

The sunny side of direct detection

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HIDDeN VIRTUAL INSTITUTE webinar

The sunny side of direct detection

- Introduction: dark matter as the key agent in cosmological structure formation
- Introduction of dark matter direct detection
- New, irreducible signal components in direct detection

Bremsstrahlung: Kouvaris, JP PRL 2016;

Migdal effect: Essig, JP, Sholapurkar, Yu PRL 2020

Solar reflection I: An, Pospelov, JP, Ritz PRL 2018

Solar reflection II: An, Nie, Pospelov, JP, Ritz PRD 2021

Dark matter: now and then



Jim Peebles, Oct 8 2019

2019 Nobel Prize “for theoretical discoveries in physical cosmology”

Early Cosmology

Predicting the cosmic microwave background

Search for origin of chemical elements => idea in the 1940's that matter passed through a hot and dense phase and all elements were made

Alpher, Bethe, Gamow 1948

Number of reactions = cross section (cm^2) \times flux of projectiles ($1/\text{cm}^2/\text{s}$) \times time (s)

key reaction $p + n \rightarrow d + \gamma$

per neutron ~ not too small number to make elements

~ not too large number so that significant amounts of D and light elements are left

=> $\langle\sigma v\rangle nt \sim 1$ at $T \sim \text{MeV} \sim 10^9 \text{ K}$ => **Universe filled with radiation**

with $\langle\sigma v\rangle \sim \text{const.}$ yields a prediction of the required baryon density $n \sim 10^{18}/\text{cm}^3$

=> With an estimate of today's density $n \sim 10^{-7}/\text{cm}^3$ together with $T \propto 1/a(t) \propto n^{1/3}$

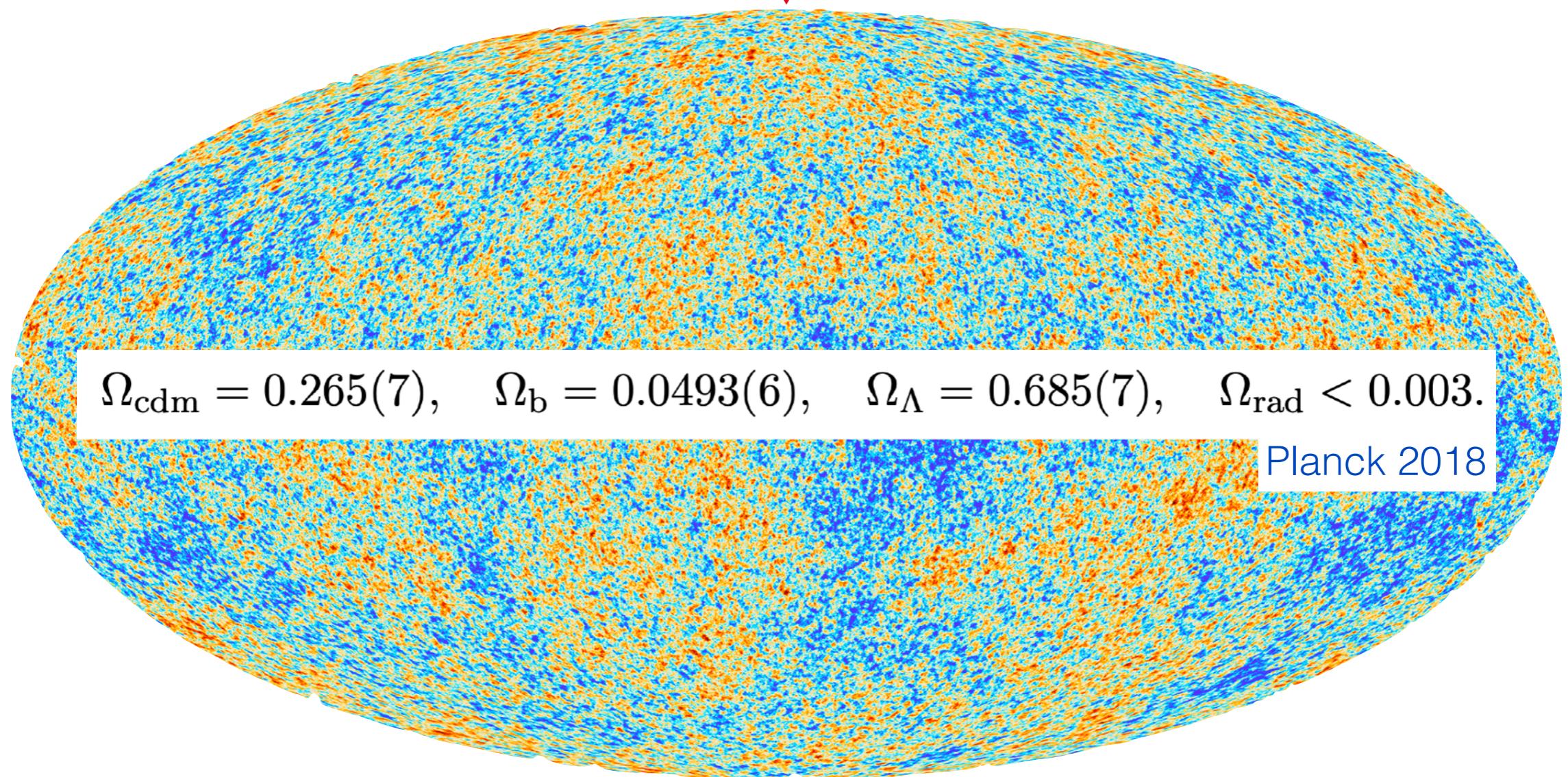
=> **T~4K today's radiation temp.**

Cosmological scales

Cosmic Microwave Background



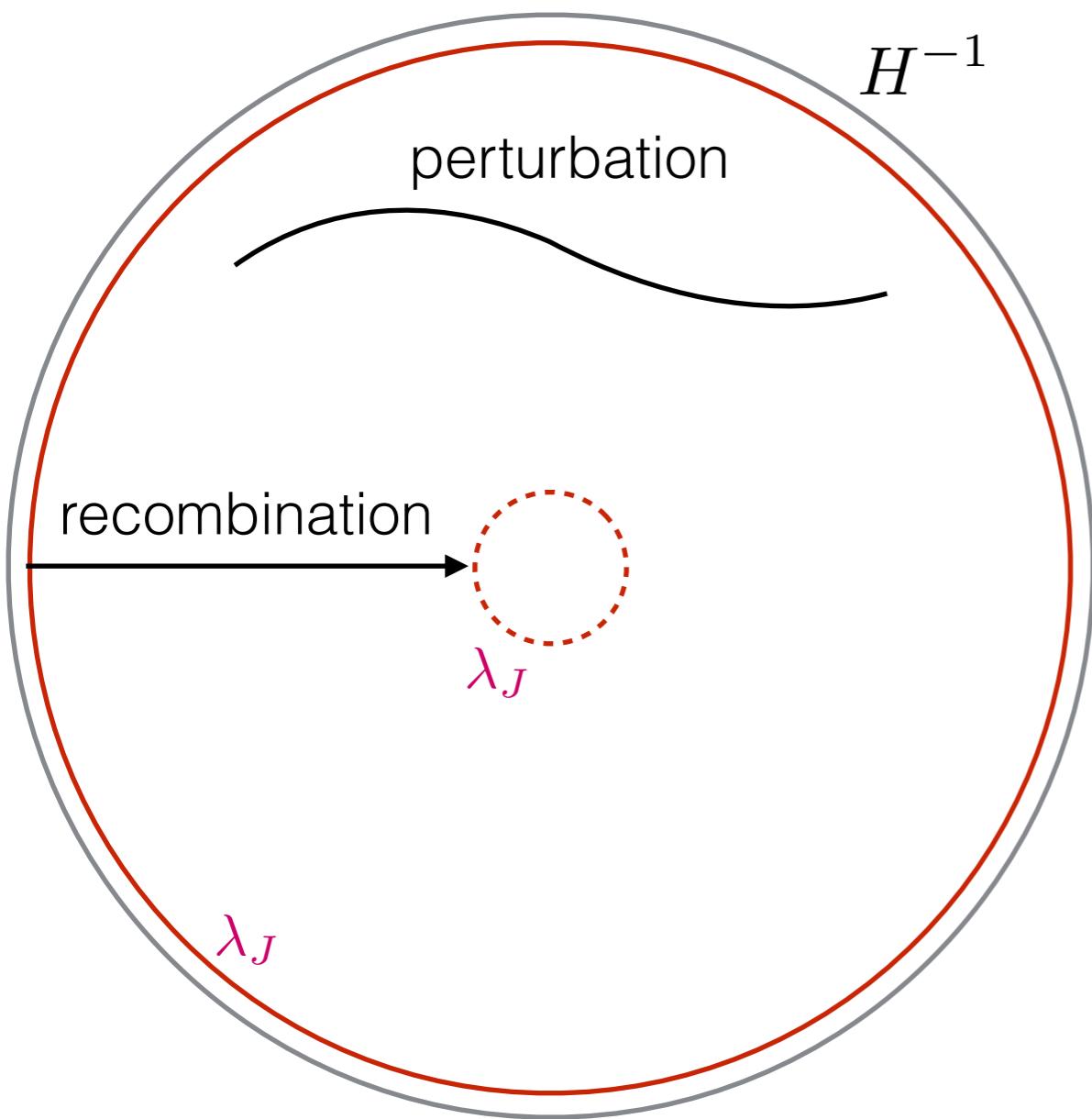
Dark Matter is the key ingredient to explain the CMB observations (linear theory)



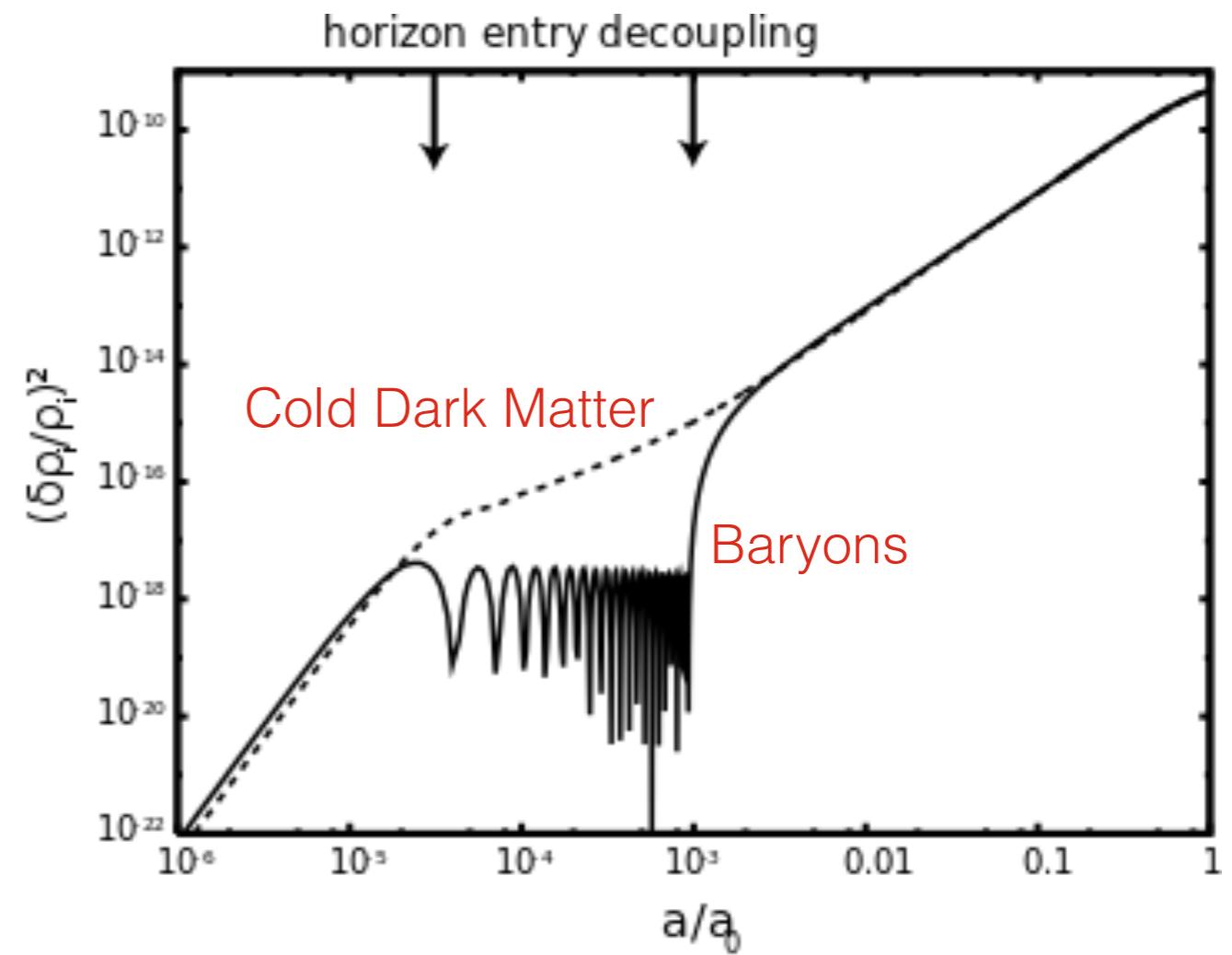
Cosmological scales

Origin of structure

$$\ddot{\delta} + [\text{Pressure} - \text{Gravity}] \delta = 0$$



baryons fall into the potential wells created by dark matter.

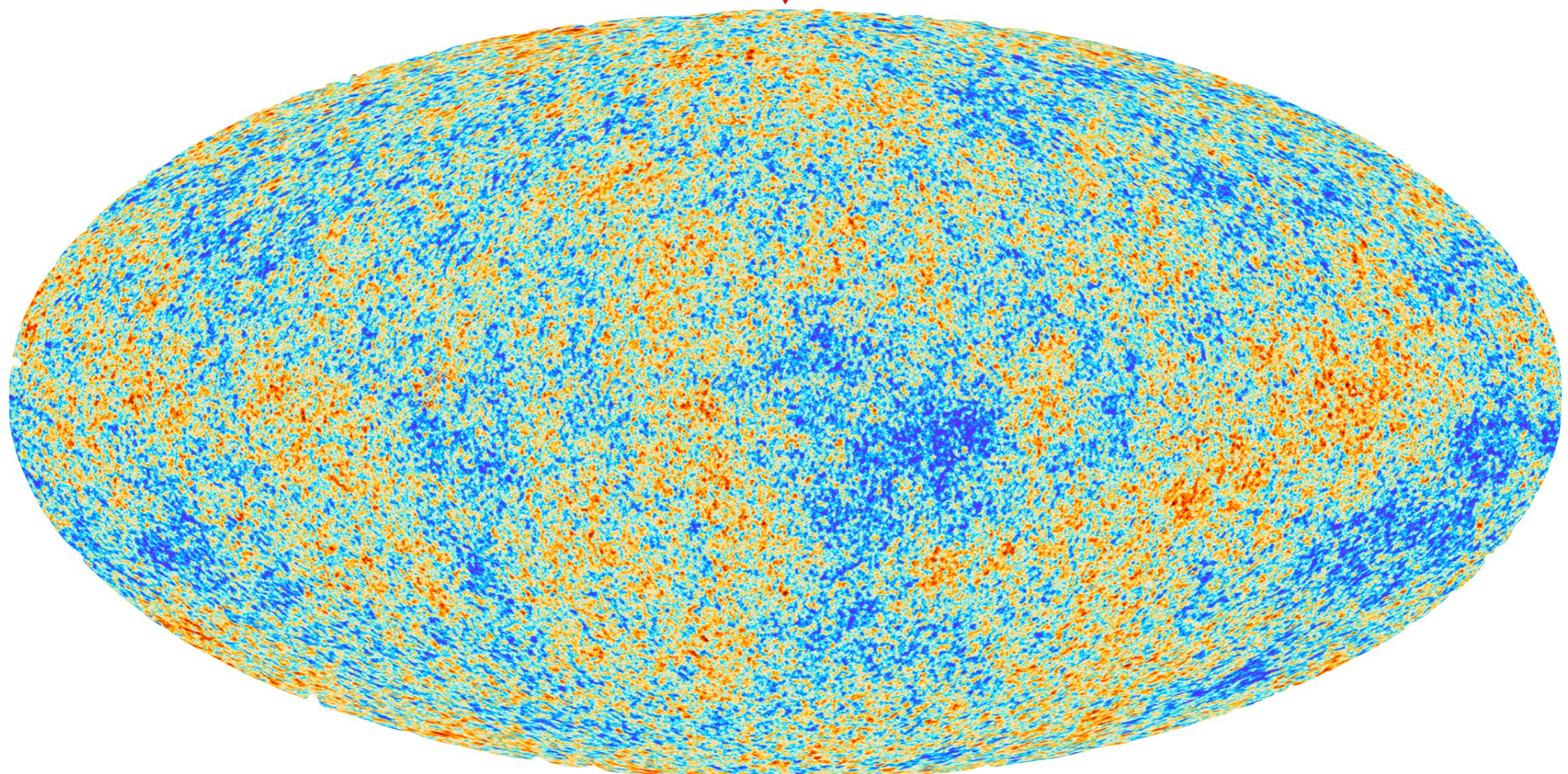


Cosmological scales

Cosmic Microwave Background



Dark Matter is the key ingredient to explain the CMB observations (linear theory)



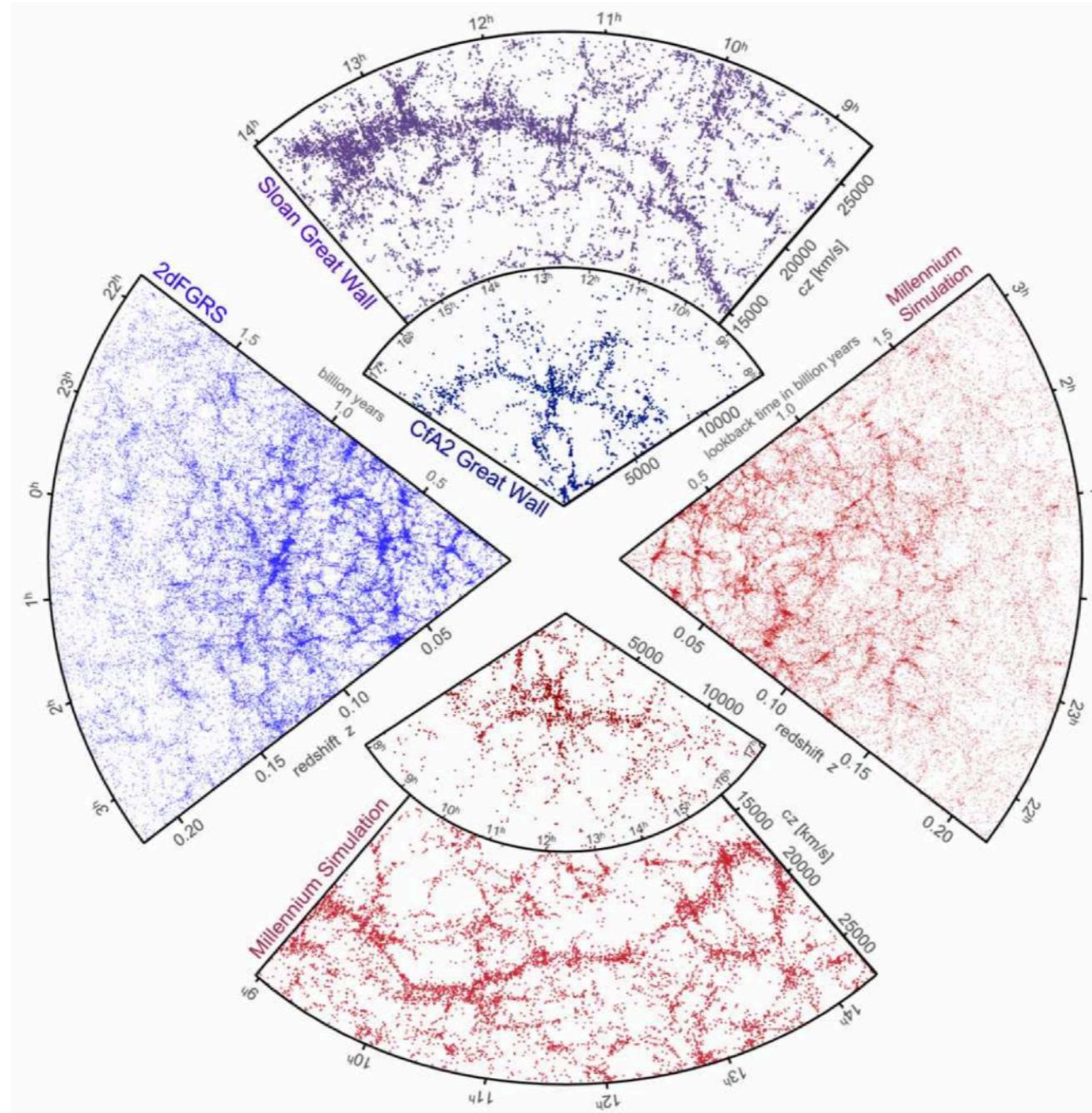
Dark Matter is the key catalyst for formation of non-linear large scale structure

Large Scale Structure

From the CMB epoch to today



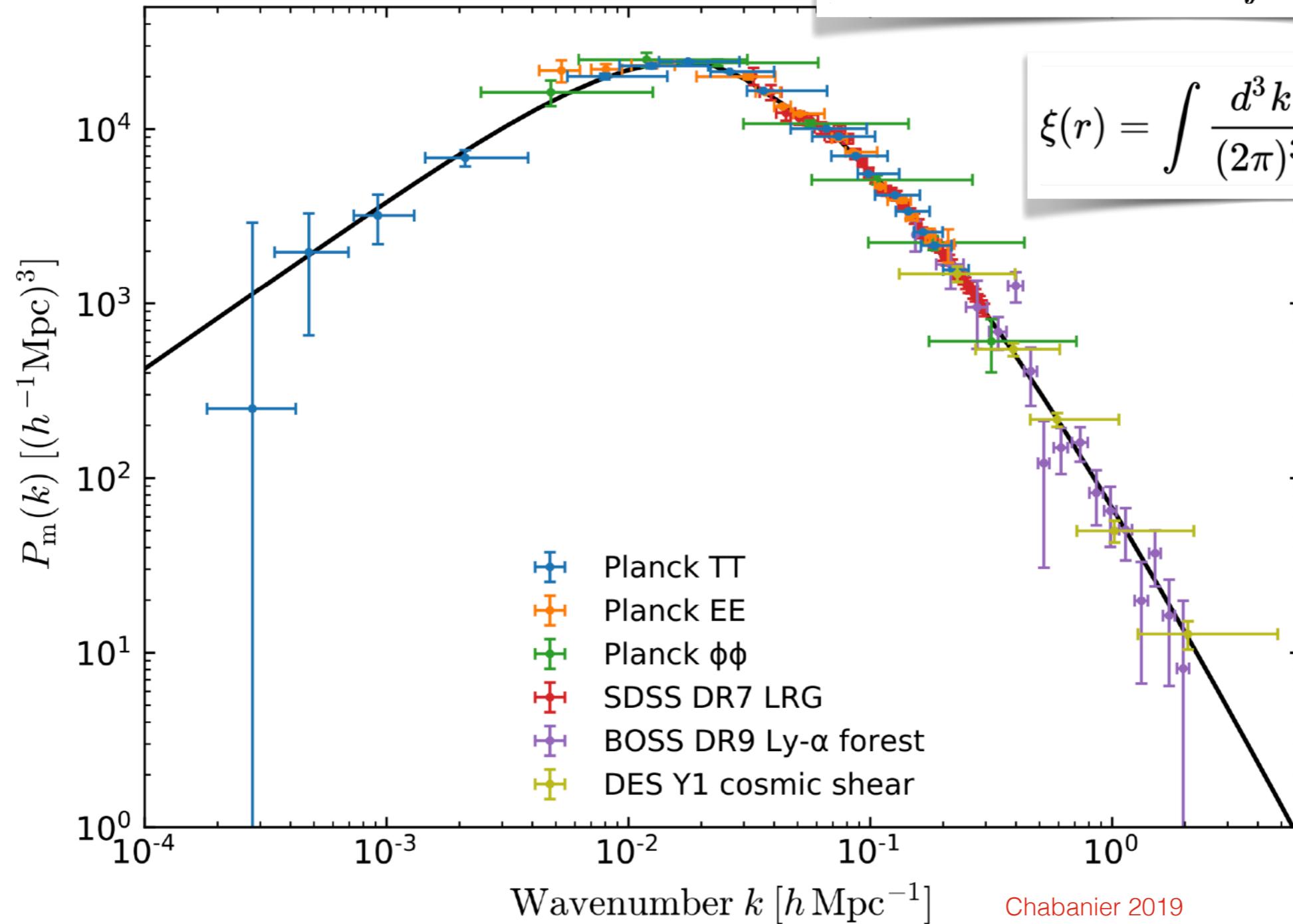
Dark Matter is the key catalyzer for formation of non-linear large scale structure



Large Scale Structure

$$\delta(\mathbf{x}) = \frac{\rho(\mathbf{x}) - \bar{\rho}}{\bar{\rho}}$$

Matter power spectrum



$$\xi(r) = \langle \delta(\mathbf{x})\delta(\mathbf{x}') \rangle = \frac{1}{V} \int d^3\mathbf{x} \delta(\mathbf{x})\delta(\mathbf{x} - \mathbf{r})$$

$$\xi(r) = \int \frac{d^3k}{(2\pi)^3} P(k) e^{i\mathbf{k}\cdot(\mathbf{x}-\mathbf{x}')}$$

Dark Matter in the Milky Way

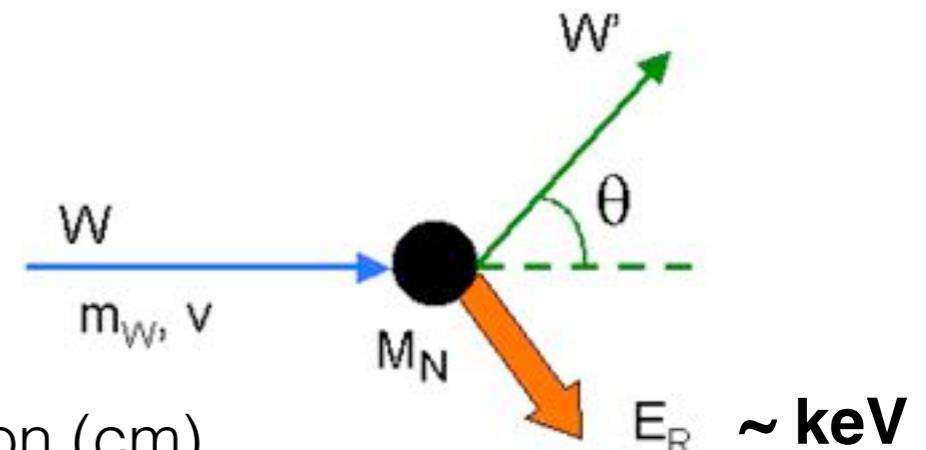


=> broader topic of this talk: non-gravitational detection of that fluffy cloud

DM Direct Detection

Basic idea

Detection Rate = particle flux ($1/\text{cm}^2/\text{sec}$) \times cross section (cm)



$$\frac{dR(t)}{dE_R} = N_T \frac{\rho_0}{m_{\text{DM}}} \int_{v \geq v_{\min}} d^3 \mathbf{v} v f_{\text{LAB}}(\mathbf{v}) \frac{d\sigma}{dE_R} \quad [\text{cpd/kg/keV}]$$

Astrophysics

local DM density



DM velocity distribution in the LAB frame



recoil cross section

Particle Physics

Contributions to $f(v)$

- virialized component
- substructure ($p < 10^{-4}$)
- debris flow, streams

Very little is known

- largely dissipationless
- stable on cosmological timescales

DM Direct Detection

Velocity distribution of DM in the halo

Maxwellian velocity distribution is found in N-body simulations

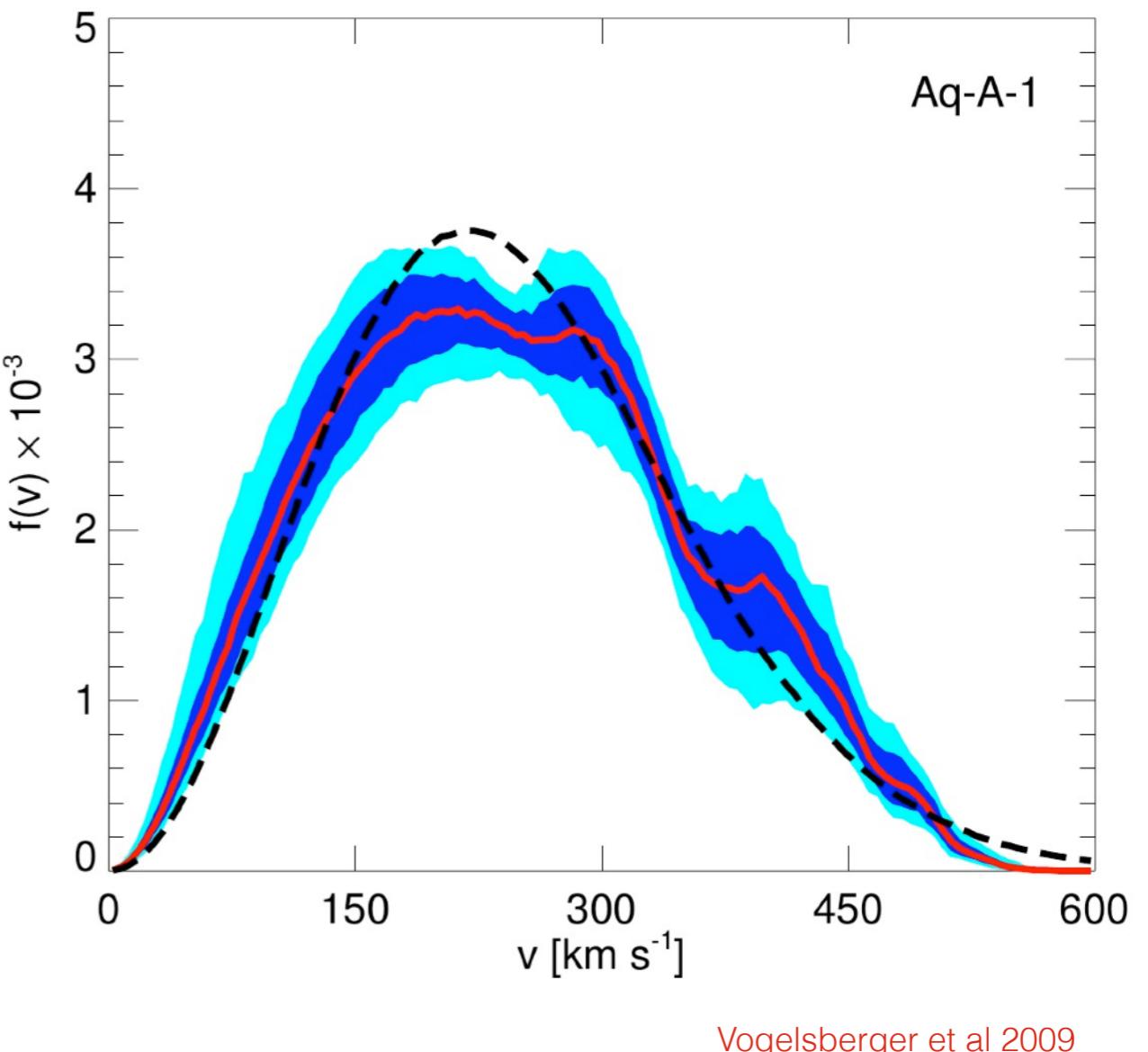
Note: not all particles of arbitrary velocity can be gravitationally bound to the halo

$$f_{\text{gal}}(\vec{v}) \approx \begin{cases} N \exp(-|\vec{v}|^2/v_0^2) & v < v_{\text{esc}} \\ 0 & v > v_{\text{esc}} \end{cases}$$

$$v_{\text{esc}} \simeq 650 \text{ km/s}$$

- Local DM flux is ($v_\chi \sim 10^{-3}c$)

$$\phi_\chi \sim \frac{\rho_0 v_\chi}{m_\chi} \sim 10^5 / \text{cm}^2/\text{s} \left(\frac{100 \text{ GeV}}{m_\chi} \right)$$



Simple options for dark sectors

Connecting new physics to the Standard Model

$$(H^\dagger H) (A\phi + \lambda\phi^2)$$

“Higgs Portal”
(a minimal model of DM)

$$LH\textcolor{red}{N}$$

“Neutrino Portal”
likely realized in nature (neutrinos have mass); sterile neutrinos

$$F_{\mu\nu}^Y \textcolor{red}{V}^{\mu\nu}$$

“Vector Portal”
kinetic mixing of abelian field strength tensors

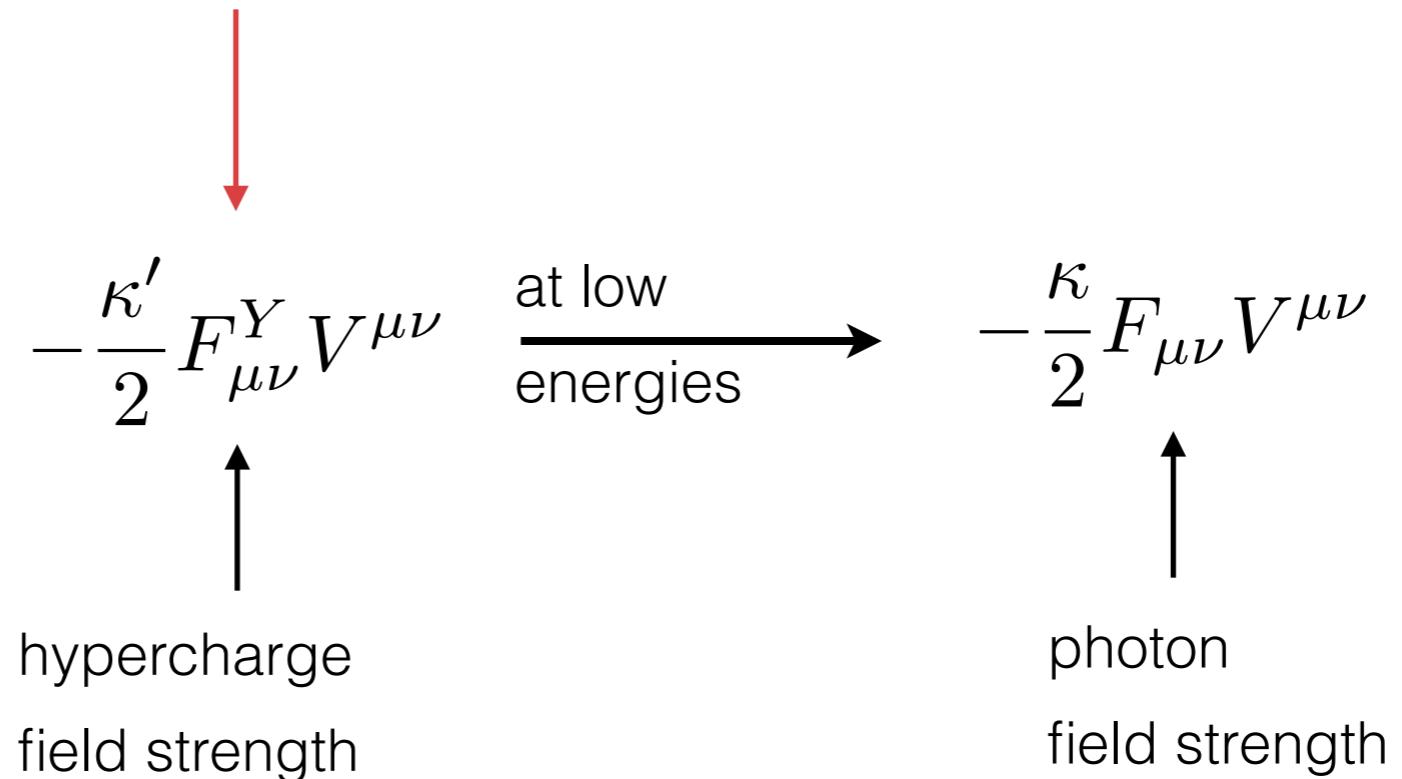
The kinetic mixing portal

“Dark Photons”

$$\text{SU}(3)_c \times \text{SU}(2)_L \times \text{U}(1)_Y \times U(1)'$$

Standard Model

× “dark sector” with vector particle V^μ



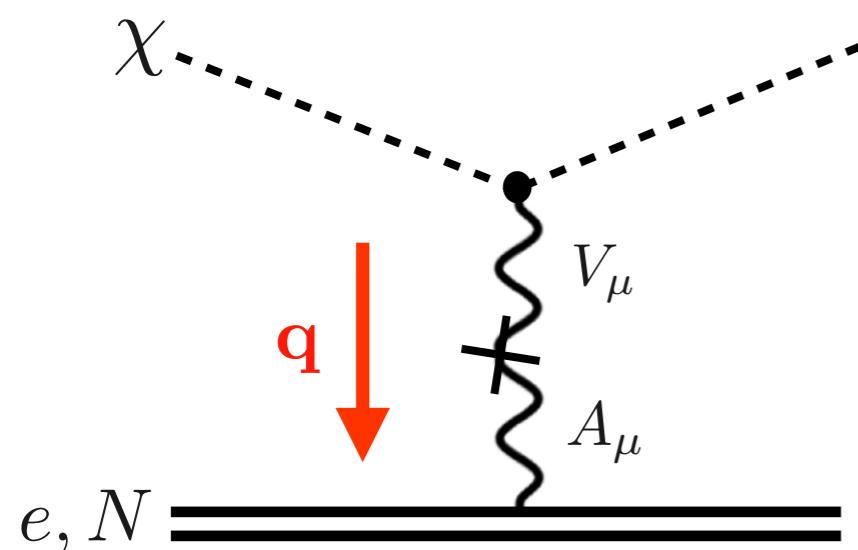
=> two parameter theory: kinetic mixing strength and mass of V

DM Direct Detection

Simple example

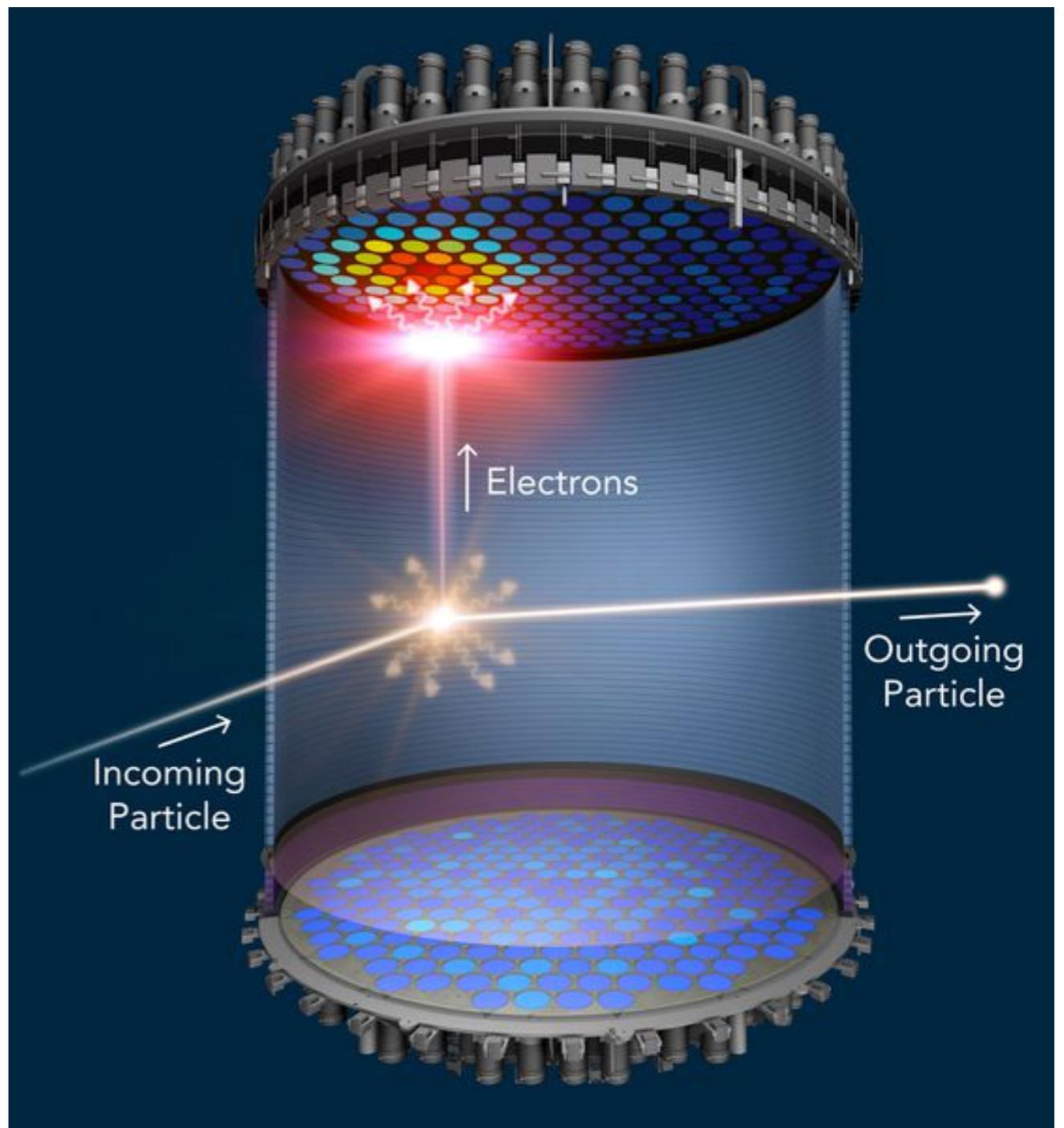
χ is the dark matter

V_μ dark photon is the “mediator”



$m_V \gg |q|$ point-like interaction

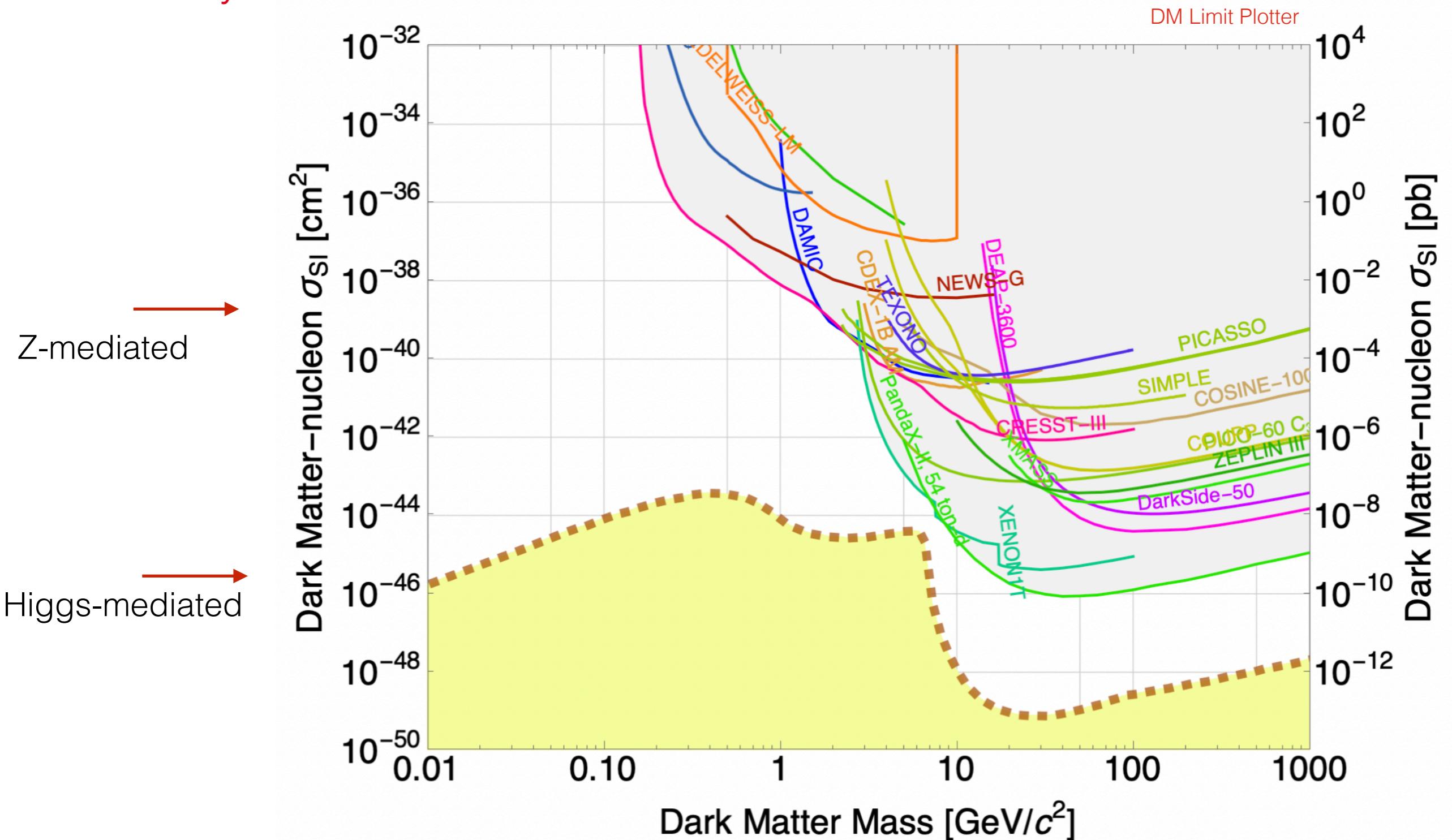
$m_V \ll |q|$ “millicharged DM”



general detection principles:
ionization, scintillation, heat, ...

DM Direct Detection

A summary of two decades of effort

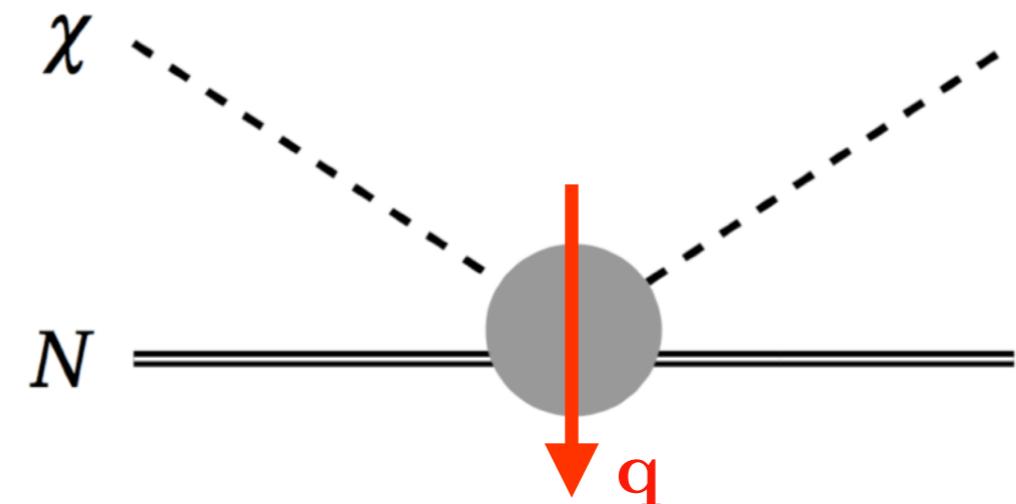


DM Direct Detection

A closer look

Nuclear kinetic recoil energy

$$E_R = \frac{\mathbf{q}^2}{2m_N} = \frac{\mu_N^2 v^2}{m_N} (1 - \cos \theta_*)$$



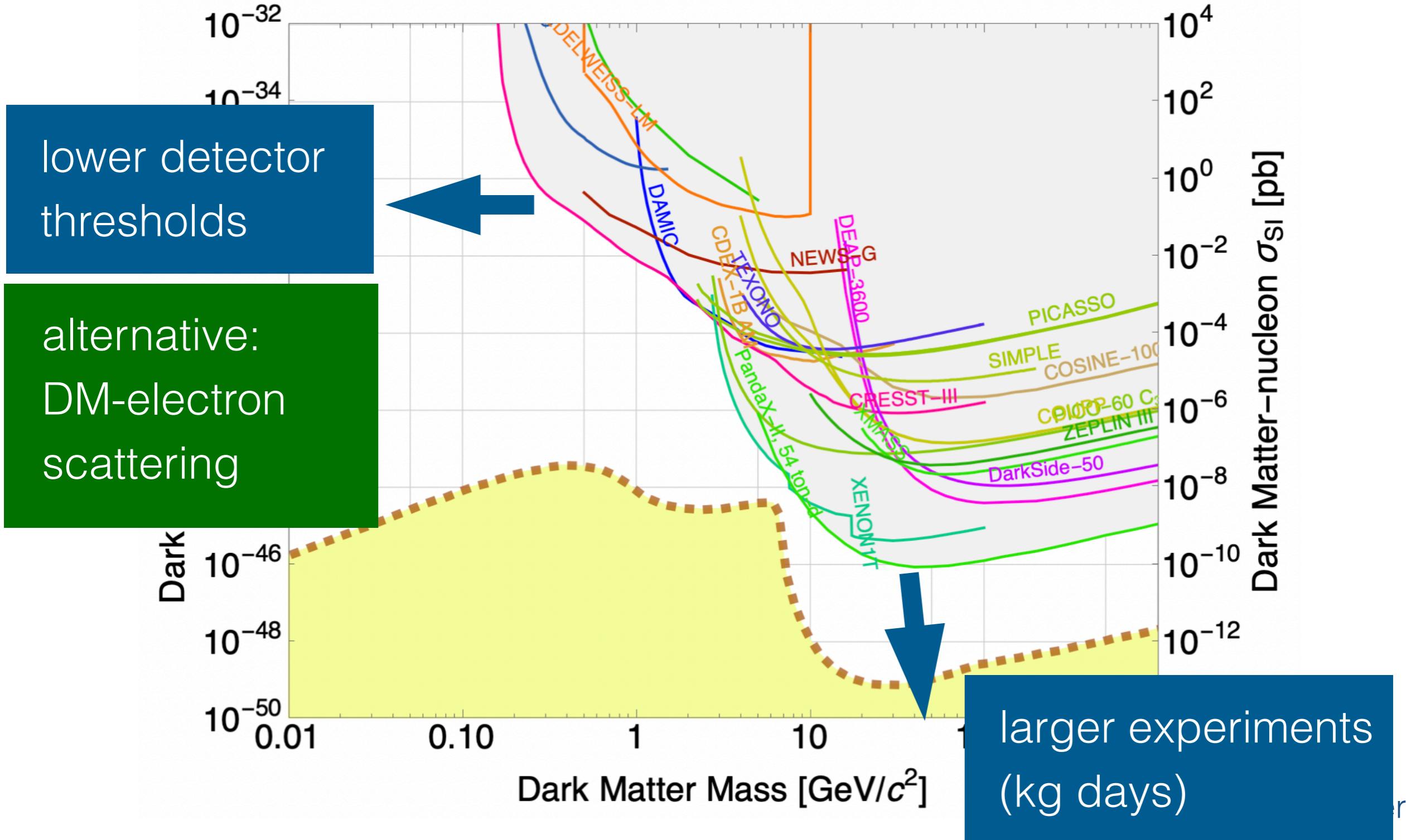
=> A given recoil, demands a minimum relative velocity

$$v_{\min} = \sqrt{\frac{m_N E_R}{2\mu_N^2}} \simeq \left(\frac{E_R}{0.5 \text{ keV}} \right)^{1/2} \frac{1 \text{ GeV}}{m_\chi} \times \begin{cases} 1700 \text{ km/s} & \text{Xenon} \\ 600 \text{ km/s} & \text{Oxygen} \end{cases}$$

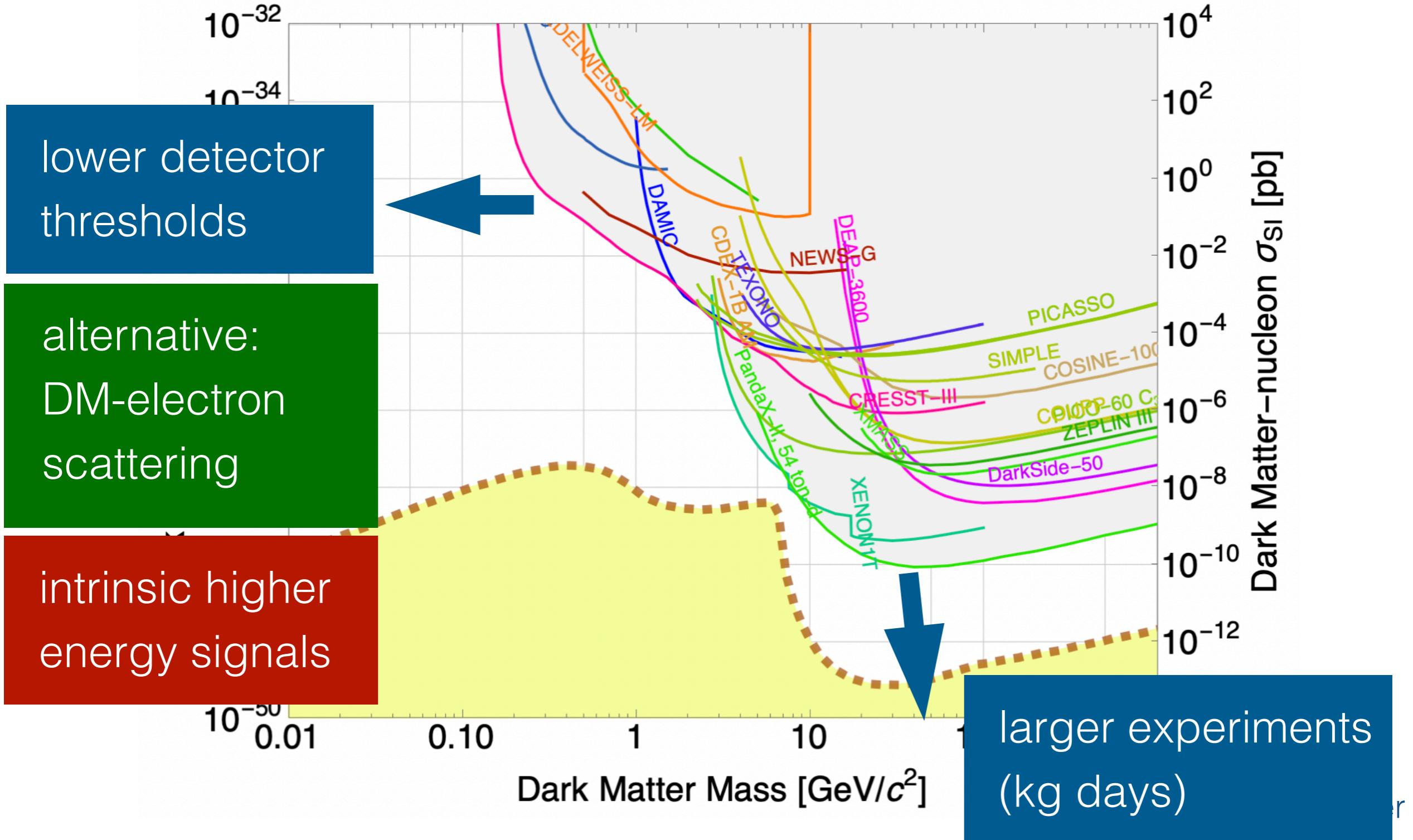
=> if $m < 1 \text{ GeV}$, then there are no particles bound to the Galaxy that could induce a 0.5 keV nuclear recoil on a Xenon atom!

“kinematical no-go theorem”

Direct detection low-mass frontier



Direct detection low-mass frontier



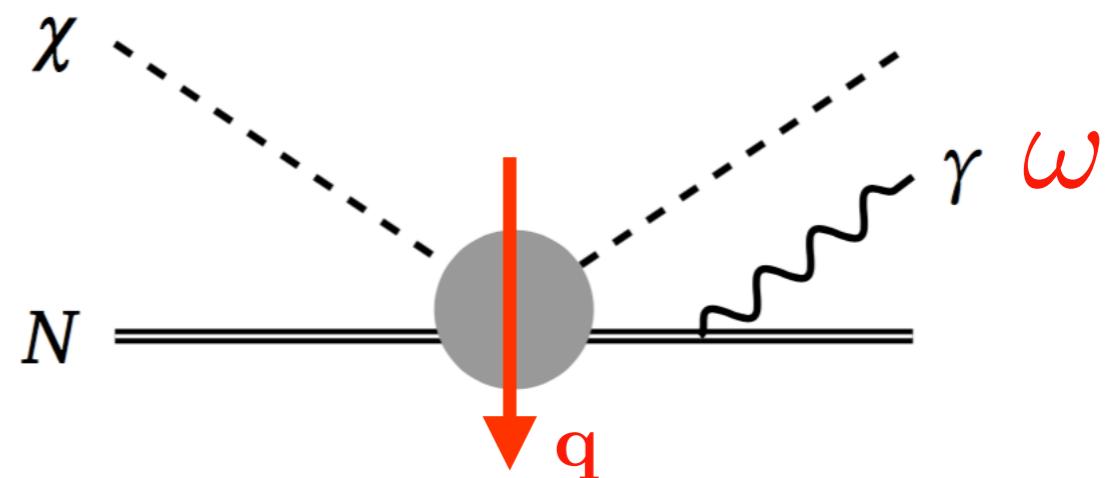
Gaining access to sub-GeV Dark Matter *through nuclear recoils*

Inelastic channel of photon emission from the nucleus

Maximum photon energy

$$\omega_{\max} \simeq \mu_N v^2 / 2 \simeq m_\chi v^2 / 2$$

$$\simeq 0.5 \text{ keV} \frac{m_\chi}{100 \text{ MeV}}$$

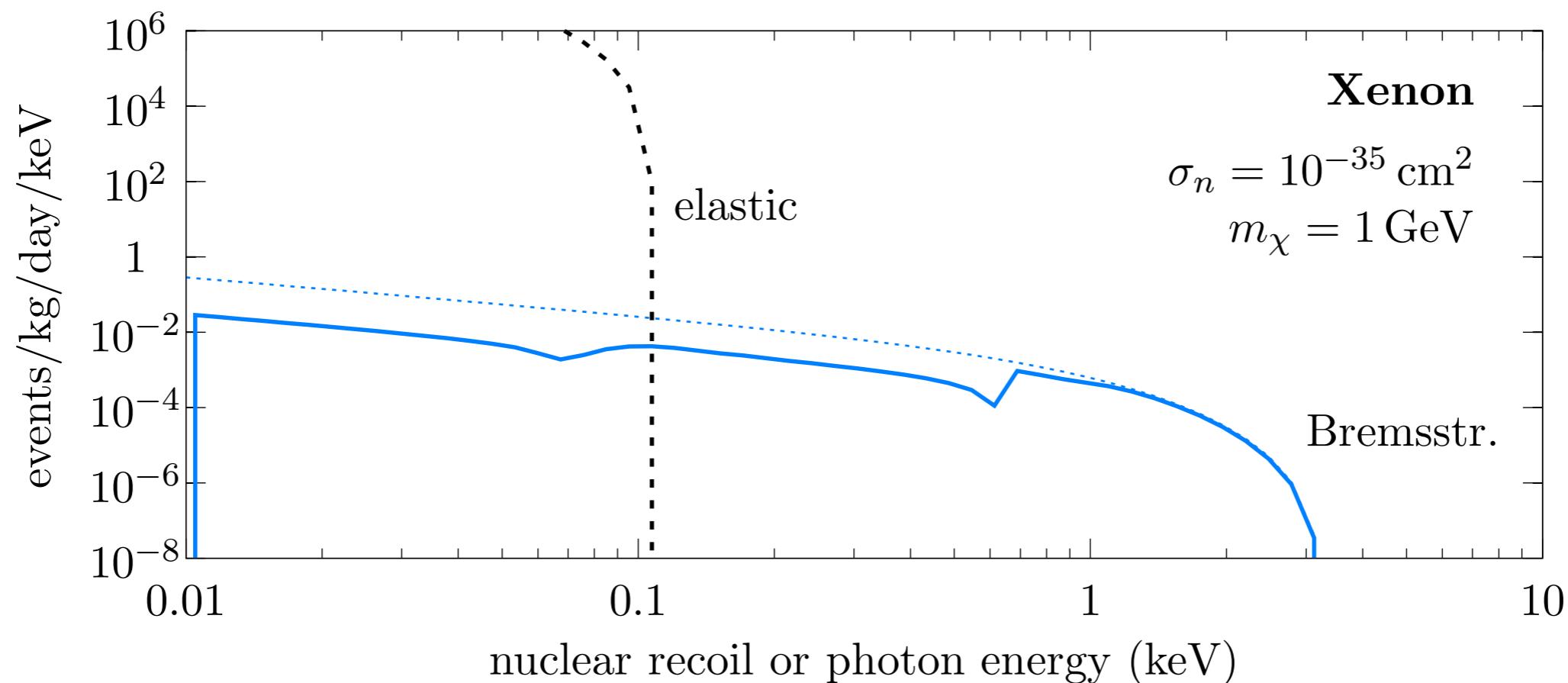
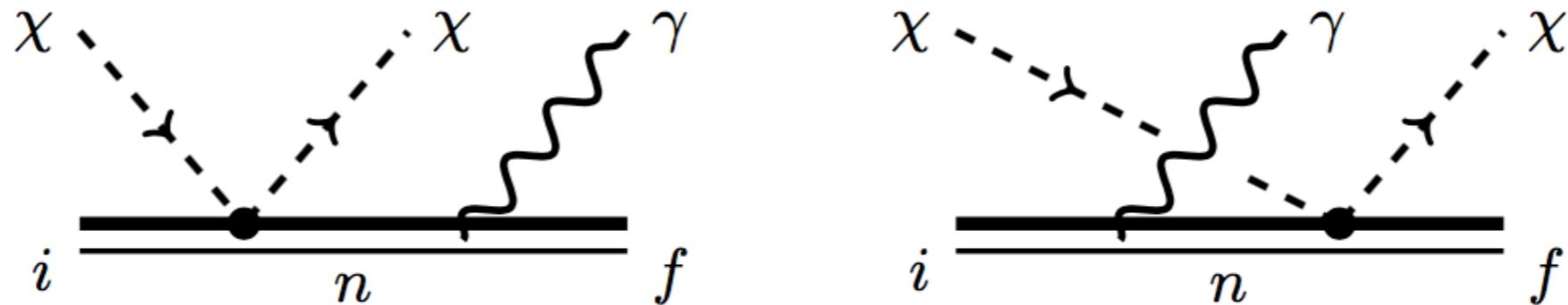


Key I: $E_{R,\max} = 4(m_\chi/m_N)\omega_{\max} \ll \omega_{\max}$ ($m_\chi \ll m_N$)

Key II: 0.5 keV nuclear recoil is easily missed (heat losses),
0.5 keV photon is never missed!

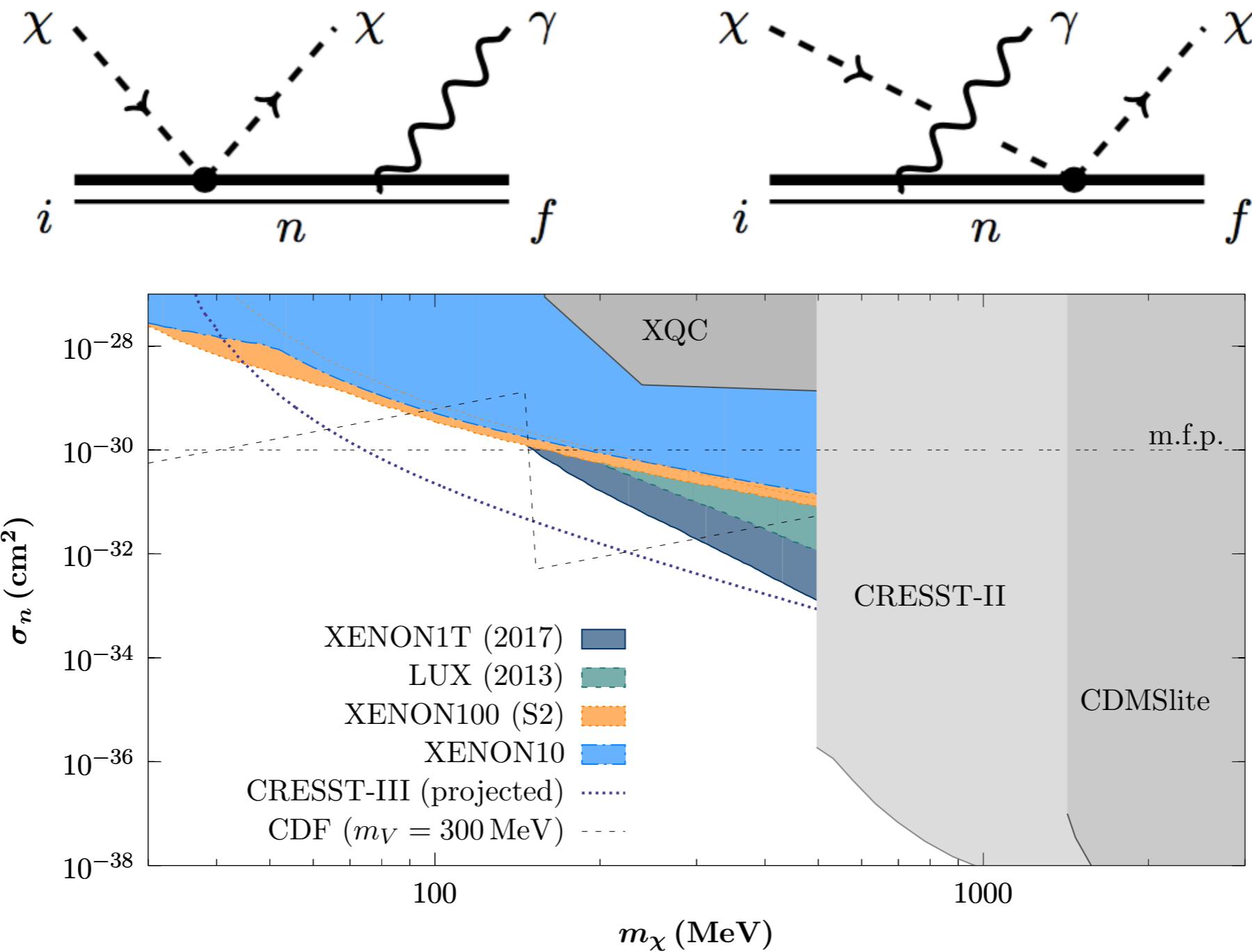
Irreducible signal components

Example Bremsstrahlung



Irreducible signal components

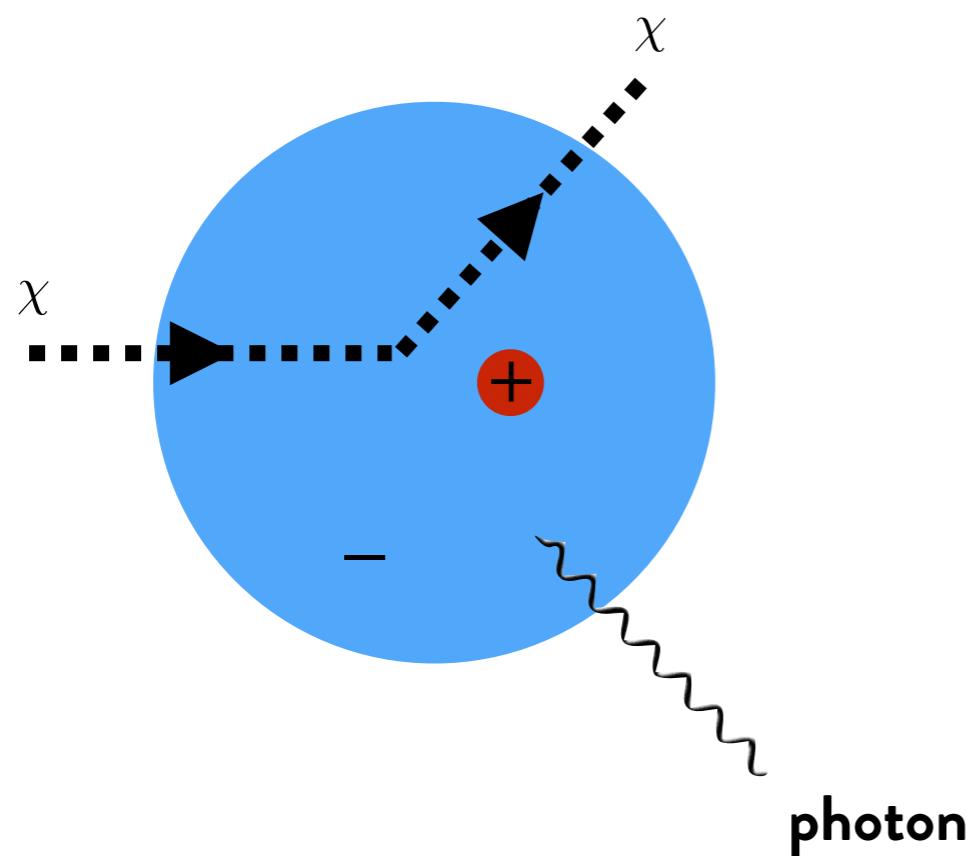
Example Bremsstrahlung



Kouvaris, JP PRL 2016

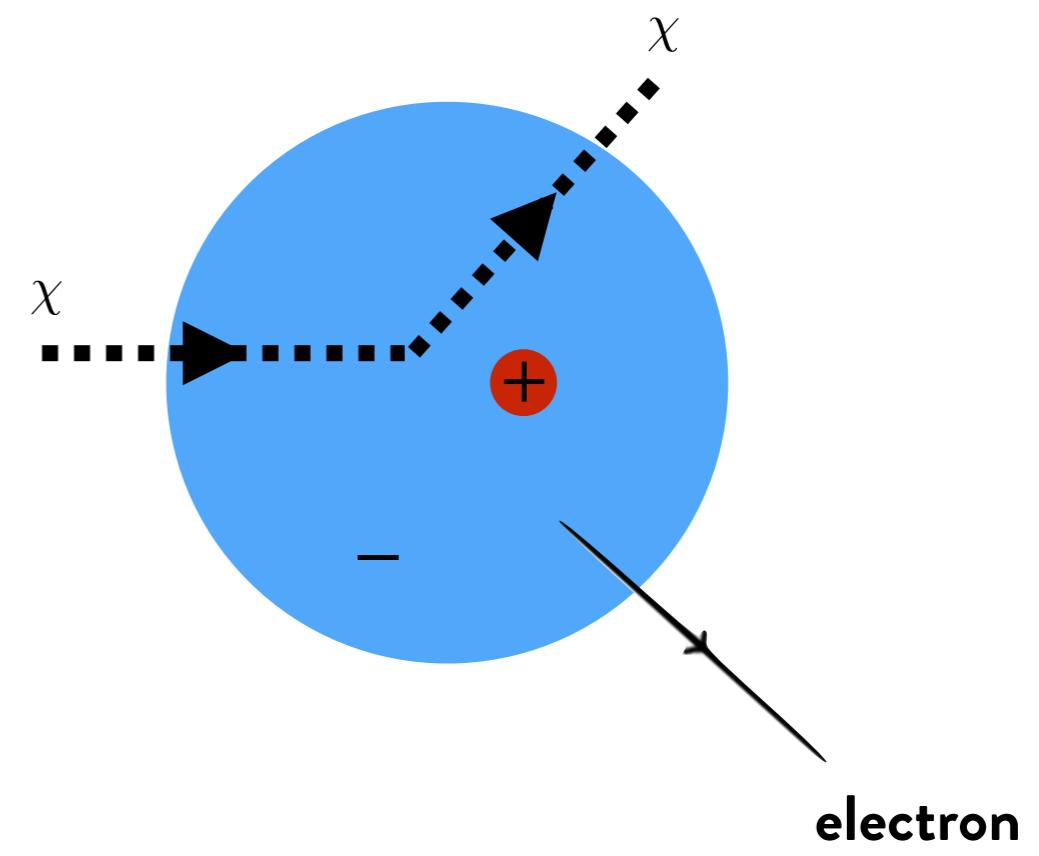
Irreducible signal components

Prompt atomic response following nuclear recoil



Bremsstrahlung

Kouvaris, JP PRL 2016



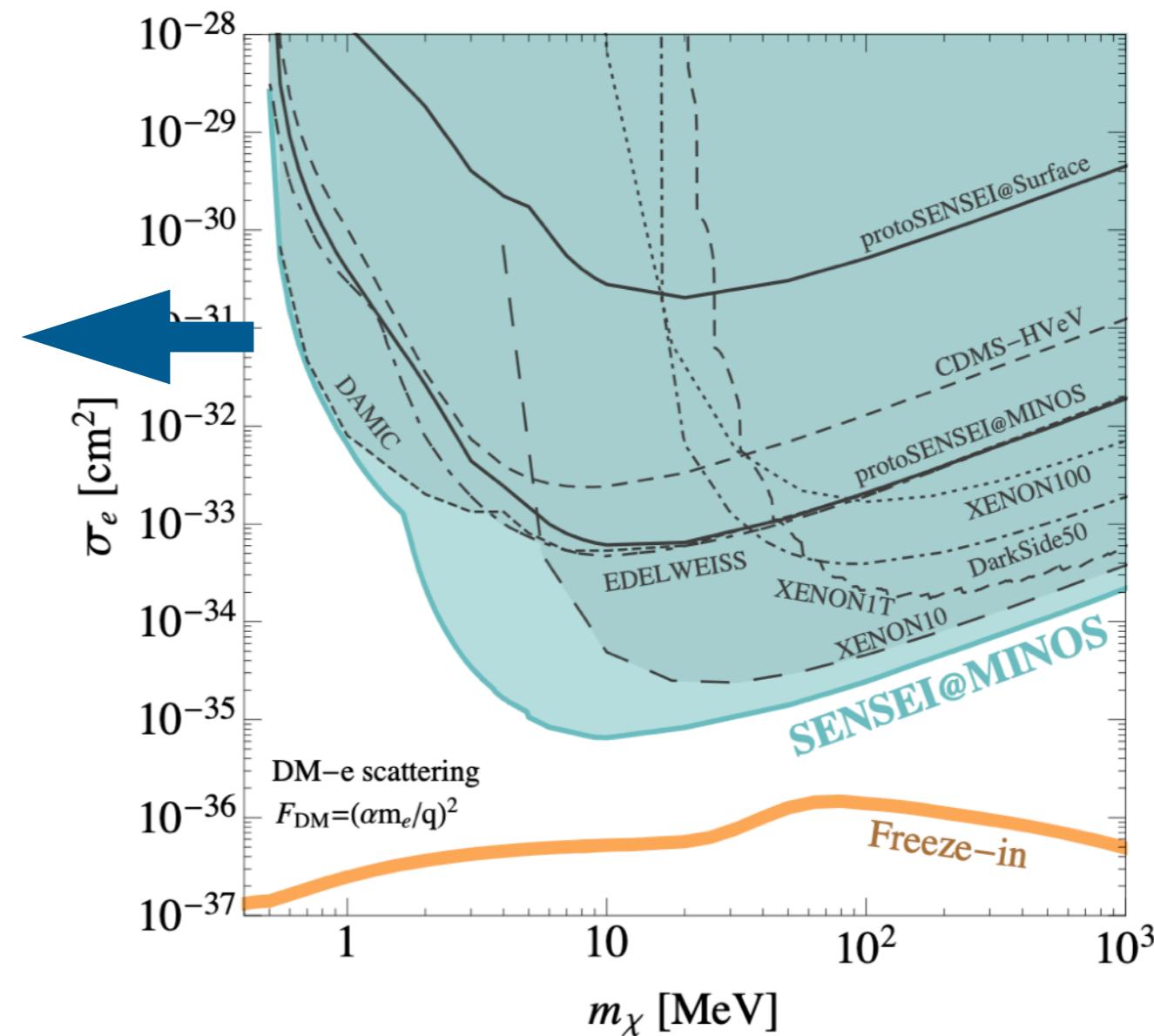
Migdal-effect

Ibe et al JCAP 2017
Essig, JP, Sholapurkar, Yu PRL 2020

Direct Detection using electrons

alternative
technology
(e.g. superconductors)

irreducible
higher energy
components



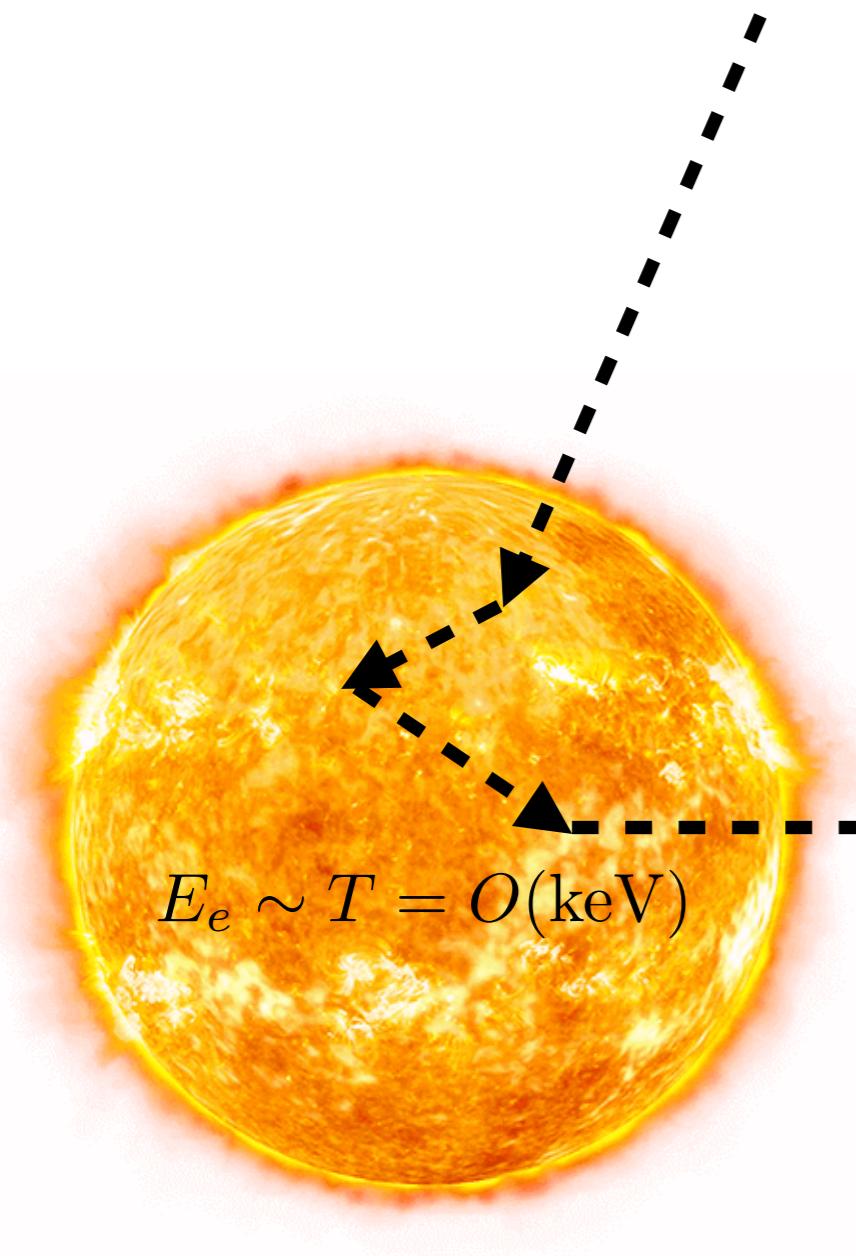
E.g. if $m < 10 \text{ MeV}$, then there are no particles bound to the Galaxy
that could ionize an outer shell Xenon electron

"kinematical no-go theorem" #2

The sun as particle accelerator

Scattering on hot electrons in the solar interior

Galactic Dark Matter $v_{\max} \sim 600 \text{ km/sec}$



Elastic scattering on electrons in the sun's interior lifts DM kinetic energy into (sub-)keV regime

$$E_{\text{DM}}^{\text{recoil}} \sim T$$

reflected spectrum

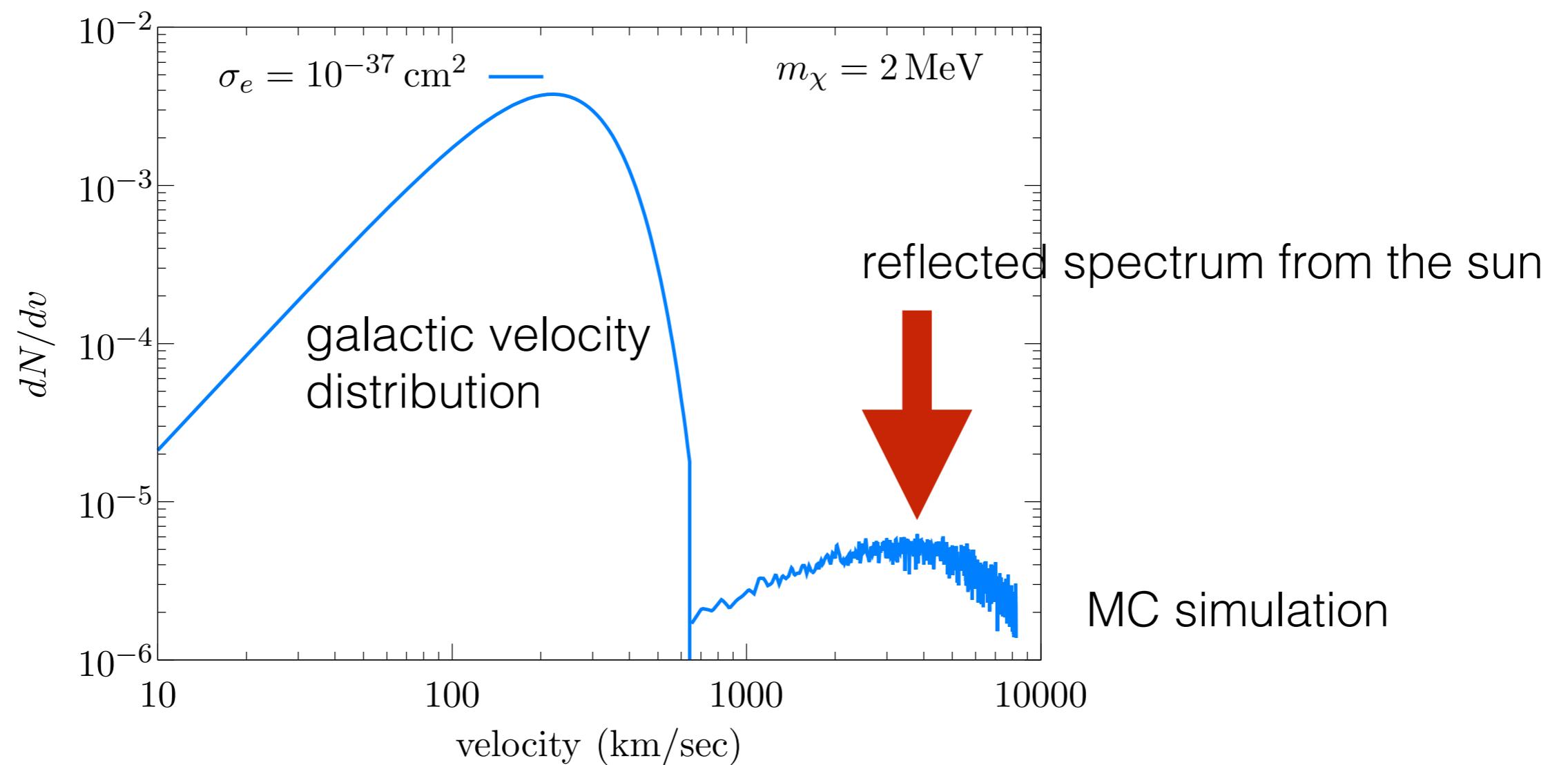
$$v_{\max} = \sqrt{2T/m}$$

$$\sim 10^4 \text{ km/sec} \times \sqrt{1 \text{ MeV}/m}$$



The sun as particle accelerator

Adding a “hot” tail to the Maxwellian



The sun as particle accelerator

Single scattering limit

mean free path of DM in the sun $l_{\text{fp}} = [n_e \langle \sigma_{\text{tot}} v_r \rangle]^{-1} \bar{v}_\chi$

probability of scattering $P_s \sim R_{\text{traj}} / l_{\text{fp}} \sim$
 $\sim \rho_{\text{core}} / m_p \times R_{\text{traj}} \times \sigma_{\text{tot}} \times \frac{\bar{v}_e}{\bar{v}_\chi}$
 $\sim \frac{\sigma_{\text{tot}}}{10^{-38} \text{ cm}^2}$ ballpark number for solar reflection

reflected flux $\frac{d\Phi_{\text{reflected}}}{dE_\chi} = \Phi_{\text{halo}} \times \frac{F_{A_\rho} A_\rho}{4\pi(\text{A.U.})^2}$ $A_\rho = \pi(4R_\odot)^2, \int dE_\chi F_{A_\rho} = 1$
=> solid angle suppression $\sim 10^{-4}$

Direct Detection of sub-MeV DM

Example model with contact interactions

UV completed through Z' where relic density is set via p-wave annihilation and safe from CMB constraints on energy injection (modulo model dependent N_{eff} contributions)

$$\mathcal{L}_{\text{int}} = G_{\chi e} \times (\bar{e} \gamma^\mu e) (i\chi^* \partial_\mu \chi - i\chi \partial_\mu \chi^*)$$

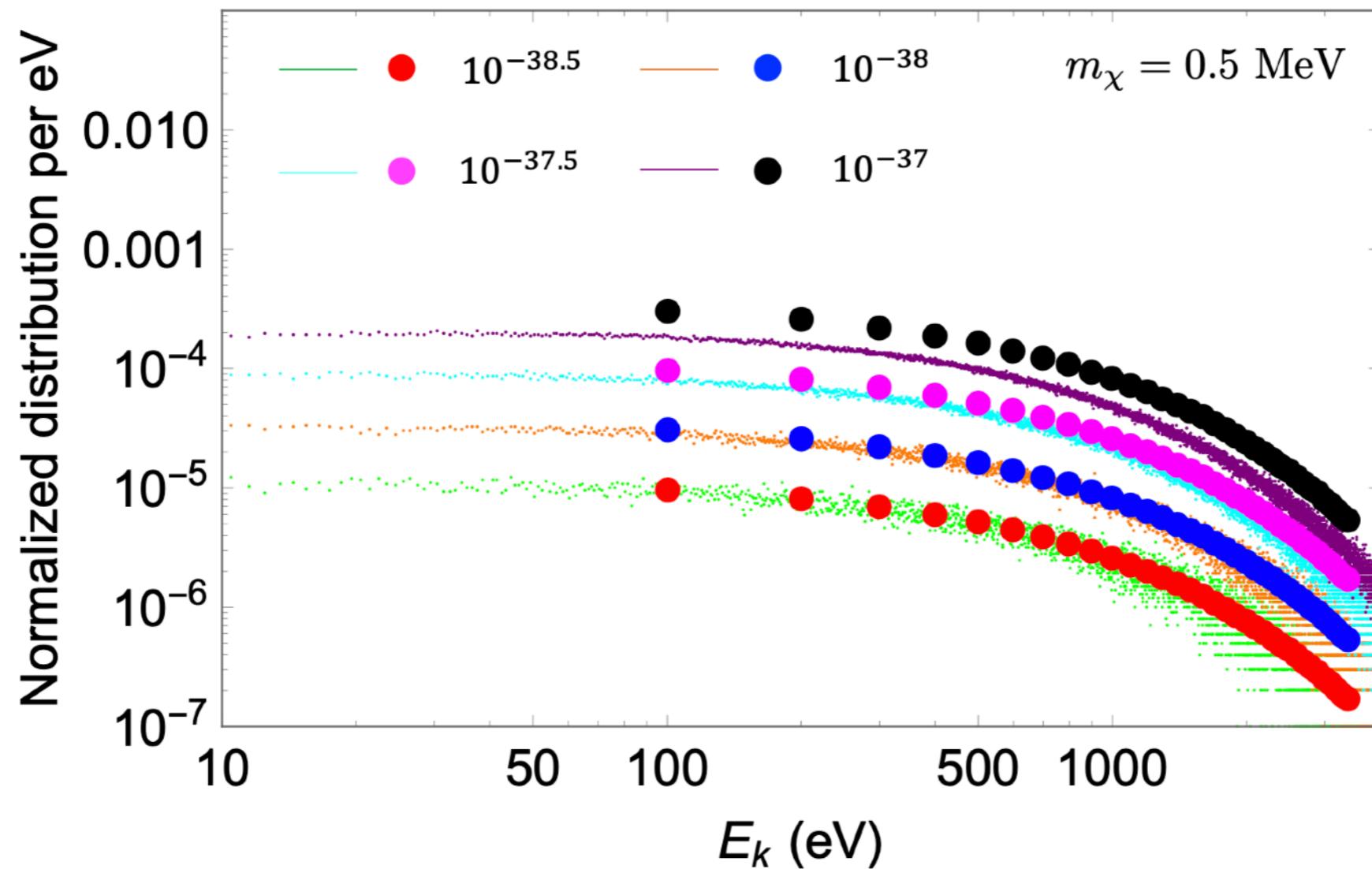
$$\sigma_{\text{ann}} v = v^2 \times \frac{G_{\chi e}^2}{12\pi} (m_e^2 + 2m_\chi^2) \sqrt{1 - \frac{m_e^2}{m_\chi^2}}$$

=> relic density requirement points to

$$\sigma_e = \frac{1}{\pi} G_{\chi e}^2 \mu_{\chi, e}^2 \rightarrow (8-9) \times 10^{-35} \text{ cm}^2 \times \frac{2\mu_{\chi, e}^2}{(2m_\chi^2 + m_e^2)v_e}$$

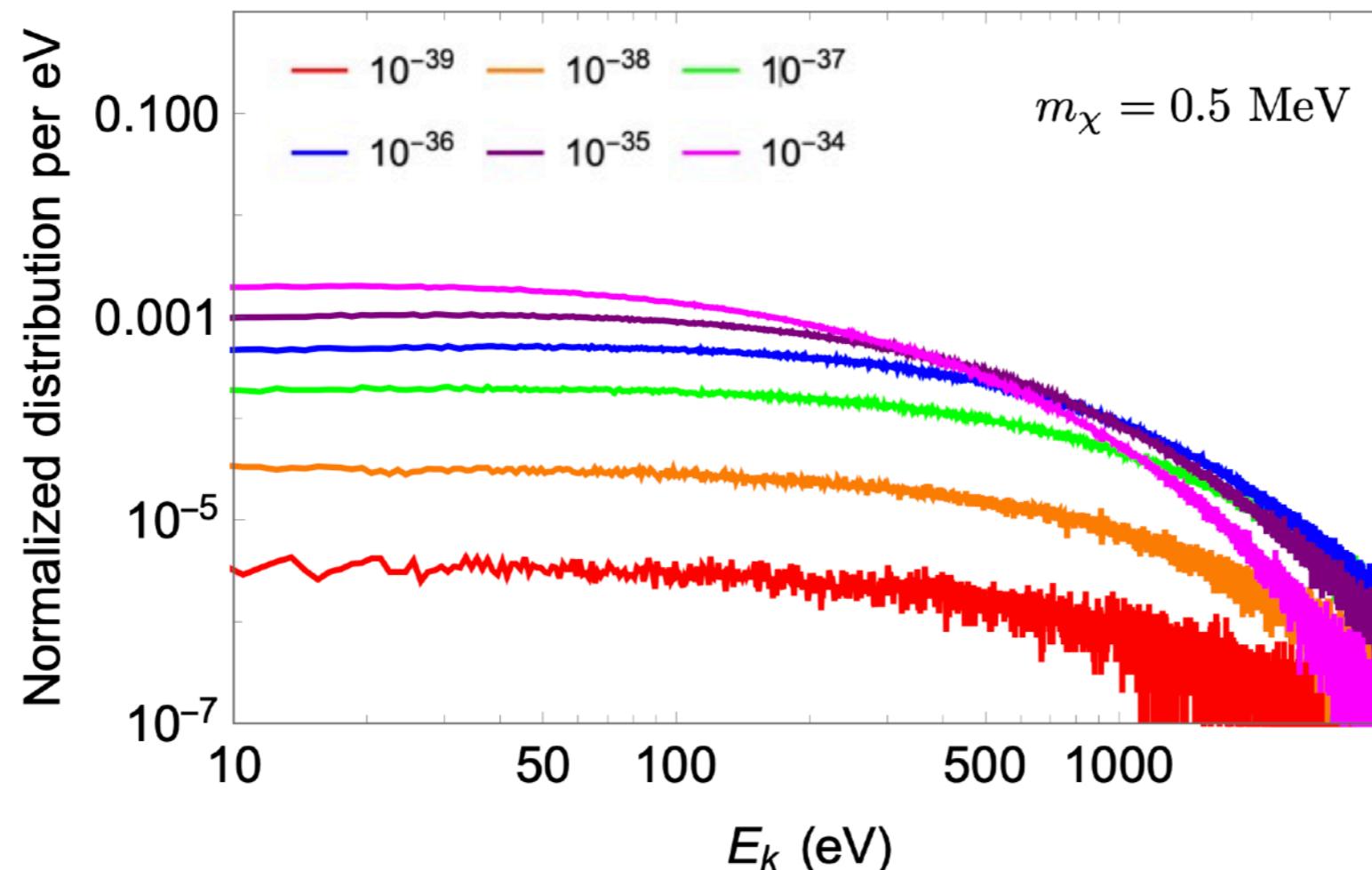
The sun as particle accelerator

Validating the Monte Carlo simulation in the single scattering limit



The sun as particle accelerator

Quantitative results for simulated fluxes



Spectra become softer for increasing cross section => reflection at larger radii

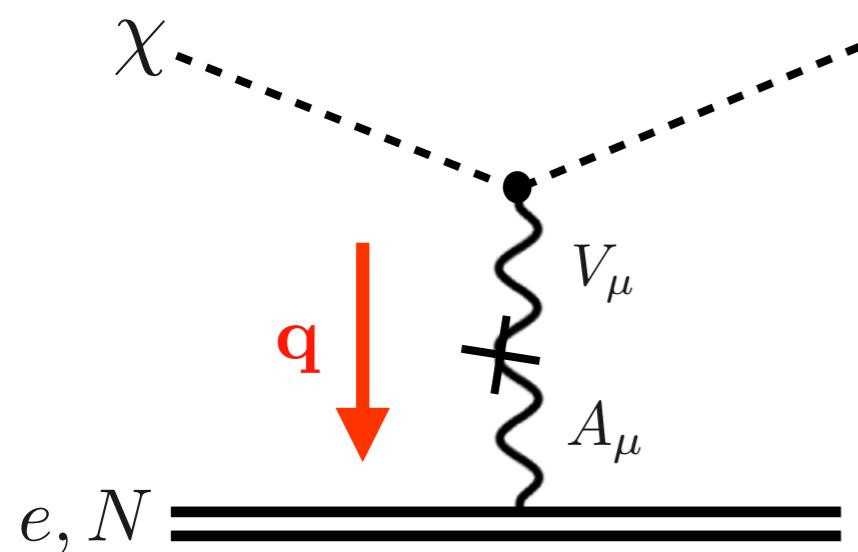
Spectra hardest for when the DM mass equals the electron mass

DM Direct Detection

Simple example

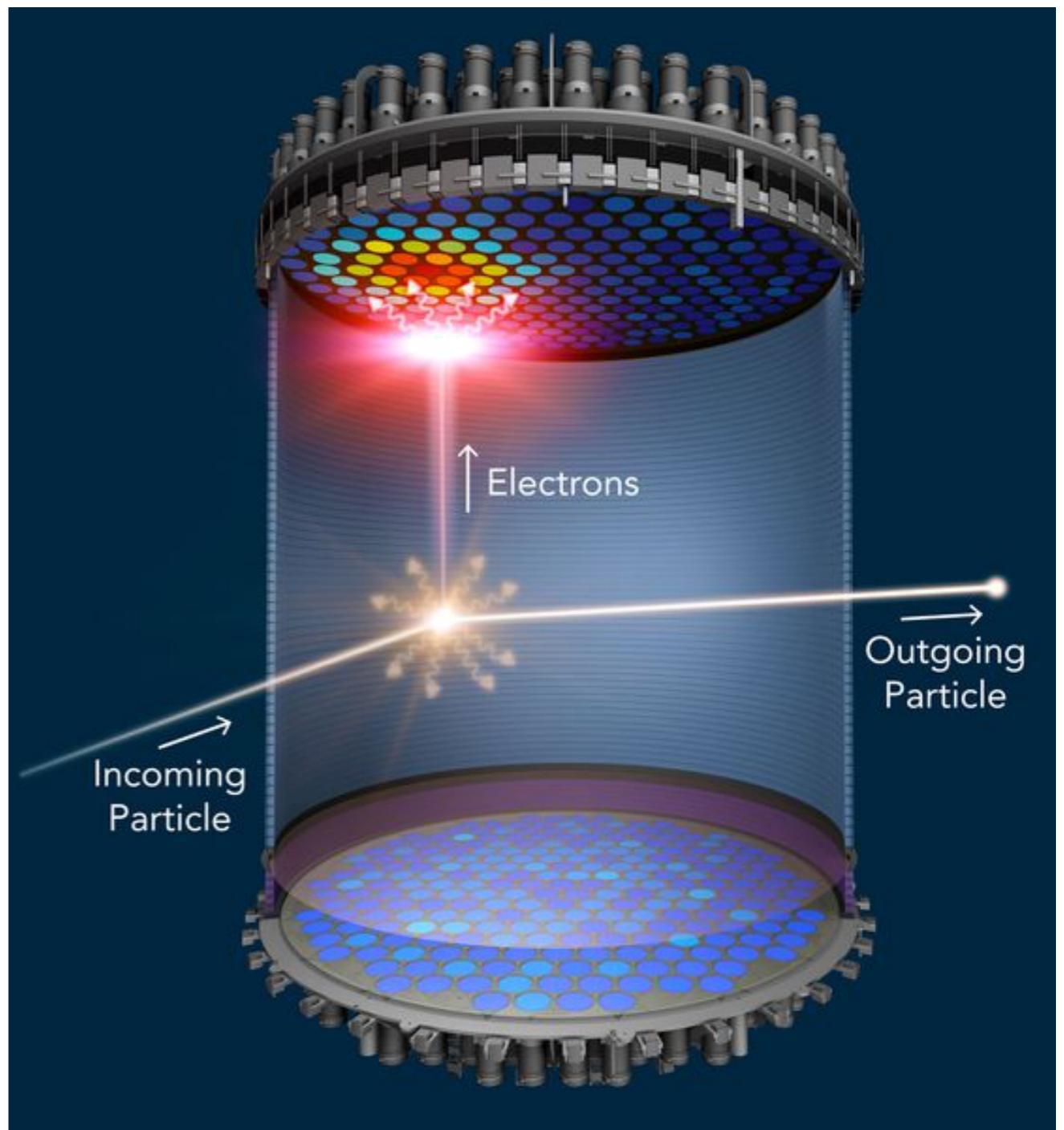
χ is the dark matter

V_μ dark photon is the “mediator”



$m_V \gg |q|$ point-like interaction

$m_V \ll |q|$ “millicharged DM”



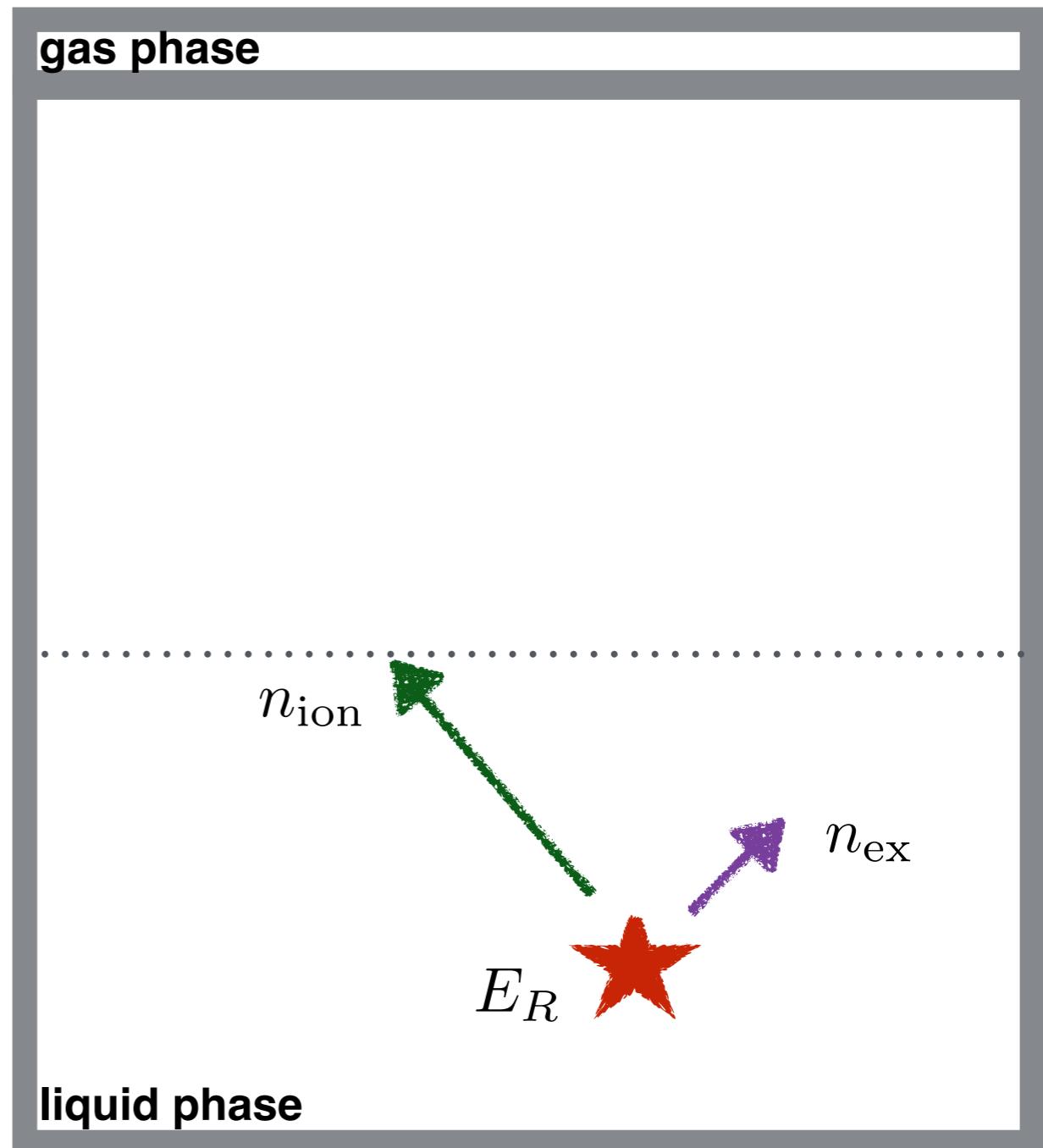
general detection principles:
ionization, scintillation, heat, ...

How does an “electron recoil” signal look in a LXe detector?

$$N_Q = \frac{E_R}{W} = n_{\text{ion}} + n_{\text{ex}}$$

$$W \simeq 13.7 \text{ eV} \quad n_{\text{ex}}/n_{\text{ion}} = \text{few \%}$$

Given energy deposition E_R , a number of quanta N_Q is produced, distributed in electron-ion pairs and excited atoms n_{ex}

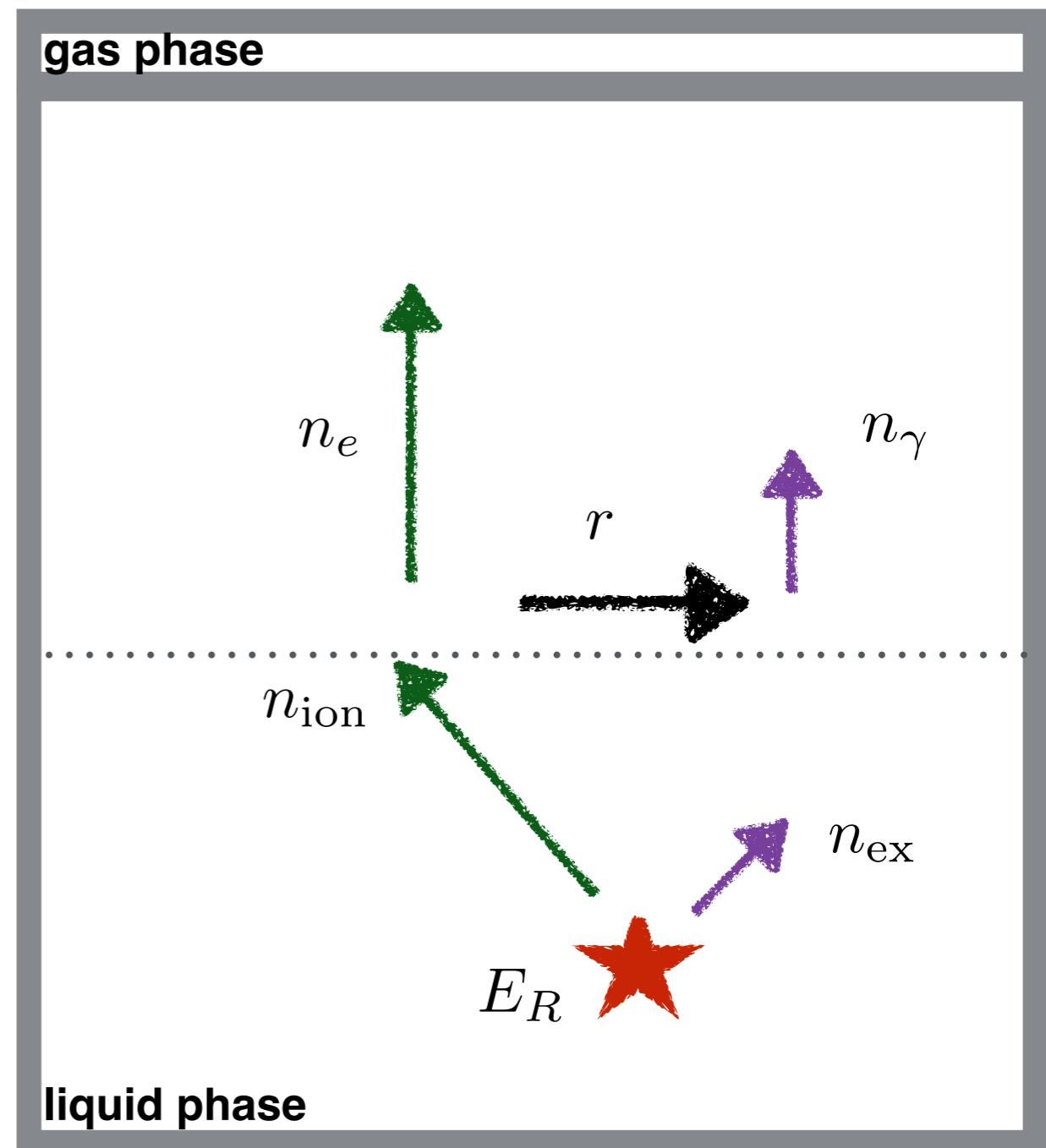


How does an “electron recoil” signal look in a LXe detector?

$$N_Q = \frac{E_R}{W} = n_{\text{ion}} + n_{\text{ex}}$$
$$= n_\gamma + n_e$$

$$n_e = n_{\text{ion}}(1 - r), \quad n_\gamma = n_{\text{ion}}r + n_{\text{ex}}$$

Observable: de-excitation photons from initial and recombined excitons n_γ and electrons that escape recombination n_e



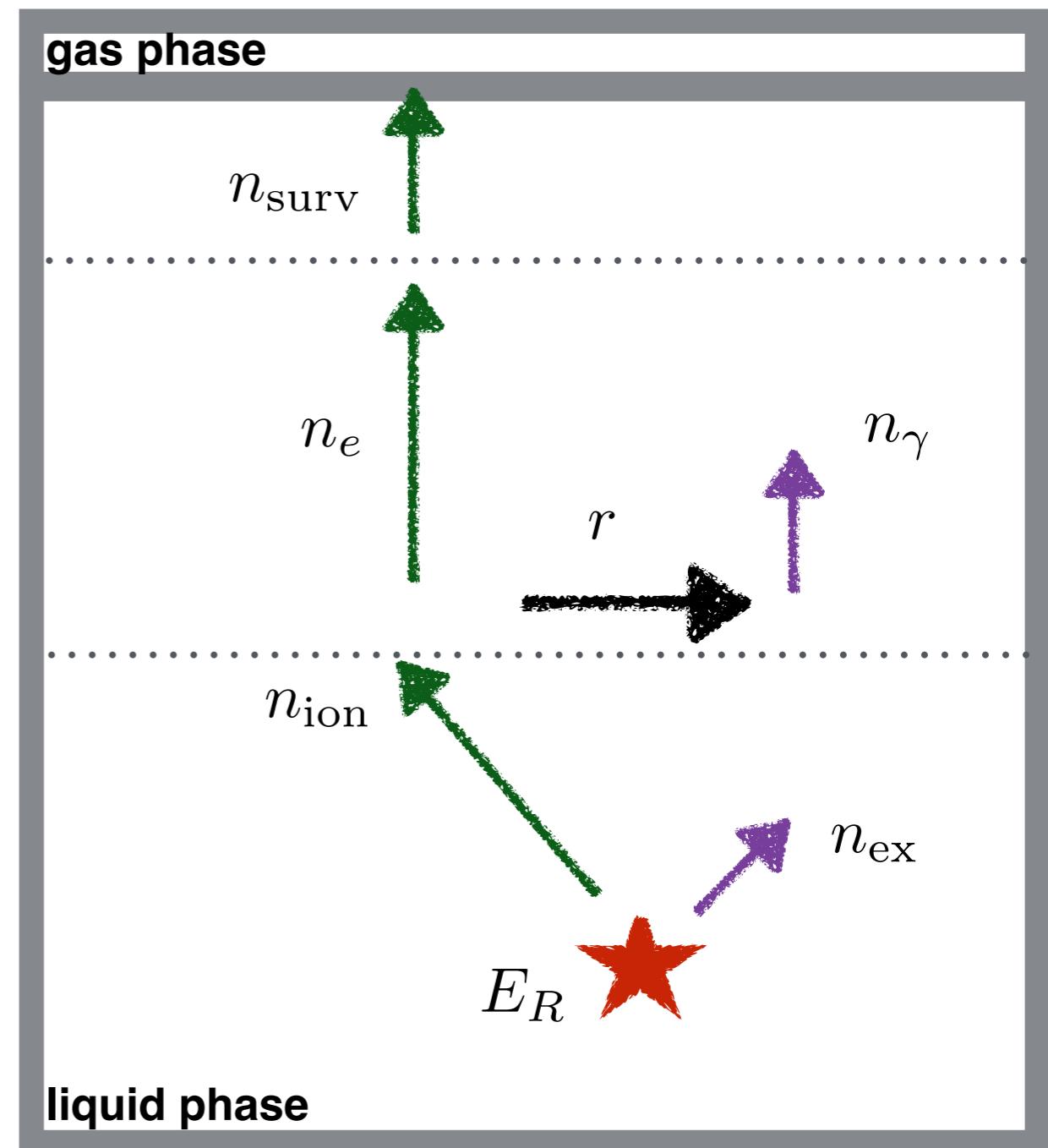
How does an “electron recoil” signal look in a LXe detector?

$$p_{\text{surv}} \simeq \exp\left(-\frac{\Delta z}{\tau v_d}\right)$$

$$v_d \simeq 1.7 \text{ mm}/\mu\text{s} \quad \tau > 1 \text{ s}$$

Electrons are drifted in the electric field towards the liquid-gas interface; depending where they are created, attenuation occurs

$$p_{\text{surv}} \sim 0.6 - 0.9$$

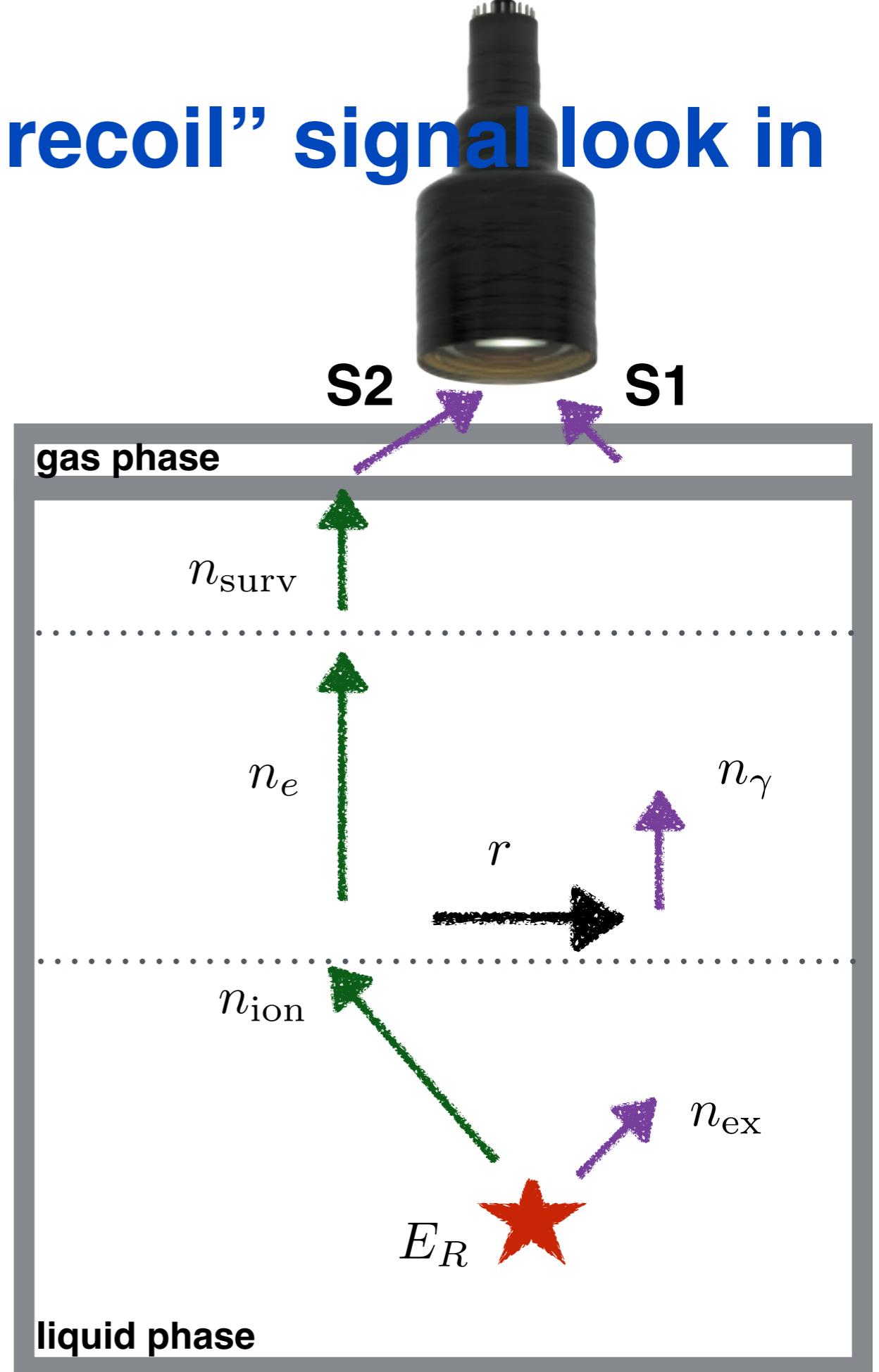


How does an “electron recoil” signal look in a LXe detector?

$$\begin{aligned}N_Q &= n_{\text{ion}} + n_{\text{ex}} \\&= n_\gamma + n_e \\&= \frac{S1}{g_1} + \frac{S2}{g_2}\end{aligned}$$

$$g_1 \simeq 0.1, \quad g_2 \simeq 10 - 50$$

An electron reaching the liquid-gas interface creates about $O(10)$ PE (S2); it takes on average 10 scintillation photons to collect 1 PE (S1)



How does an “electron recoil” signal look in a LXe detector?

detector specific

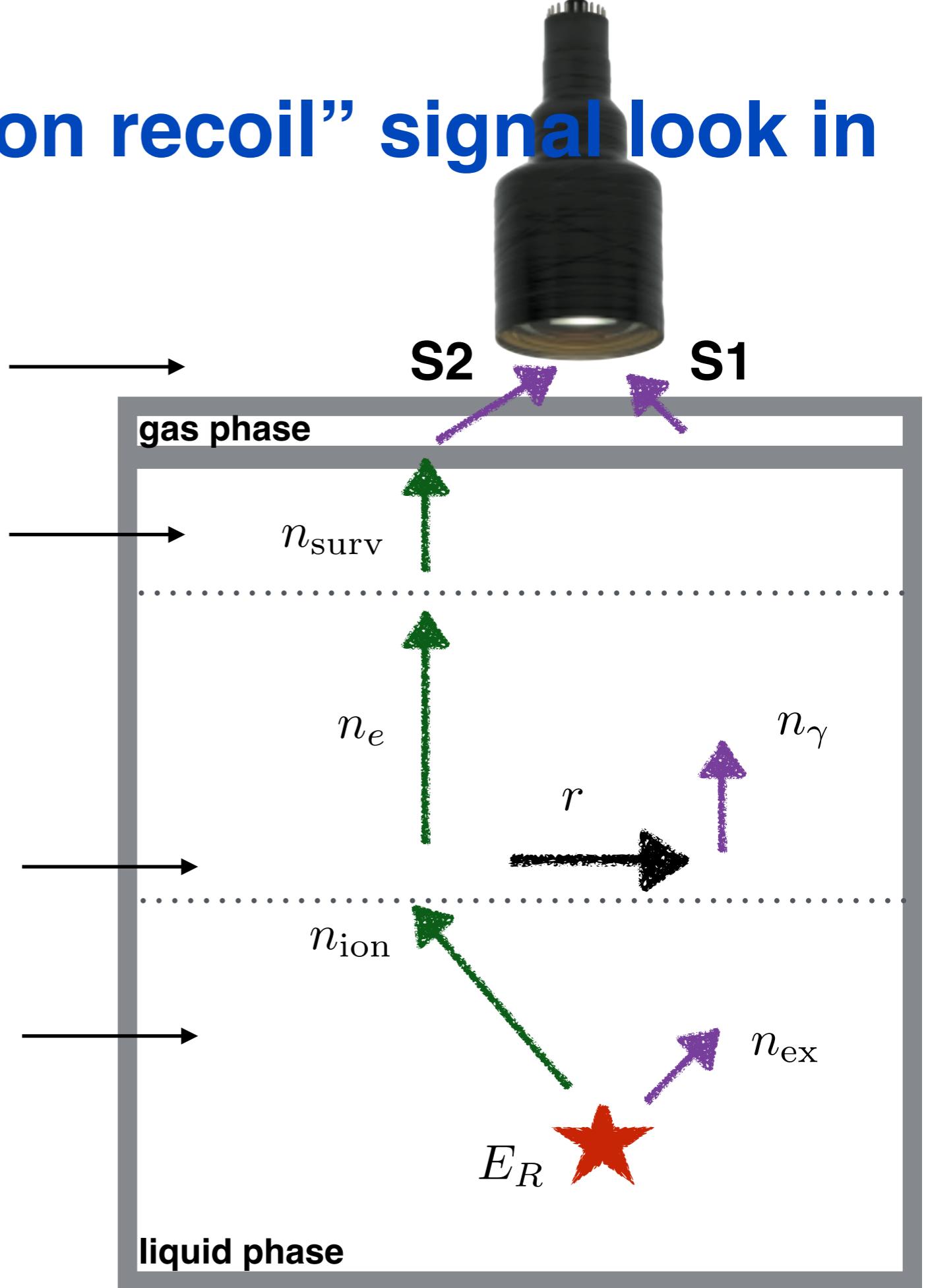
{ fluctuates

{ fluctuates

LXE universal
(given E-field)

{ fluctuates

=> model $\text{PDF}(\text{S1}, \text{S2} | \text{E}_{\text{dep}})$

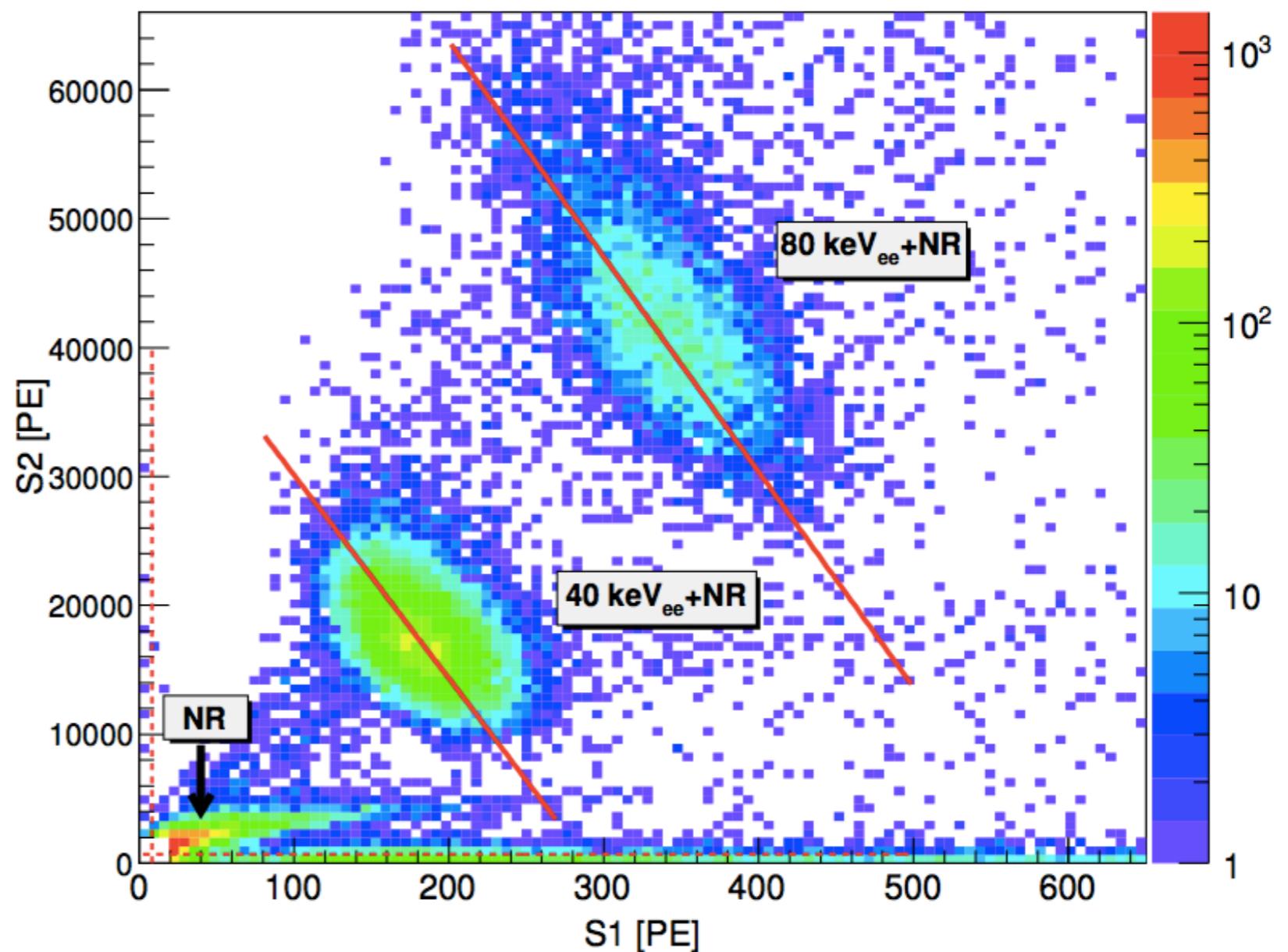


How does an “electron recoil” signal look in a LXe detector?

e.g. PandaX

$$\begin{aligned}N_Q &= n_{\text{ion}} + n_{\text{ex}} \\&= n_\gamma + n_e \\&= \frac{S1}{g_1} + \frac{S2}{g_2}\end{aligned}$$

note the anti-correlation between S1 and S2



How does an “electron recoil” signal look in a LXe detector?

S2-only search

$$\text{PDF} \quad P(S2|E_R) = \sum_{n_e^{\text{surv}}} \sum_{n_e} P(S2|n_e^{\text{surv}}) P(n_e^{\text{surv}}|n_e) P(n_e|\langle n_e \rangle)$$

For example: $P(n_{e,\gamma}|\langle n_{e,\gamma} \rangle) = \text{binom}(n_{e,\gamma}|N_Q, f_{e,\gamma})$

$$\langle n_e \rangle = E_{\text{dep.}} Q_y \quad \text{with charge yield } Q_y \text{ measured or modelled}$$

$$P(n_e^{\text{surv}}|n_e) \simeq \langle p_{\text{surv}} \rangle = 0.8$$

$$P(S2|n_e^{\text{surv}}) = \text{gauss}(S2|g_2 n_e^{\text{surv}}, \sigma_{S2})$$

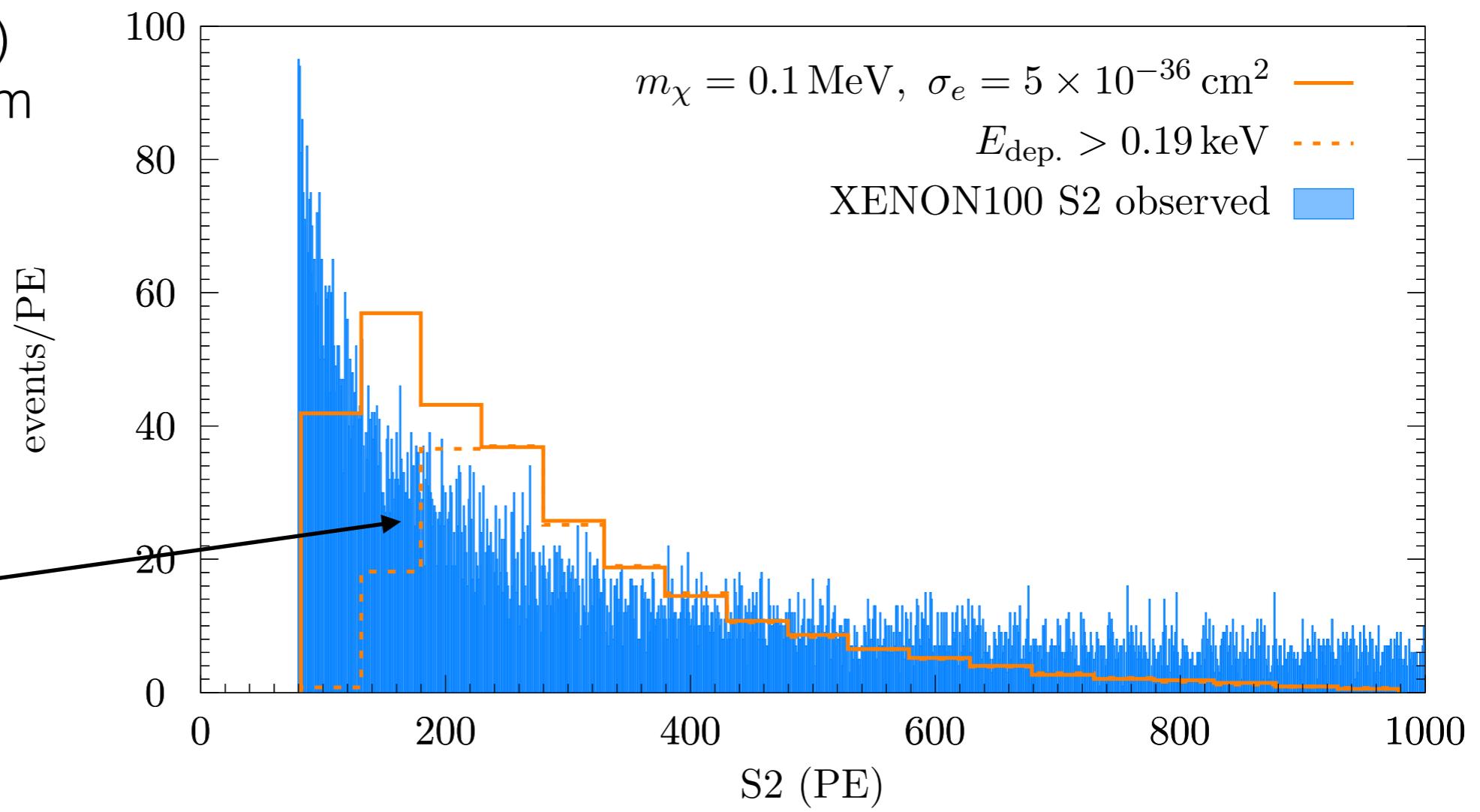
Experimental rate: $\frac{dR}{dS2} = \varepsilon(S2) \int dE_R \quad P(S2|E_R) \frac{dR}{dE_R}$

S2 only spectrum

Example XENON100

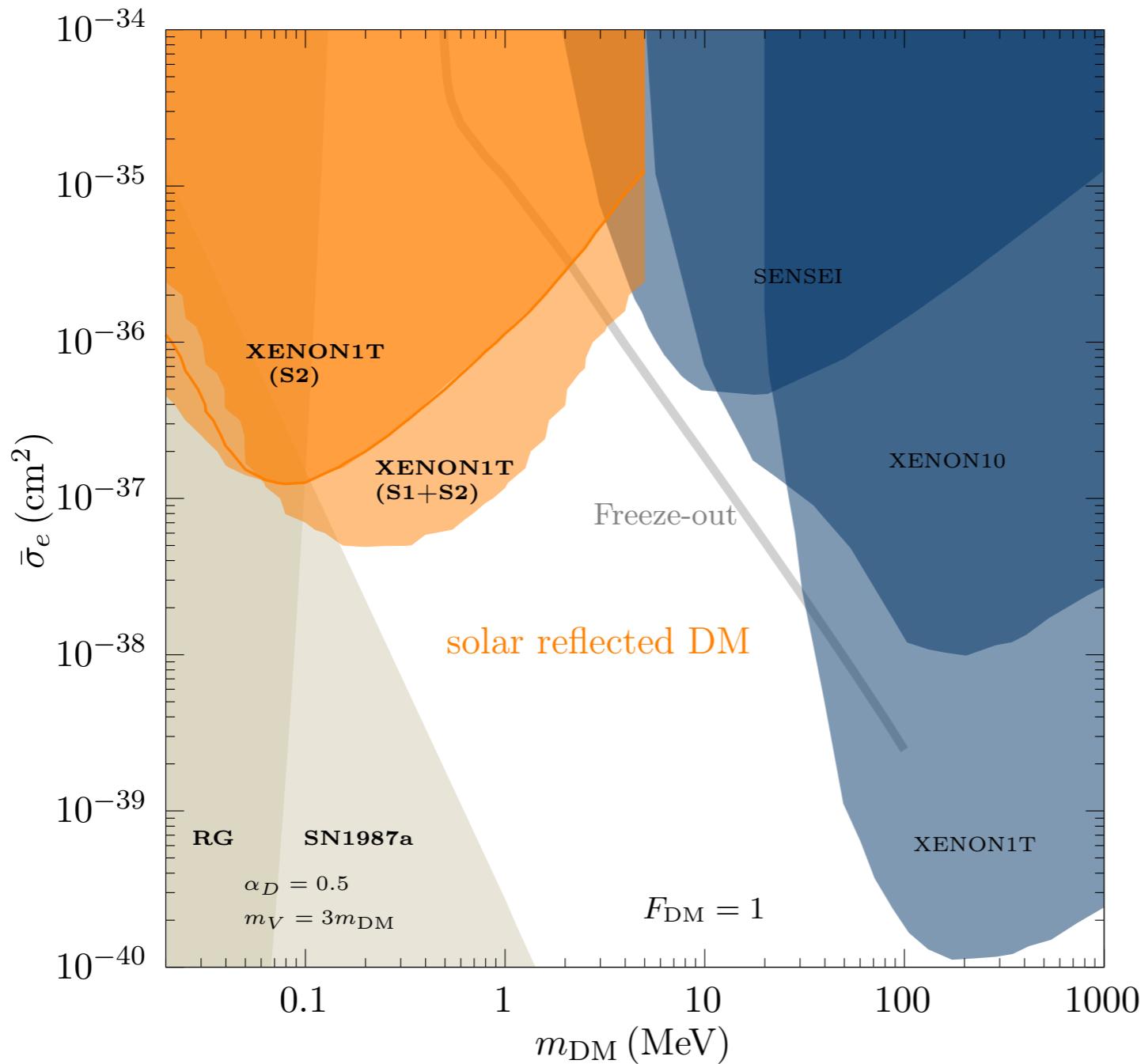
XENON100 (2016)
ionization spectrum

for expt. limits:
demand minimum
deposited energy
of 0.19 keV
=> data driven
approach



Solar Reflection of DM

Contact interactions with electrons



An, Pospelov, JP, Ritz PRL 2018

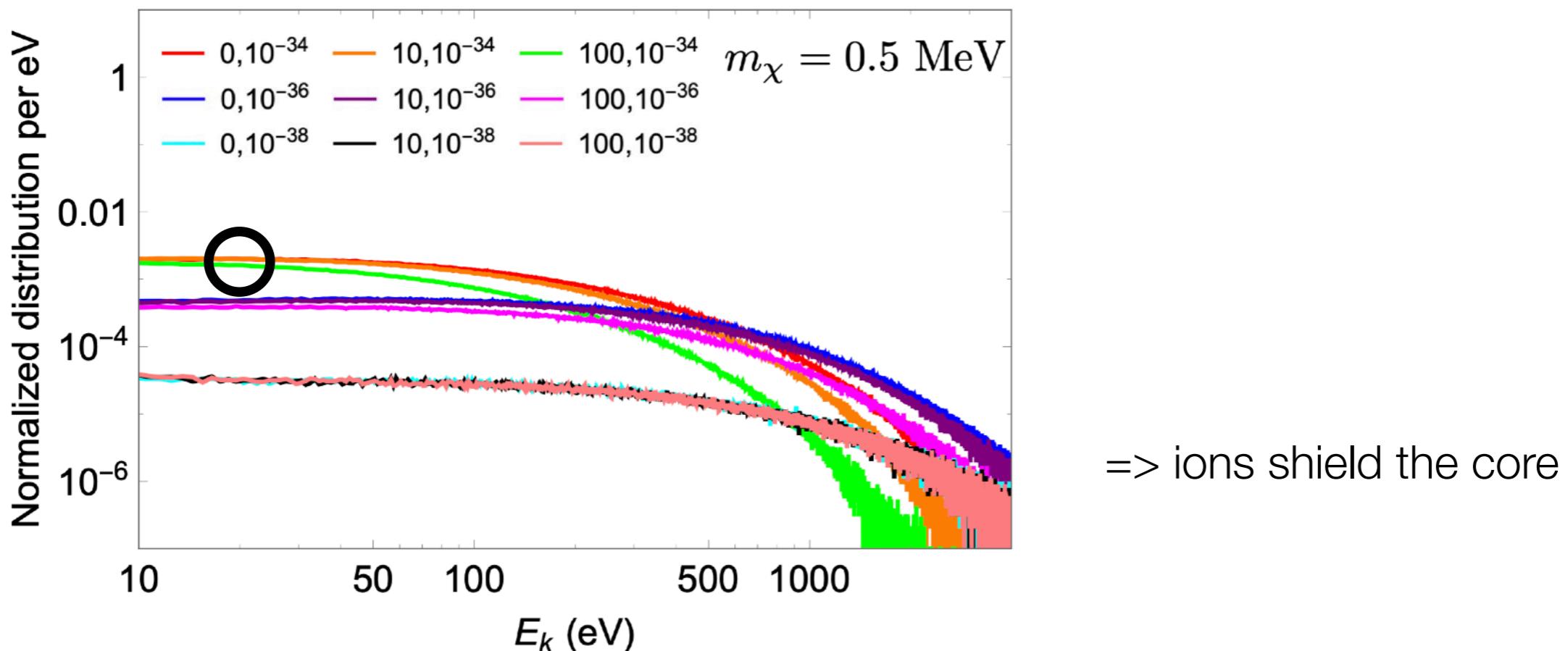
An, Nie, Pospelov, JP, Ritz PRD 2021

Solar Reflection of DM

Role of ions?

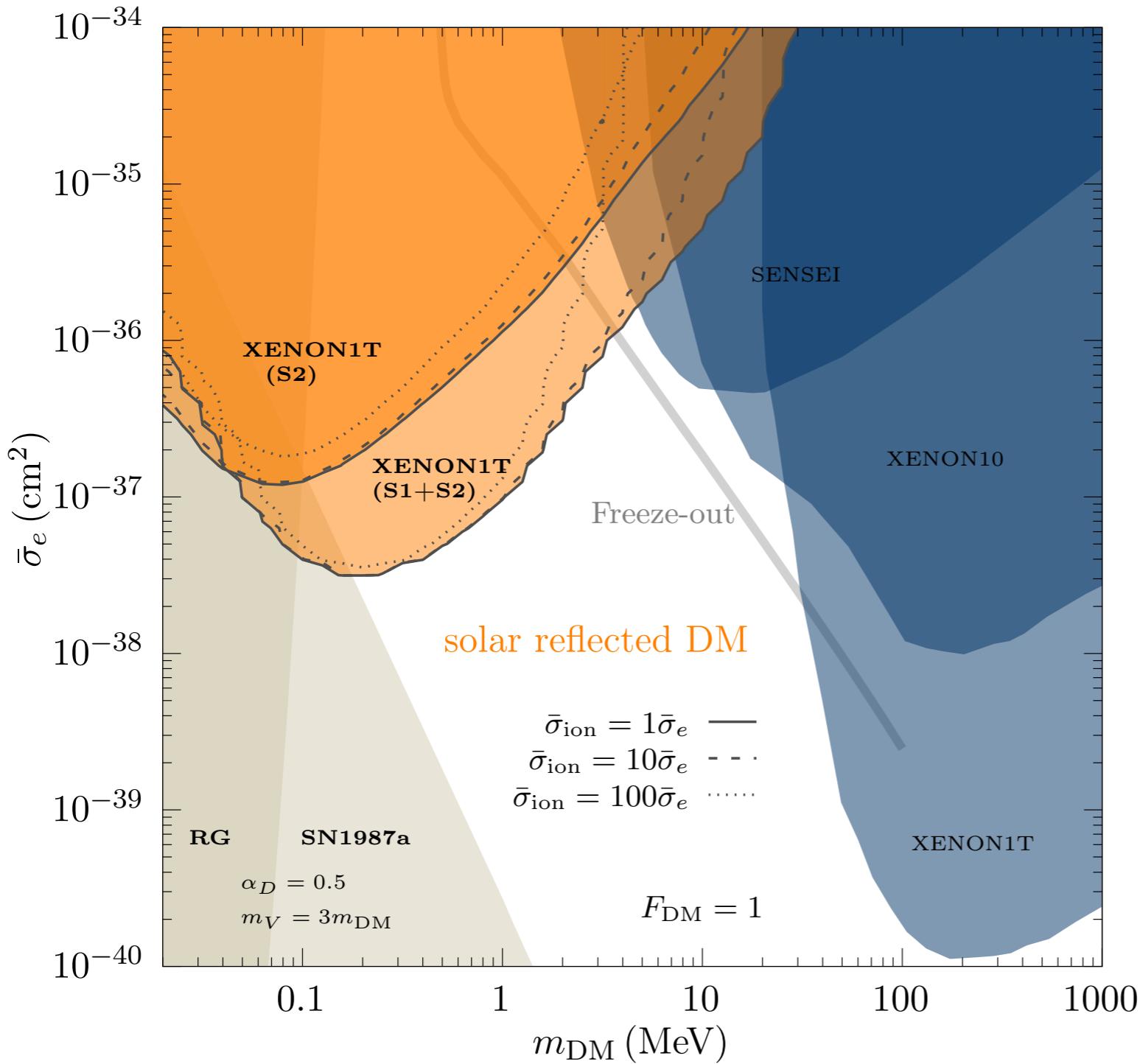
$m_\chi \ll m_{\text{nucleus}}$ collisions with ions only change direction but not energy

2 options: either ions shield the hot solar core from DM, or they may turn around DM that has already entered and increase chance of further upscattering



Solar Reflection of DM

Role of ions



An, Pospelov, JP, Ritz PRL 2018

An, Nie, Pospelov, JP, Ritz PRD 2021

Solar Reflection of DM

Millicharged DM

Difficult/expensive to treat because of long-range interactions

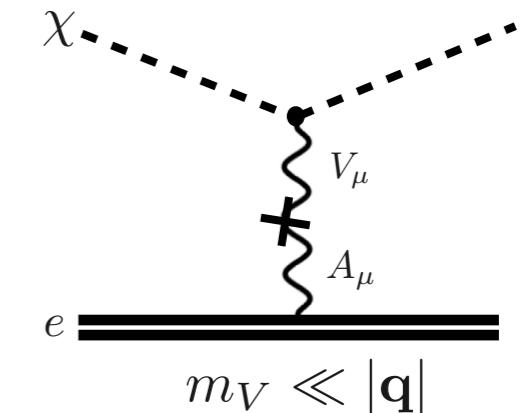
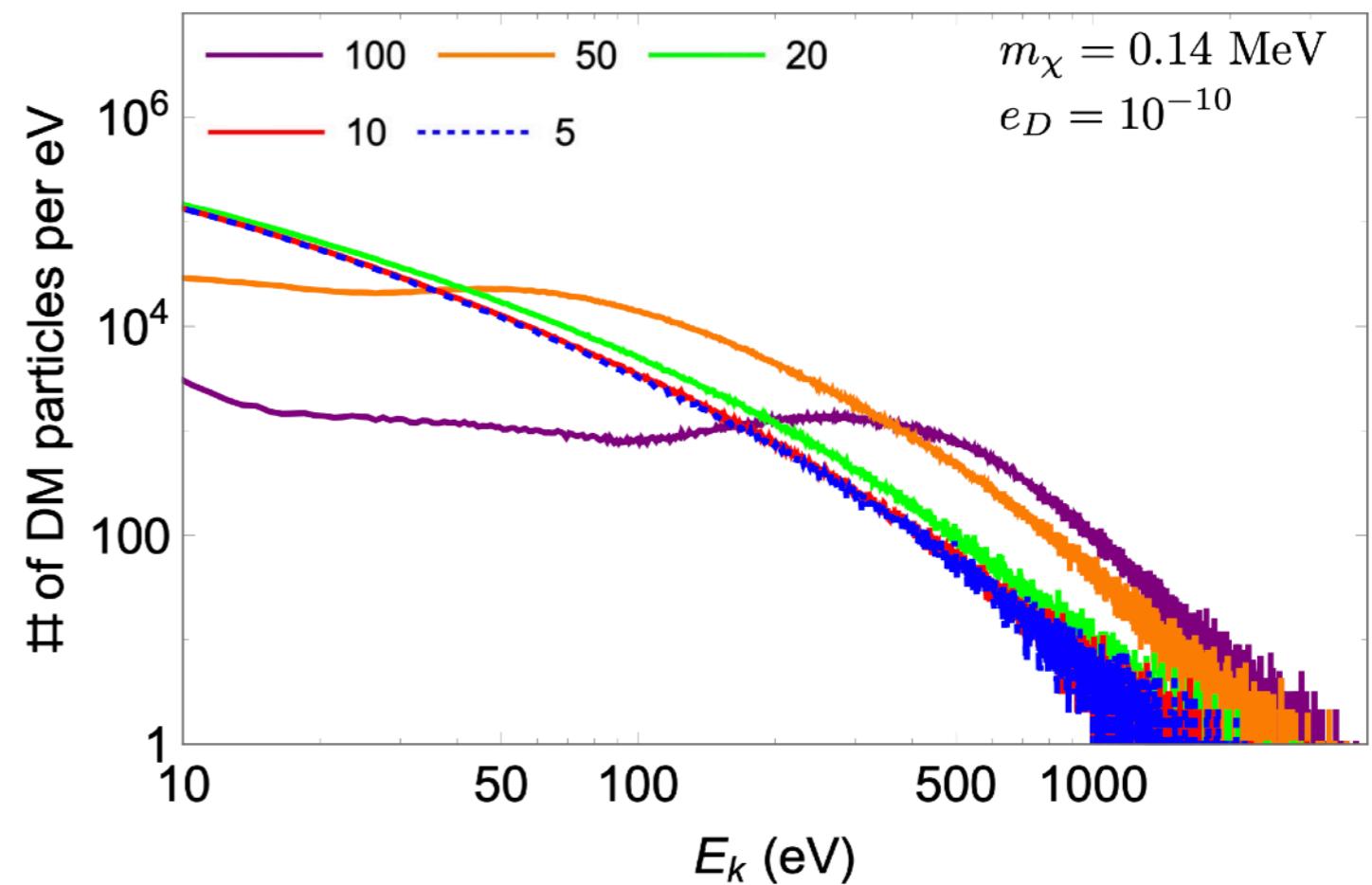
=> forward scattering biased; eventually “Debye-screened”

2 effects:

“hard” large-angle scatterings
accelerate DM;

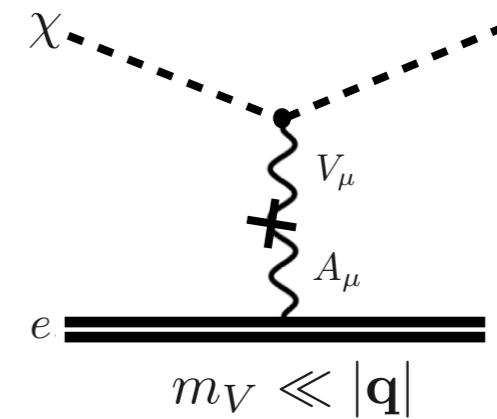
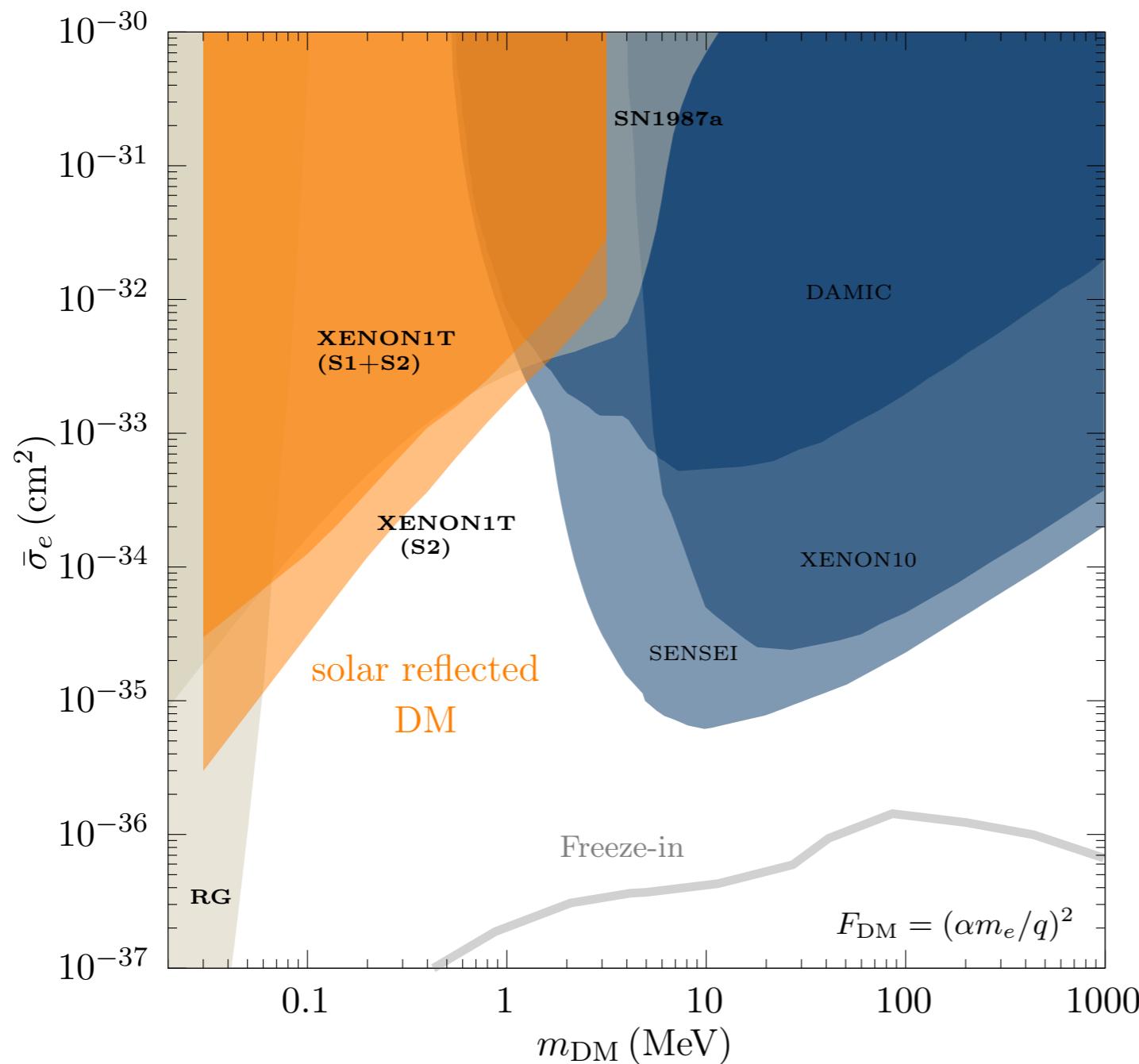
“soft” small-angle scatterings
effectively friction/viscosity

$$x_q = \frac{q}{\sqrt{m_e T}} = \frac{\Delta v_1 m_\chi}{\sqrt{m_e T}} > \frac{\zeta v_1 m_\chi}{\sqrt{m_e T}}$$



Solar Reflection of DM

Millicharged DM

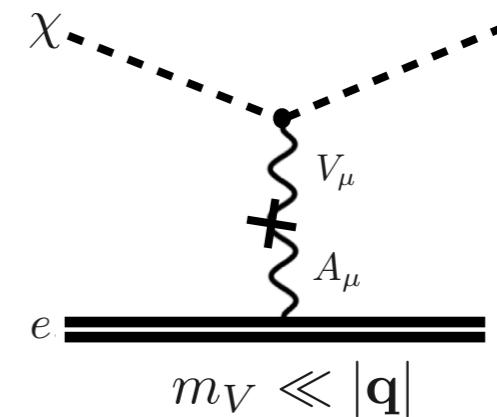
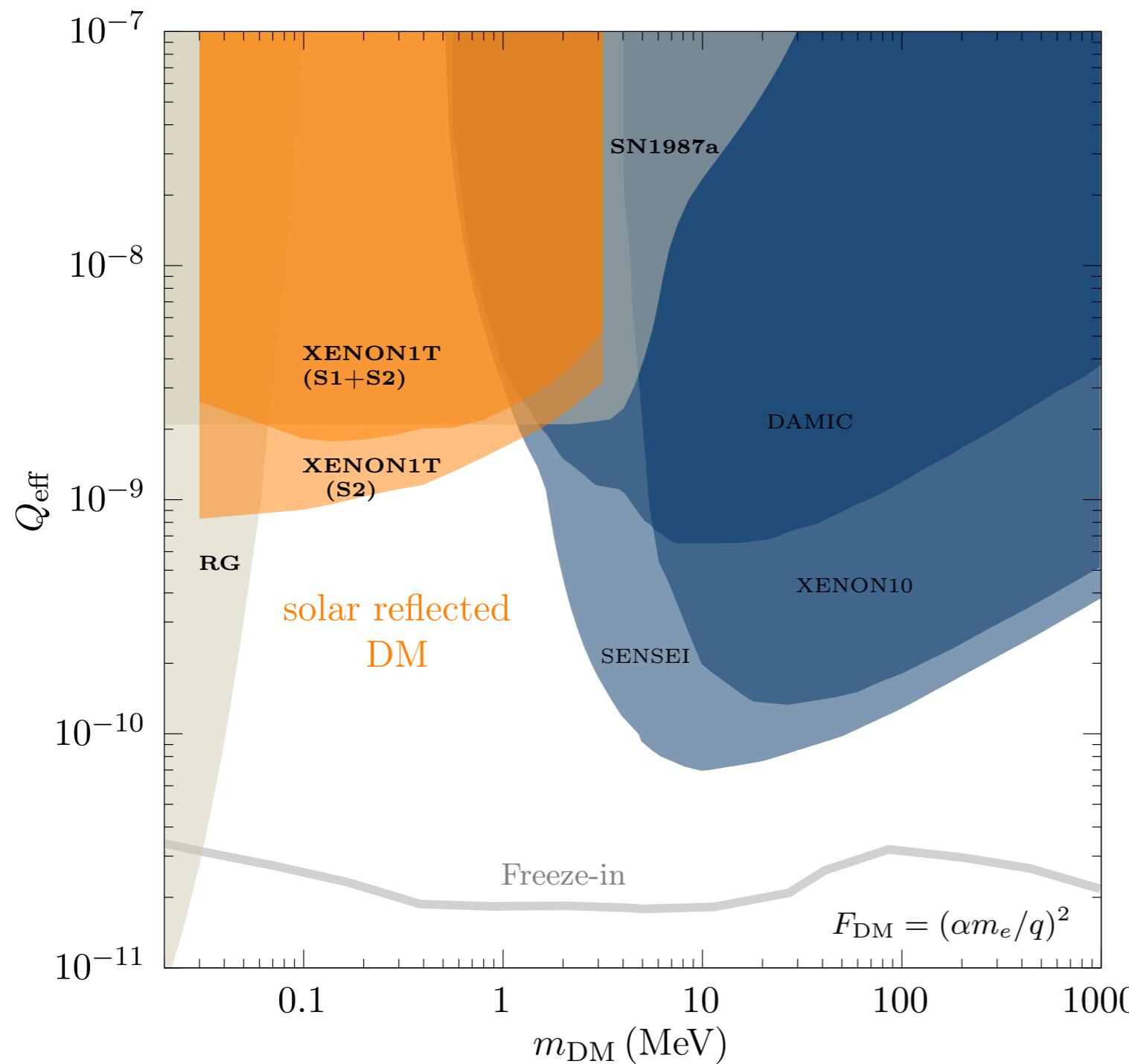


$$Q_{\text{eff}} = e_D \kappa / e$$

An, Nie, Pospelov, JP, Ritz PRD 2021

Solar Reflection of DM

Millicharged DM

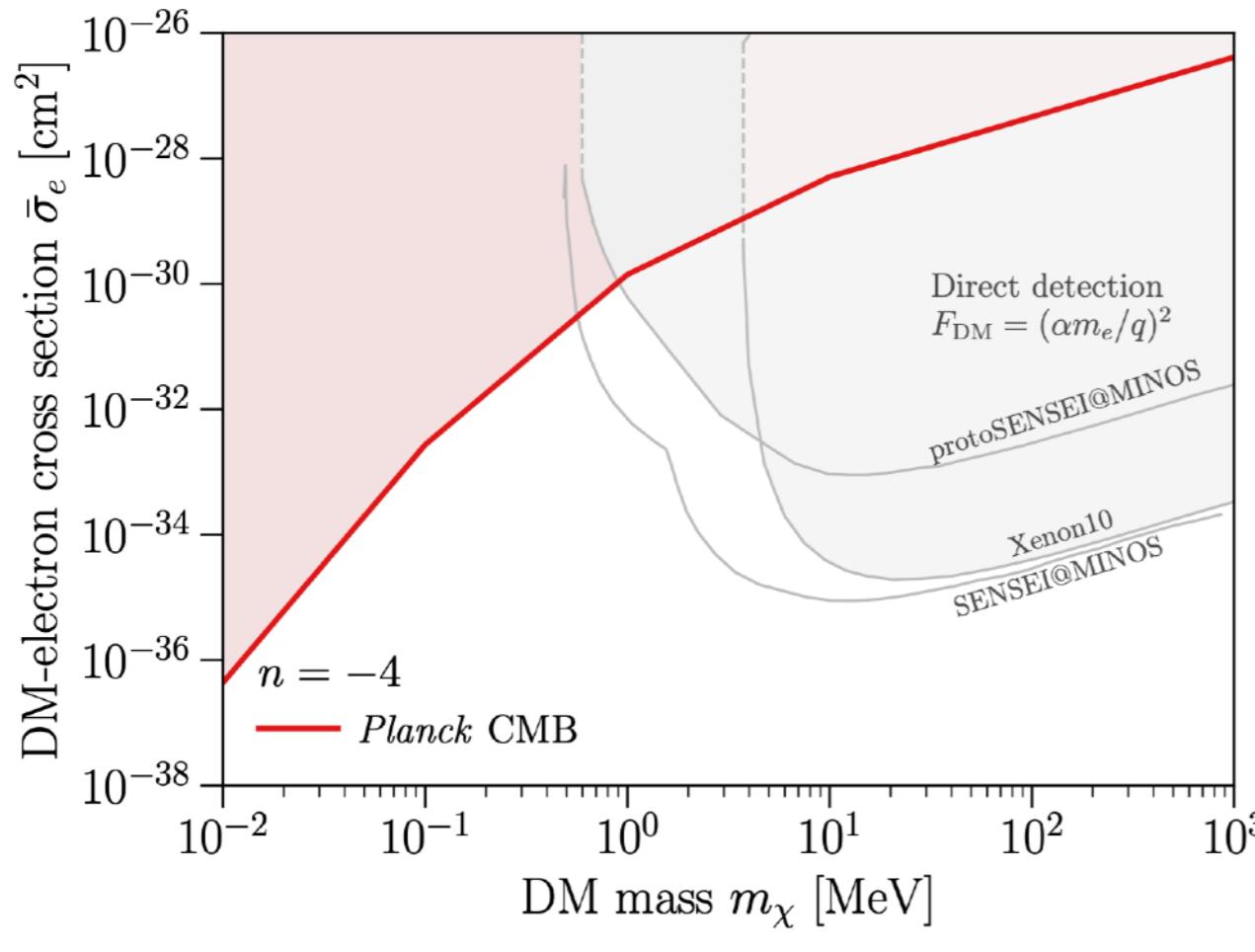


$$Q_{\text{eff}} = e_D \kappa / e$$

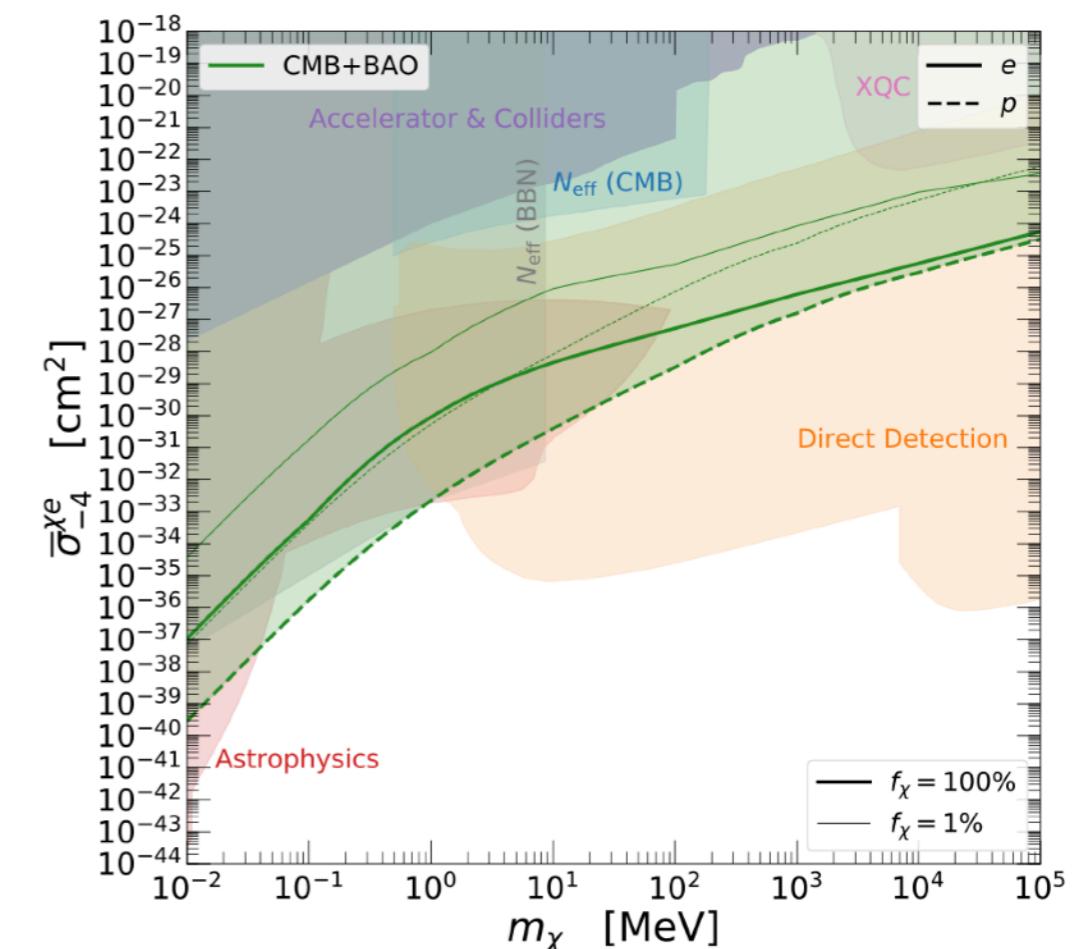
An, Nie, Pospelov, JP, Ritz PRD 2021

Cosmological limits

Similar sensitivity from cosmology



Nguyen et al arXiv:2107.12380



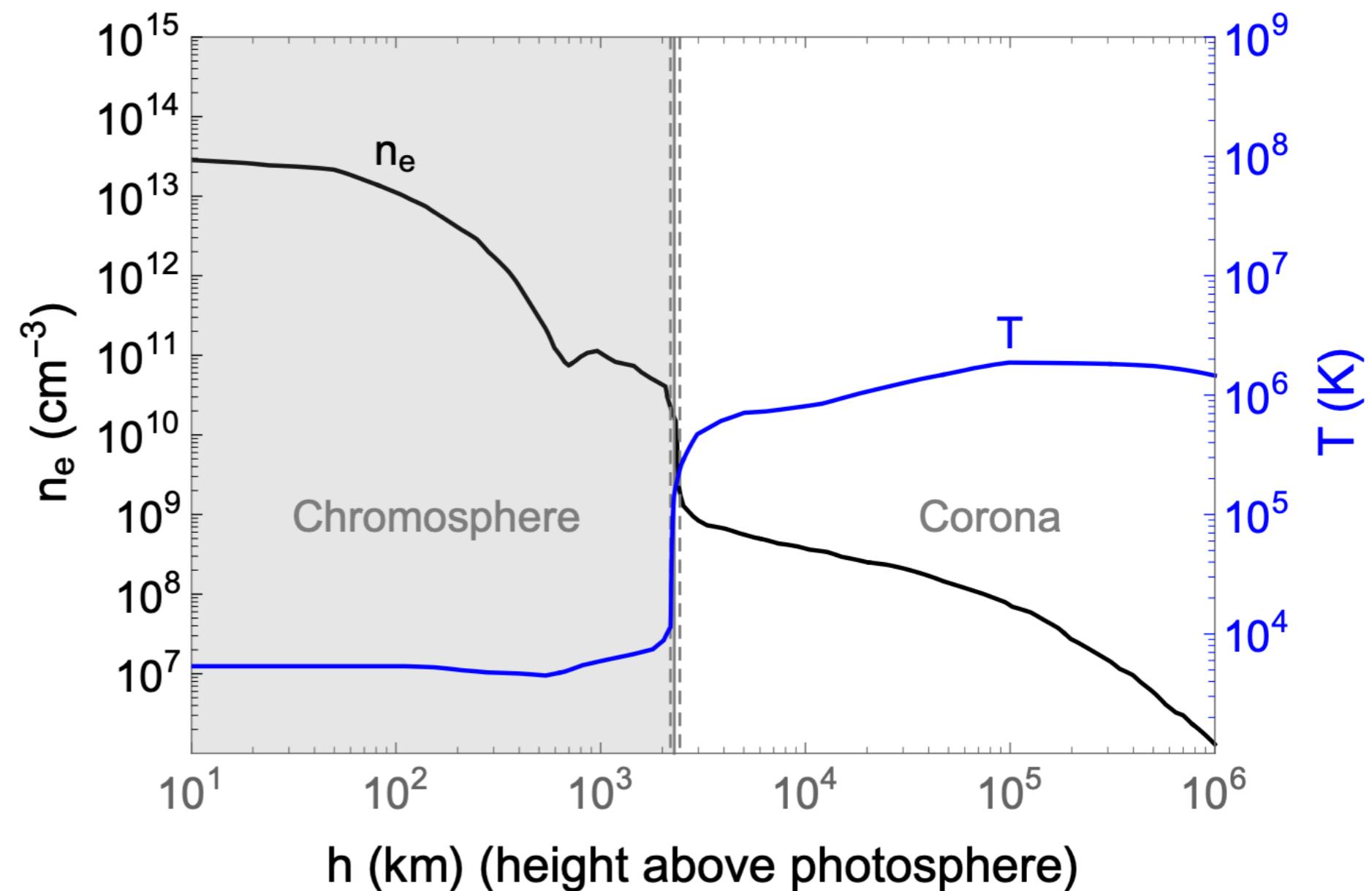
Buen-Abad et al arXiv:2107.12377v1

Reflection of DM from the Corona

Sensitivity to large couplings? (where particles don't enter the sun)

Solar corona dilute,
but hot!

Very small but energetic
enough flux possible



Reflection of DM from the Corona

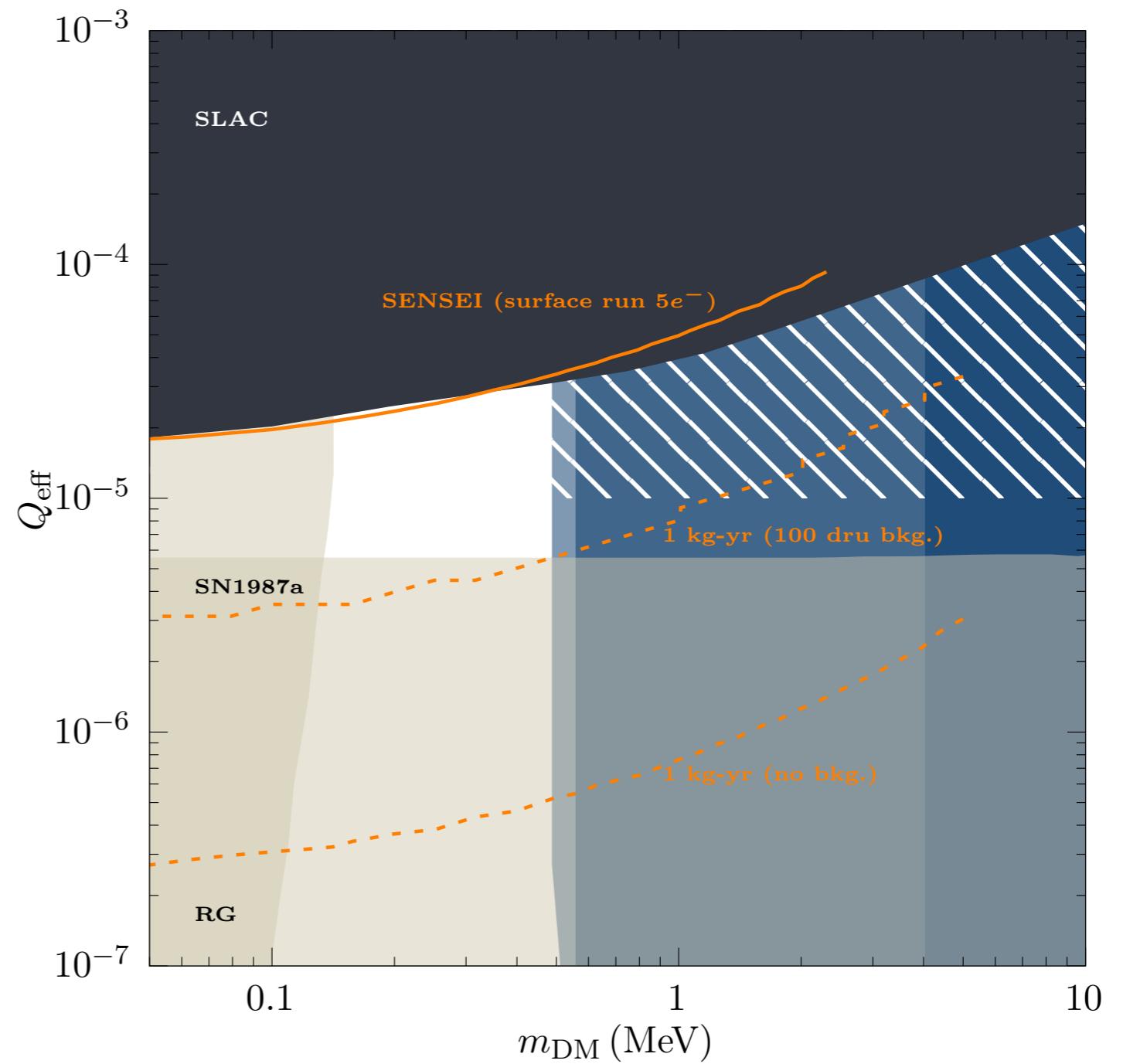
Sensitivity to large couplings? (where particles don't enter the sun)

Solar corona dilute,
but hot!

Very small but energetic
enough flux possible

particles will not reach deep
underground => take surface
runs

(we are neglecting
complications here; rough
estimate)



Summary

- the number of possibilities for particle DM appears daunting (no insight on mass); however there are very well motivated cases that can serve as “prototype models” coupled through portal interactions
- direct detection aims at registering DM-atom interaction; lowering threshold, one gains exponentially
- we may harvest irreducible signal components (Bremsstrahlung, Migdal electrons, solar reflected DM) to extend the physics reach of those experiments without extra cost
- solar reflection of MeV-mass DM with couplings to electrons
 - => O(10^{-4}) component to the DM flux at earth from solid angle
 - => extends the reach in the “electron-scattering” channel to MeV DM mass range with optimum sensitivity at electron mass direct detection
 - => cosmological bounds on DM-electron scattering for light mediator are similar but complementary; for sub-% fractional abundance, cosmological limits may disappear

Summary

