mathcal{L}_{\not\!L} = y_N(\tilde{H}\bar{L})N_R + \frac{1}{2}M_N\bar{N}_R^cN_R + \text{h.c.}$$

• Coupling between an axion and the lepton number

$$\mathcal{L}_{a-L} = rac{c_L \partial_\mu a}{f_a} j_L^\mu, \qquad j_L^\mu \equiv \bar{L} \gamma^\mu L \qquad \mathcal{L} = \mathcal{L}_{a-j} + \mathcal{L}_L$$

• Moving non-trivial profile for the axion breaks both CP and CPT

$$\mathcal{L}_{a-L} = \dot{\theta} j_L^0 \simeq \mu j_L^0 \,, \qquad \qquad \dot{\theta} \equiv \frac{\dot{a}_{\rm DW}(t,z)}{f_a} = \frac{2m_a \gamma_w v_w}{\cosh\left[m_a \gamma_w v_w \left(t - t_{\rm passage}\right)\right]}$$

• See-saw model: breaks L number

$$\mathcal{L}_{\not\!L} = y_N(\tilde{H}\bar{L})N_R + \frac{1}{2}M_N\bar{N}_R^cN_R + \text{h.c.}$$

• Coupling between an axion and the lepton number

$$\mathcal{L}_{a-L} = rac{c_L \partial_\mu a}{f_a} j_L^\mu, \qquad j_L^\mu \equiv \bar{L} \gamma^\mu L \qquad \mathcal{L} = \mathcal{L}_{a-j} + \mathcal{L}_L$$

• Moving non-trivial profile for the axion breaks both CP and CPT

$$\mathcal{L}_{a-L} = \dot{\theta} j_L^0 \simeq \mu j_L^0 \,, \qquad \qquad \dot{\theta} \equiv \frac{\dot{a}_{\rm DW}(t,z)}{f_a} = \frac{2m_a \gamma_w v_w}{\cosh\left[m_a \gamma_w v_w\left(t - t_{\rm passage}\right)\right]}$$

Boltzmann equations deduced

$$\frac{dY_{\Delta L}}{dt} = -\left(\frac{\gamma_D}{n_L^{\rm eq}}(Y_{\Delta L} + Y_{\Delta L}^{\rm eq}(t)) + 2\frac{\gamma_{2\to2}}{n_L^{\rm eq}}\left(Y_{\Delta L} + Y_{\Delta L}^{\rm eq}(t)\right)\right), \quad Y_{\Delta L}^{\rm eq}(t) \equiv \frac{n_L^{\rm eq}}{s}\frac{2c_L\dot{a}}{f_a T}$$

Miguel Vanvlasselaer (VUB, IIHE and COST Action)

Intuition and computation

• Rotate away to the coupling to the DW

$$\begin{split} L &\to e^{ic_L a/f_a} L \qquad \Rightarrow \qquad i\bar{L}\partial\!\!\!/ L \to i\bar{L}\partial\!\!\!/ L - c_L \frac{\partial_\mu a}{f_a} j_L^\mu \,, \qquad y_N(\tilde{H}\bar{L})N_R \to e^{-ic_L a/f_a} y_N(\tilde{H}\bar{L})N_R \,. \\ &\Rightarrow \qquad \mathcal{L}_{\not\!\!L} = e^{-ic_L a/f_a} y_N(\tilde{H}\bar{L})N_R + \frac{1}{2} M_N \bar{N}_R^c N_R + \text{h.c.} \,. \end{split}$$

Intuition and computation

 $\times []$

• Rotate away to the coupling to the DW

$$L \to e^{ic_L a/f_a} L \qquad \Rightarrow \qquad i\bar{L}\partial \!\!\!/ L \to i\bar{L}\partial \!\!/ L - c_L \frac{\partial_\mu a}{f_a} j_L^\mu, \qquad y_N(\tilde{H}\bar{L})N_R \to e^{-ic_L a/f_a} y_N(\tilde{H}\bar{L})N_R.$$

$$\Rightarrow \qquad \mathcal{L}_{\not\!L} = e^{-ic_L a/f_a} y_N(\tilde{H}\bar{L})N_R + \frac{1}{2}M_N\bar{N}_R^c N_R + \mathrm{h.c.} \, .$$

• New Feynman rules from the vertex : $\delta(\sum_i E_i) \rightarrow \delta(\sum_i E_i + c_L \dot{a}/f_a)$

$$\dot{n}_L + 3Hn_L = \int \frac{d^3p_L}{(2\pi)^3} \left(\mathcal{C}^{\text{decay}}[f_L] + \mathcal{C}^{\text{scatt}}[f_L] \right) \equiv I^{\text{decay}} + I^{\text{scatt}} ,$$
$$I^{\text{decay}} = \int d\Pi_L d\Pi_N d\Pi_H (2\pi)^4 \delta^3 (p_N - p_L - p_H) \delta(E_N - E_L - E_H - c_L \dot{a}/f_a)$$
$$\mathcal{M}_{N \to LH} |^2 f_N - |\mathcal{M}_{LH \to N}|^2 f_L f_H]$$

Intuition and computation

• Rotate away to the coupling to the DW

$$L \to e^{ic_L a/f_a} L \qquad \Rightarrow \qquad i\bar{L}\partial \!\!\!/ L \to i\bar{L}\partial \!\!/ L - c_L \frac{\partial_\mu a}{f_a} j_L^\mu, \qquad y_N(\tilde{H}\bar{L})N_R \to e^{-ic_L a/f_a} y_N(\tilde{H}\bar{L})N_R.$$

$$\Rightarrow \qquad \mathcal{L}_{\not\!L} = e^{-ic_L a/f_a} y_N(\tilde{H}\bar{L})N_R + \frac{1}{2}M_N\bar{N}_R^c N_R + \mathrm{h.c.} \, .$$

• New Feynman rules from the vertex : $\delta(\sum_i E_i) \rightarrow \delta(\sum_i E_i + c_L \dot{a}/f_a)$

$$\dot{n}_L + 3Hn_L = \int \frac{d^3p_L}{(2\pi)^3} \left(\mathcal{C}^{\text{decay}}[f_L] + \mathcal{C}^{\text{scatt}}[f_L] \right) \equiv I^{\text{decay}} + I^{\text{scatt}} ,$$

$$I^{\text{decay}} = \int d\Pi_L d\Pi_N d\Pi_H (2\pi)^4 \delta^3 (p_N - p_L - p_H) \delta(E_N - E_L - E_H - c_L \dot{a}/f_a)$$

 $\times \left[|\mathcal{M}_{N \to LH}|^2 f_N - |\mathcal{M}_{LH \to N}|^2 f_L f_H \right]$ • Shift in the distributions: like a chemical potential = $c_L \dot{a}/f_a$:

$$f_L f_H \approx e^{-\frac{E_H + E_L}{T}} = e^{-\frac{E_N - c_L \dot{a}/f_a}{T}}$$

[1504.07917]:Daido, Kitajima, Takahashi, [2411.13494]: Mariotti, Nagels, Rase, MV



Production and wash-outs coming from $HL \rightarrow N$, $LH \rightarrow H\bar{L}$ (usually called wash-outs)!

Domain wall leptogenesis



Domain wall leptogenesis



Observed abundance $Y_{\Delta L} \sim 10^{-10}$ requires $T_{\rm ann} \gtrsim 10^9$ GeV.

Parameter space for DW-genesis I

 $v_m = 0.4, \quad c\alpha = 1, \quad \Delta\theta = 10^{-6}$

• Thermalisation condition inside the DW:

 $L_{\rm thermalisation} \ll L_{\rm DW}$ $\mathcal{O}(10)\alpha_w^2 \gtrsim \frac{\gamma_w v_w m_a}{T_{\text{ann}}} \qquad T_{\text{ann}} \gg 10^2 m_a \gamma_w v_w$ $\mathcal{O}(10)\alpha_s^2 \gtrsim \frac{\gamma_w v_w m_a}{T_{\text{ann}}} \qquad T_{\text{ann}} \gg 10 m_a \gamma_w v_w$ • Dilution from the axion decay: $\rho_a = H(T_{ann})\sigma$: $Y_{\Delta L} = Y_{\Delta L}^0 \times D$, with

$$D = \mathsf{Min}\left[1, 0.57 \frac{g_*(T_{\mathsf{ann}})}{g_*(T_{\mathsf{alp dec}})^{1/4}} \frac{\sqrt{M_{\mathsf{Pl}}\Gamma} \ T_{\mathsf{ann}}^3}{\Delta V}\right]$$



 10^{16}

Parameter space for DW-genesis II



Suppression
$$\approx \mathcal{O}(1-10) \times \Delta \theta$$

[2411.13494]: Mariotti, Nagels, Rase, MV



• Gravitational waves from a mass distribution

$$h_{ij}(r,t) = \frac{2G}{r} \ddot{I}_{ij}(t-r), \qquad I_{ij} \equiv \int d^3 r \rho_{dw} \left(r_i r_j - \frac{r^2 \delta_{ij}}{3} \right)$$

• Gravitational waves from a mass distribution

$$h_{ij}(r,t) = \frac{2G}{r}\ddot{I}_{ij}(t-r), \qquad I_{ij} \equiv \int d^3r\rho_{dw} \left(r_i r_j - \frac{r^2 \delta_{ij}}{3}\right)$$

• In the case of DW network in scaling

$$R \sim D \sim H^{-1} \sim k^{-1} , \qquad v_{dw} \sim 1$$

• Gravitational waves from a mass distribution

$$h_{ij}(r,t) = \frac{2G}{r}\ddot{I}_{ij}(t-r), \qquad I_{ij} \equiv \int d^3r\rho_{dw} \left(r_i r_j - \frac{r^2 \delta_{ij}}{3}\right)$$

• In the case of DW network in scaling

•

$$R \sim D \sim H^{-1} \sim k^{-1} , \qquad v_{dw} \sim 1$$

$$I_{ij} \sim \rho_{dw} R^5 \qquad \Rightarrow \qquad h_{ij}(r,t) \sim \frac{G}{R} \frac{\rho_{dw} R^5}{R^2} \sim G \sigma H^{-1}$$

• Gravitational waves from a mass distribution

$$h_{ij}(r,t) = \frac{2G}{r} \ddot{I}_{ij}(t-r), \qquad \qquad I_{ij} \equiv \int d^3 r \rho_{dw} \left(r_i r_j - \frac{r^2 \delta_{ij}}{3} \right)$$

• In the case of DW network in scaling

•

•

$$R \sim D \sim H^{-1} \sim k^{-1} , \qquad v_{dw} \sim 1$$

$$I_{ij} \sim \rho_{dw} R^5 \qquad \Rightarrow \qquad h_{ij}(r,t) \sim \frac{G}{R} \frac{\rho_{dw} R^5}{R^2} \sim G\sigma H^{-1}$$
$$\frac{d\rho_{gw}}{d\log k} = \frac{k^3 |\dot{h}|^2}{2(2\pi)^3 Ga^2} \sim G\sigma^2 Ha^{-2}$$

• Gravitational waves from a mass distribution

$$h_{ij}(r,t) = \frac{2G}{r}\ddot{I}_{ij}(t-r), \qquad I_{ij} \equiv \int d^3r\rho_{dw} \left(r_i r_j - \frac{r^2 \delta_{ij}}{3}\right)$$

• In the case of DW network in scaling

$$R \sim D \sim H^{-1} \sim k^{-1} , \qquad v_{dw} \sim 1$$

$$\begin{split} I_{ij} &\sim \rho_{dw} R^5 \qquad \Rightarrow \qquad h_{ij}(r,t) \sim \frac{G}{R} \frac{\rho_{dw} R^5}{R^2} \sim G \sigma H^{-1} \\ &\frac{d\rho_{\rm gw}}{d\log k} = \frac{k^3 |\dot{h}|^2}{2(2\pi)^3 G a^2} \sim G \sigma^2 H a^{-2} \\ \rho_{\rm gw} &= \epsilon_{\rm gw} G \mathcal{A}^2 \sigma^2 \sim \text{CONST} \qquad \qquad \Rightarrow \Omega_{\rm GW} = \frac{G \mathcal{A}^2 \sigma^2}{3 H^2 M_{\rm pl}^2} \end{split}$$

•

•

a

GW signal and leptogenesis

Incompatibility of leptogenesis and detectable GWs

 10^{15} 10^{14} 10^{13} ******* f_a/GeV 10^{-5} 10- 10^{12} 10^{11} 10^{10} 107 10^{8} 10^{9} 10^{10} 10^{11} m_a/GeV

 $v_w = 0.4, \quad c\alpha = 1, \quad \Delta\theta = 10^{-7.3}$

- \bullet amplitude GW signal maximal if $T_{\rm ann} \sim T_{\rm dom}.$
- Suppressions scales like T_{ann} :

$$\Delta \theta \equiv \frac{\Delta V}{V_0} \sim \frac{T_{\rm ann}^2}{m_a M_{\rm pl}} \xrightarrow{T=T_{\rm dom}} \frac{v^2}{M_{\rm pl}^2} \ll 1$$

$$D = \mathrm{Min}\left[1, 0.57 \frac{g_*(T_{\mathrm{ann}})}{g_*(T_{\mathrm{alp}\;\mathrm{dec}})^{1/4}} \frac{\sqrt{M_{\mathrm{Pl}}\Gamma} \; T_{\mathrm{ann}}^3}{\Delta V}\right]$$

• Competition between GW suppression and lepto suppression

[2411.13494]: Mariotti, Nagels, Rase, MV

• Baryon-DM coincidence

 $\Omega_{\rm DM} \sim 5\Omega_b$

• Baryon-DM coincidence

 $\Omega_{\rm DM}\sim 5\Omega_b$

• Typical solution Cogenesis: relate the mechanism of baryogenesis and DM production

$$n_b \sim n_{\rm DM} \qquad \Rightarrow \Omega_{\rm DM} \sim 5\Omega_b \qquad \Rightarrow \frac{m_{\rm DM}}{m_b} \sim 5$$

• Baryon-DM coincidence

 $\Omega_{\rm DM}\sim 5\Omega_b$

• Typical solution Cogenesis: relate the mechanism of baryogenesis and DM production

$$n_b \sim n_{\rm DM} \qquad \Rightarrow \Omega_{\rm DM} \sim 5\Omega_b \qquad \Rightarrow \frac{m_{\rm DM}}{m_b} \sim 5$$

• Domain wall co-genesis

$$\mathcal{L} \supset \underbrace{y(\tilde{H}\bar{L})N_R + \frac{1}{2}M_N\bar{N}_R^c N_R}_{\text{see-saw sector}} + \underbrace{y_D(\phi^{\dagger}\bar{\chi})N_R + m_{\chi}\bar{\chi}\chi}_{\text{DM sector}} + h.c. ,$$

$$\mathcal{L}_{a-j} = c_L \frac{\partial_{\mu}a}{f_a} j_L^{\mu} + c_{\chi} \frac{\partial_{\mu}a}{f_a} j_{\chi}^{\mu}$$

• Baryon-DM coincidence

 $\Omega_{\rm DM}\sim 5\Omega_b$

• Typical solution Cogenesis: relate the mechanism of baryogenesis and DM production

$$n_b \sim n_{\rm DM} \qquad \Rightarrow \Omega_{\rm DM} \sim 5\Omega_b \qquad \Rightarrow \frac{m_{\rm DM}}{m_b} \sim 5$$

• Domain wall co-genesis

$$\mathcal{L} \supset \underbrace{y(\tilde{H}\bar{L})N_R + \frac{1}{2}M_N\bar{N}_R^cN_R}_{\text{see-saw sector}} + \underbrace{y_D(\phi^{\dagger}\bar{\chi})N_R + m_{\chi}\bar{\chi}\chi}_{\text{DM sector}} + h.c.,$$

$$\mathcal{L}_{a-j} = c_L\frac{\partial_{\mu}a}{f_a}j_L^{\mu} + c_{\chi}\frac{\partial_{\mu}a}{f_a}j_{\chi}^{\mu}$$
endows and $(\mu = R > 1 < 1 = 1)$

• Different hierarchies: $y_D/y \equiv R > 1, < 1, = 1.$

$$\mathsf{Br}_{N \to \phi \chi} = \frac{R^2}{1 + R^2}, \qquad \mathsf{Br}_{N \to HL} = \frac{1}{1 + R^2} \qquad \Gamma_{\rm tot}^D = \frac{y^2 (1 + R^2) M_N}{16\pi}.$$

[2411.13494]: Mariotti, Nagels, Rase, MV



[2411.13494]: Mariotti, Nagels, Rase, MV





• DW-genesis: generation of baryon number independent on CP-violation in the couplings



- DW-genesis: generation of baryon number independent on CP-violation in the couplings
- Suppressions from matter domination and opposite DWs



- DW-genesis: generation of baryon number independent on CP-violation in the couplings
- Suppressions from matter domination and opposite DWs
- DWs are copious source of GWs



- DW-genesis: generation of baryon number independent on CP-violation in the couplings
- Suppressions from matter domination and opposite DWs
- DWs are copious source of GWs
- observable GW signal and successful DW-genesis seem mutually exclusive



- DW-genesis: generation of baryon number independent on CP-violation in the couplings
- Suppressions from matter domination and opposite DWs
- DWs are copious source of GWs
- observable GW signal and successful DW-genesis seem mutually exclusive
- Model of cogenesis with $Y_B \neq Y_{\chi}$!