

WIMP dark matter as radiative neutrino mass messenger

Joaquim Palacio
Thursday 7th November, 2013

9th MultiDark Consolider Workshop Alcalá de Henares

M. Hirsch, R.A. Lineros, S. Morisi, J. Palacio, N. Rojas, J.W.F. Valle.
[10.1007/JHEP10\(2013\)149](https://arxiv.org/abs/10.1007/JHEP10(2013)149)



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DE VALÈNCIA



Outline

Dark Matter Evidences

(Radiative) Neutrino Mass Models

Dark Matter & Neutrino Masses

The Model

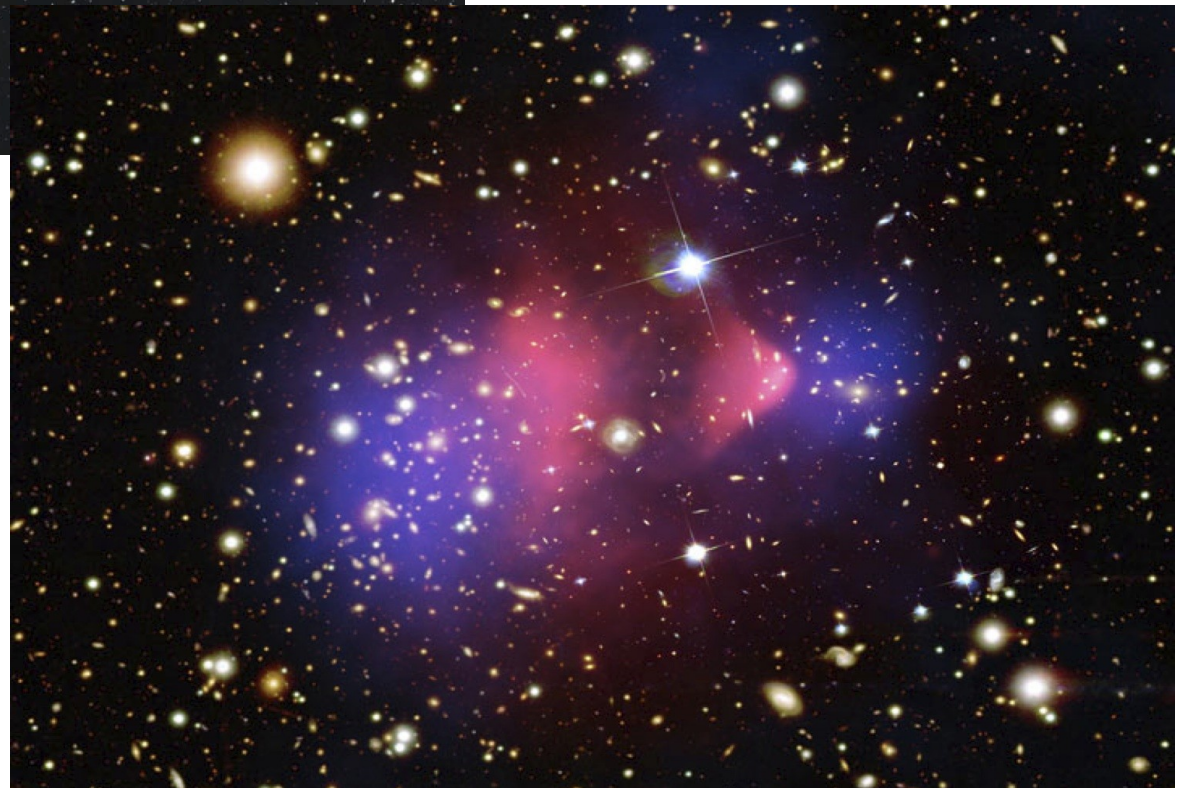
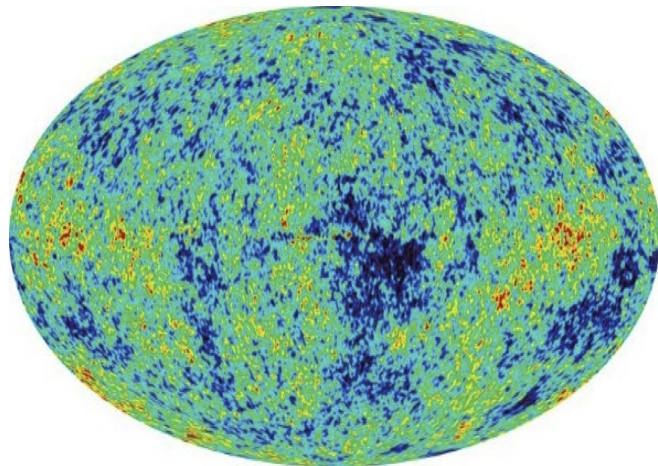
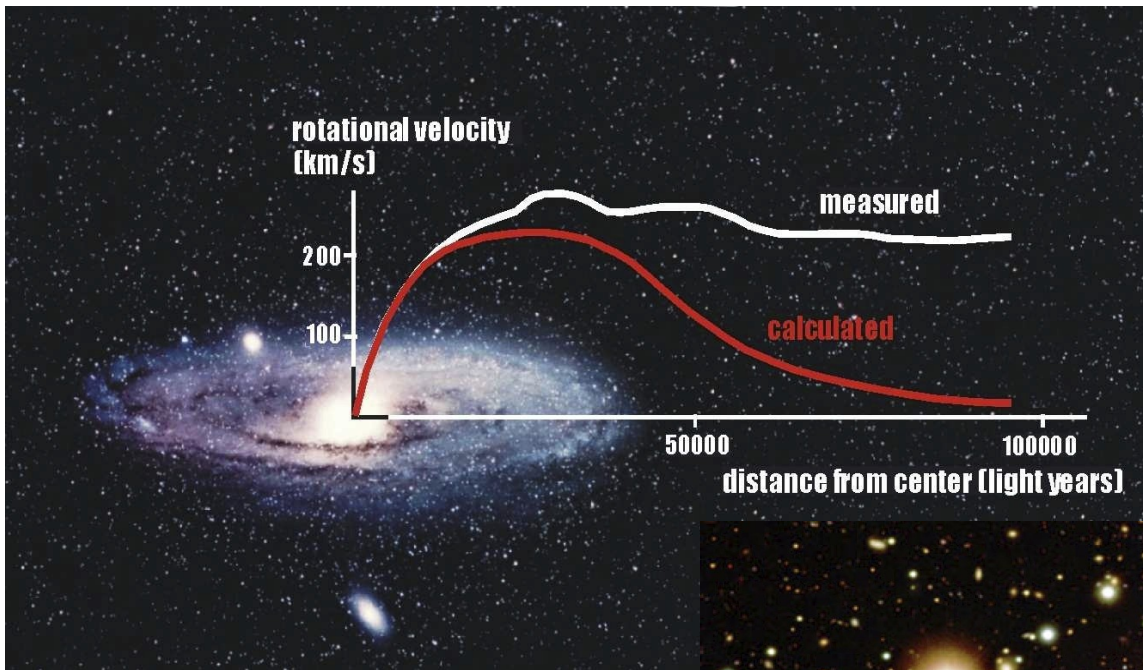
Particle content

Study of the Model

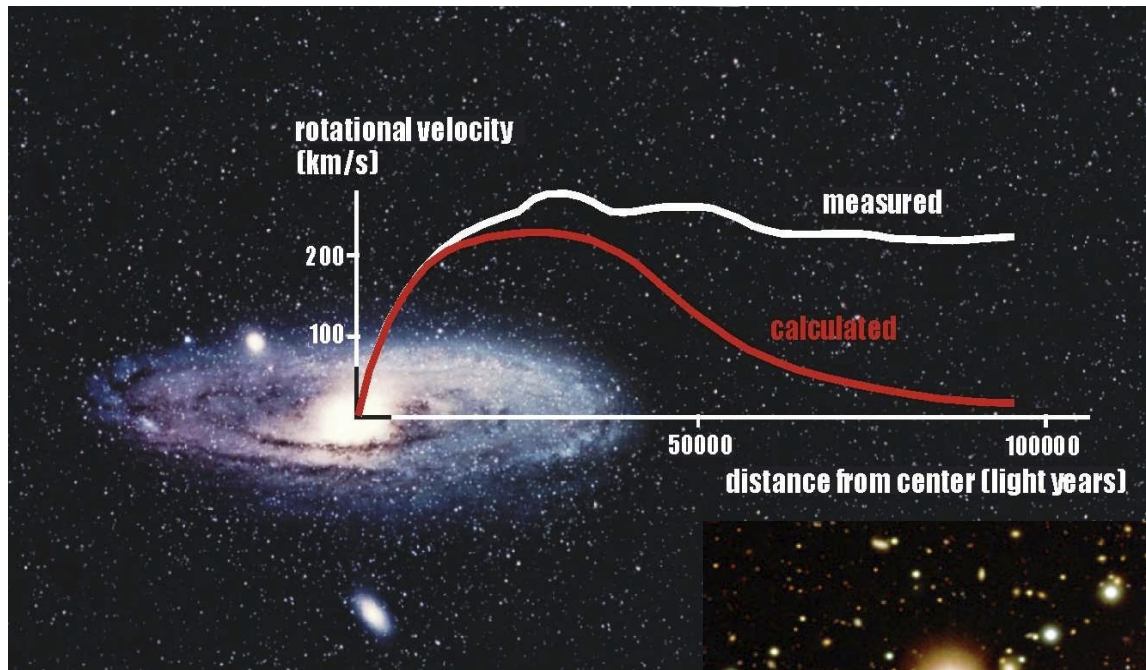
Detection prospects

Conclusions

DM evidences

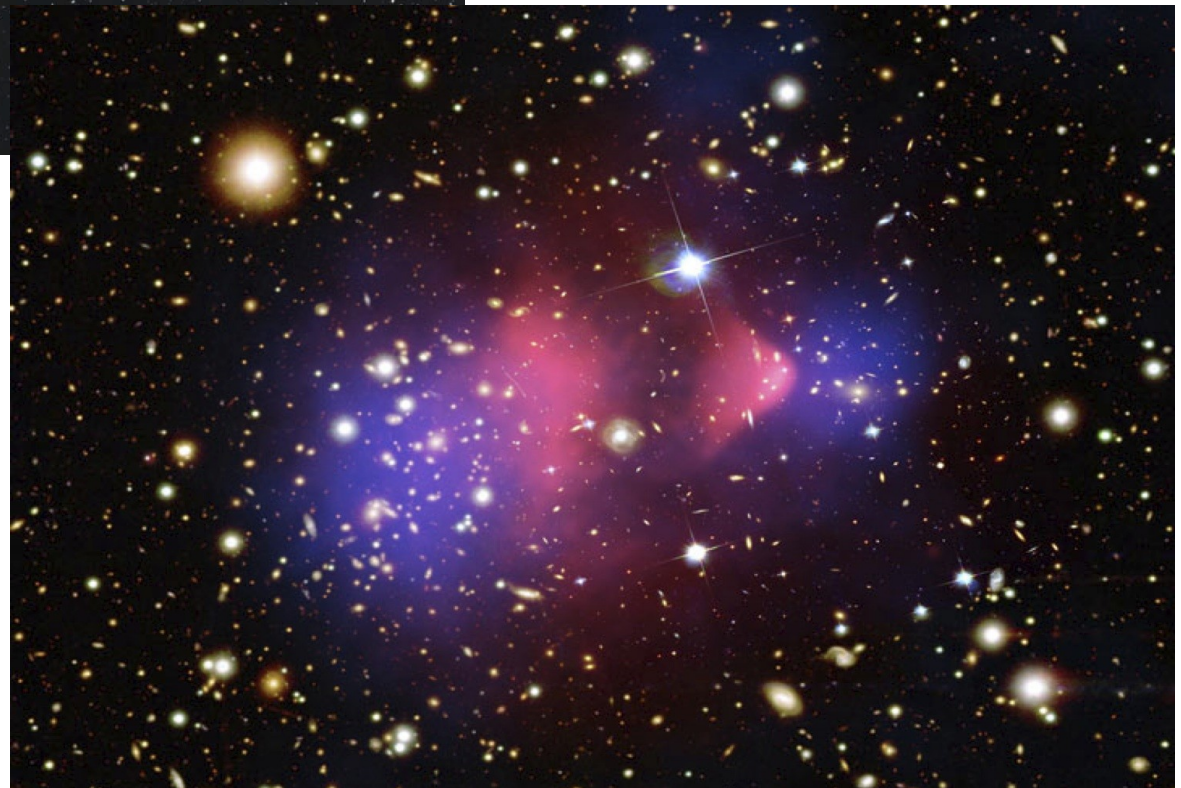
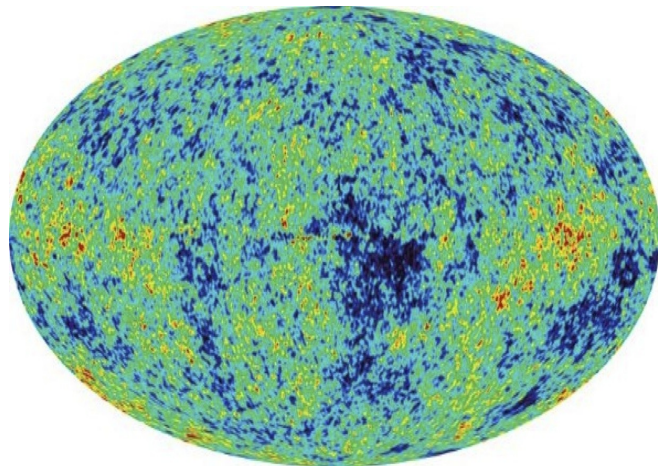


DM evidences



Need of extra matter

→ ¿WIMP?



Neutrino mass models

Neutrinos are massless in the SM

FORERO-TORTOLA-VALLE
Phys. Rev. D 86, 073012 (2012)

parameter	best fit	1σ range	2σ range	3σ range
$\Delta m_{21}^2 [10^{-5} \text{eV}^2]$	7.62	7.43–7.81	7.27–8.01	7.12–8.20
$ \Delta m_{31}^2 [10^{-3} \text{eV}^2]$	2.55	2.46 – 2.61	2.38 – 2.68	2.31 – 2.74
	2.43	2.37 – 2.50	2.29 – 2.58	2.21 – 2.64
$\sin^2 \theta_{12}$	0.320	0.303–0.336	0.29–0.35	0.27–0.37
$\sin^2 \theta_{23}$	0.613 (0.427) ^a	0.400-0.461 & 0.573–0.635	0.38–0.66	0.36–0.68
	0.600	0.569–0.626	0.39–0.65	0.37–0.67
$\sin^2 \theta_{13}$	0.0246	0.0218–0.0275	0.019–0.030	0.017–0.033
	0.0250	0.0223–0.0276	0.020–0.030	
δ	0.80π	$0 - 2\pi$	$0 - 2\pi$	$0 - 2\pi$
	-0.03π			

dirac mass term require a **very small yukawa...**

¿small masses = new physics?

Neutrino mass models

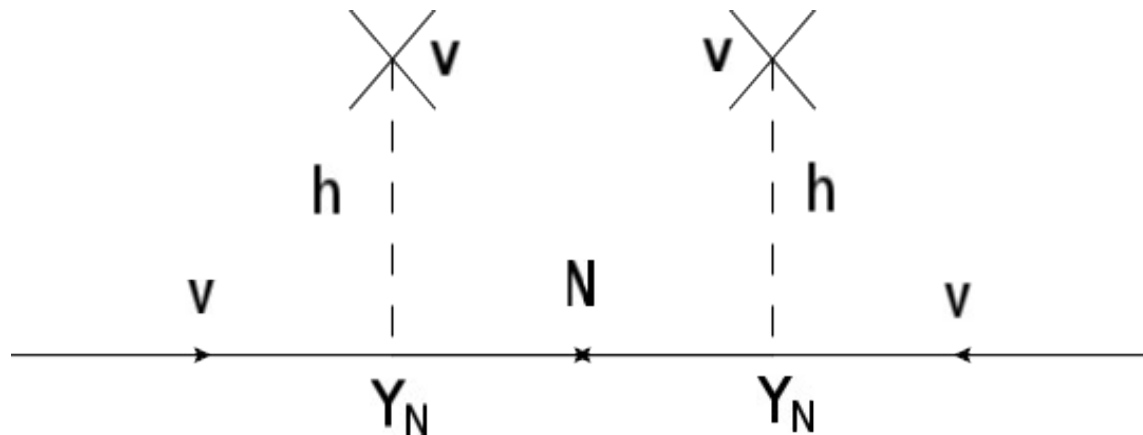
Effective dim 5 operator: Weinberg Operator $(\mathcal{O}_{ij}) = \frac{1}{\Lambda} L_{iL}^c \tilde{\phi}^* \tilde{\phi}^\dagger L_{jL}$
 where $L_i = (\nu_e, e), (\nu_\mu, \mu), (\nu_\tau, \tau)$;

	Field	Spin	SU(2)	Y
Type 1	N	1/2	1	0
Type 2	Δ	0	3	-2
Type 3	Σ	1/2	3	0

Neutrino mass models

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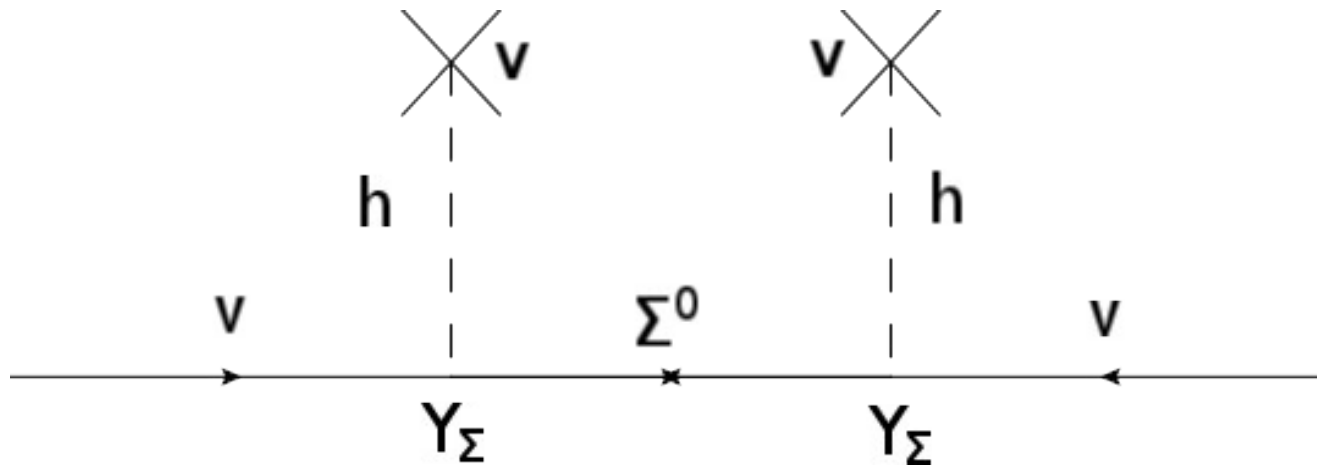
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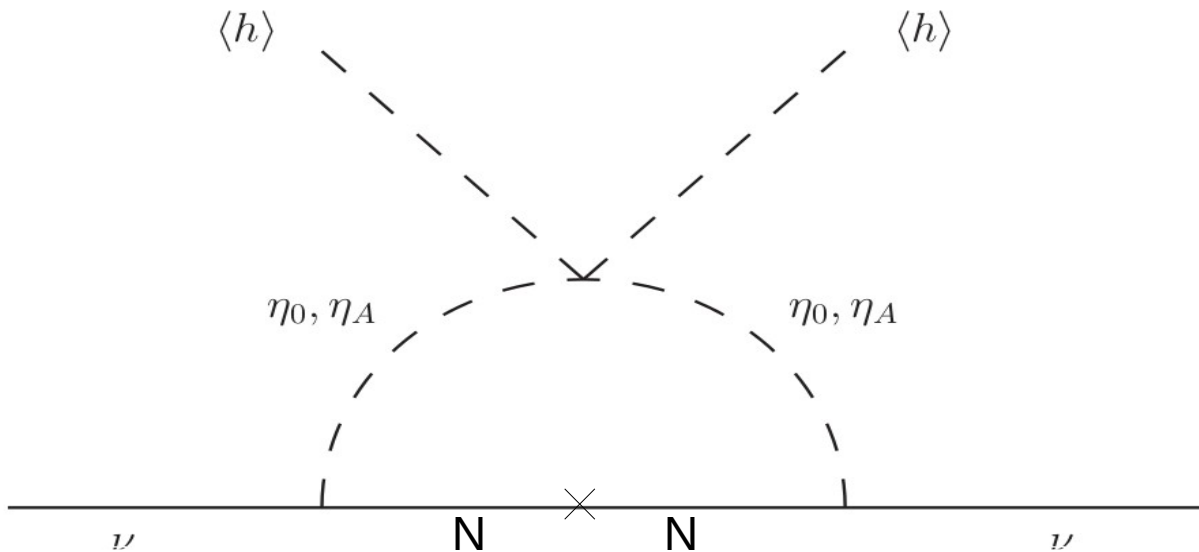
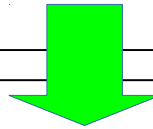
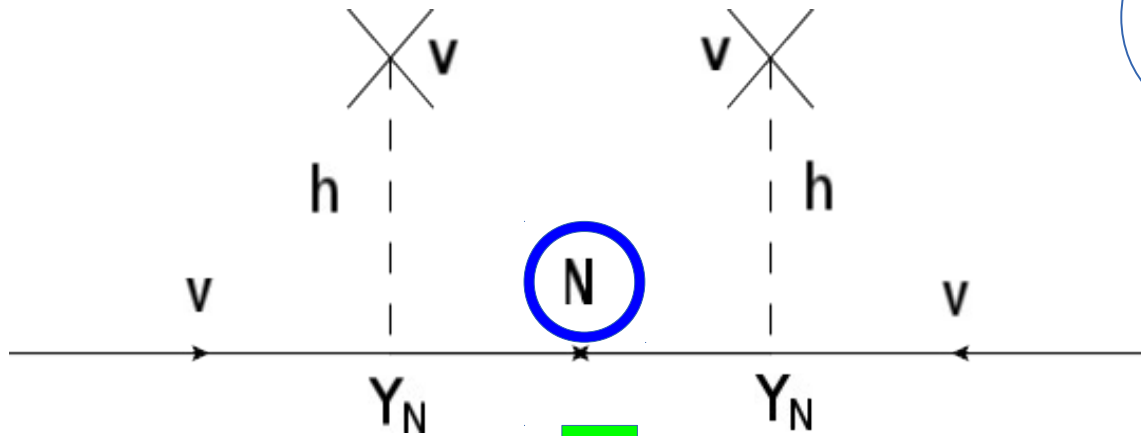
	Field	Spin	SU(2)	Y
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It was then realised that based on the same matter content, neutrino masses could arise at **loop** level, providing an **interesting link** between Dark Matter and Neutrino masses generating mechanism.

Neutrino mass models

As an example, Radiative Type 1

Z2



Scalar:

-lot of freedom

Fermion:

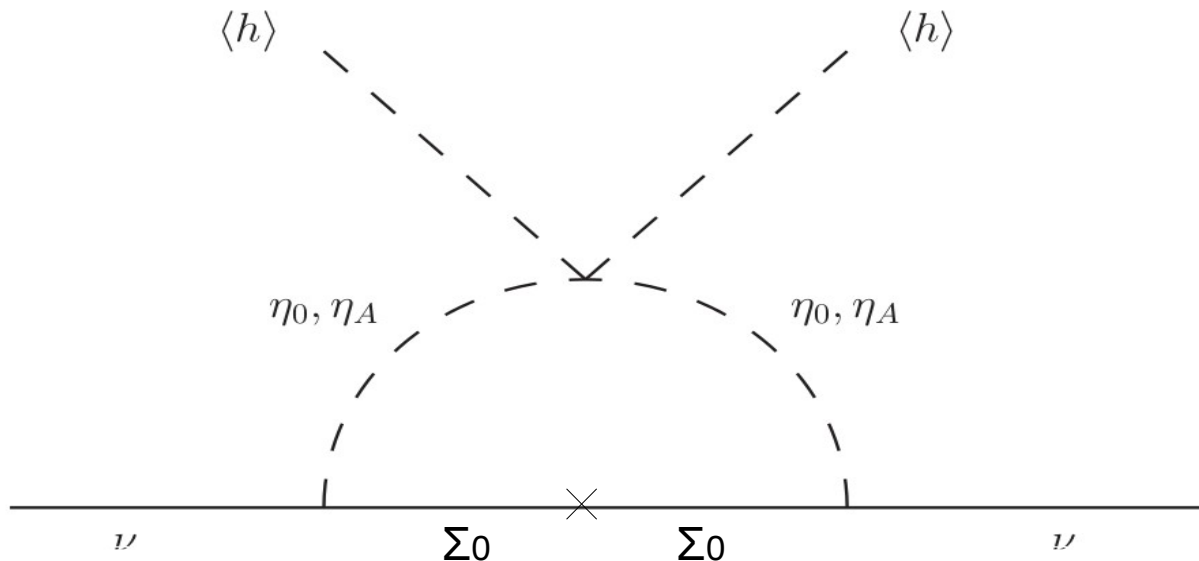
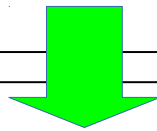
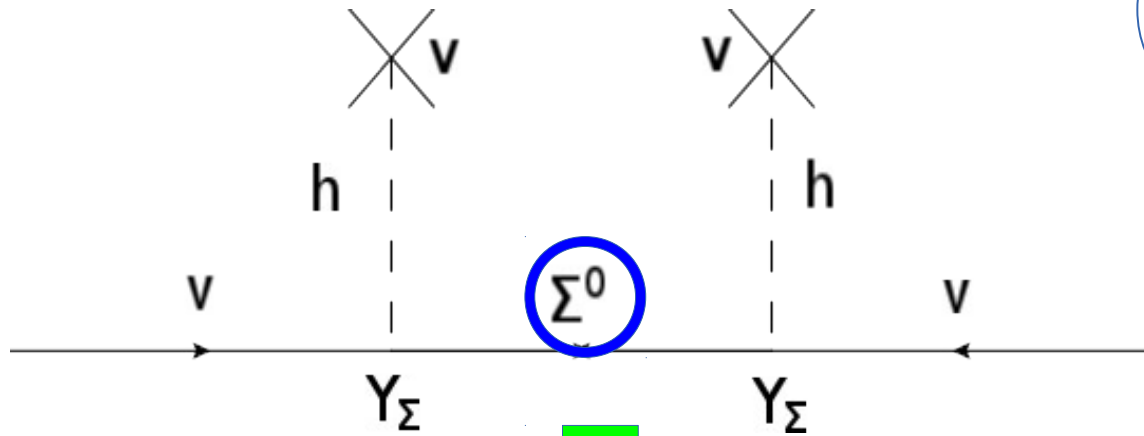
-easy light DM

-poor phenomenology

Neutrino mass models

And Radiative Type 3 case

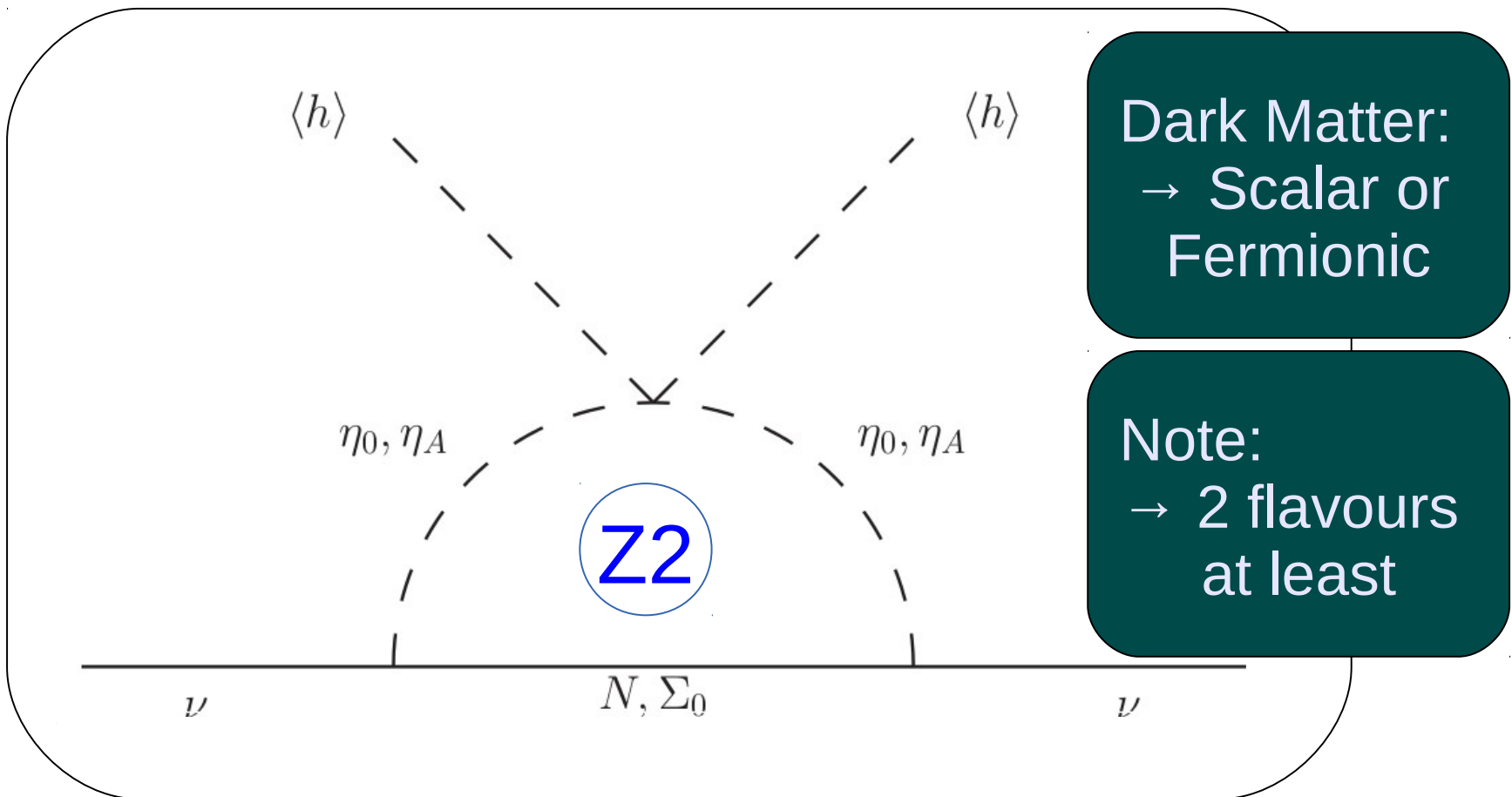
Z2



Scalar:
-lot of freedom

Fermion:
-rich phenomenology
-very constrained
DM mass

Neutrino mass models



We would like to join de **advantatges** of both scenario:

- >Light DM for the singlet
- >Rich phenomenology

The Model

M. Hirsch, R.A. Lineros, S. Morisi, J. Palacio, N. Rojas, J.W.F. Valle.
10.1007/JHEP10(2013)149

Z₂

Fields	Standard Model			Fermions		Scalars	
	L	e	ϕ	Σ	N	η	Ω
$SU(2)$	2	1	2	3	1	2	3
Y	-1	-2	1	0	0	1	0
Z_2	+	+	+	-	-	-	+

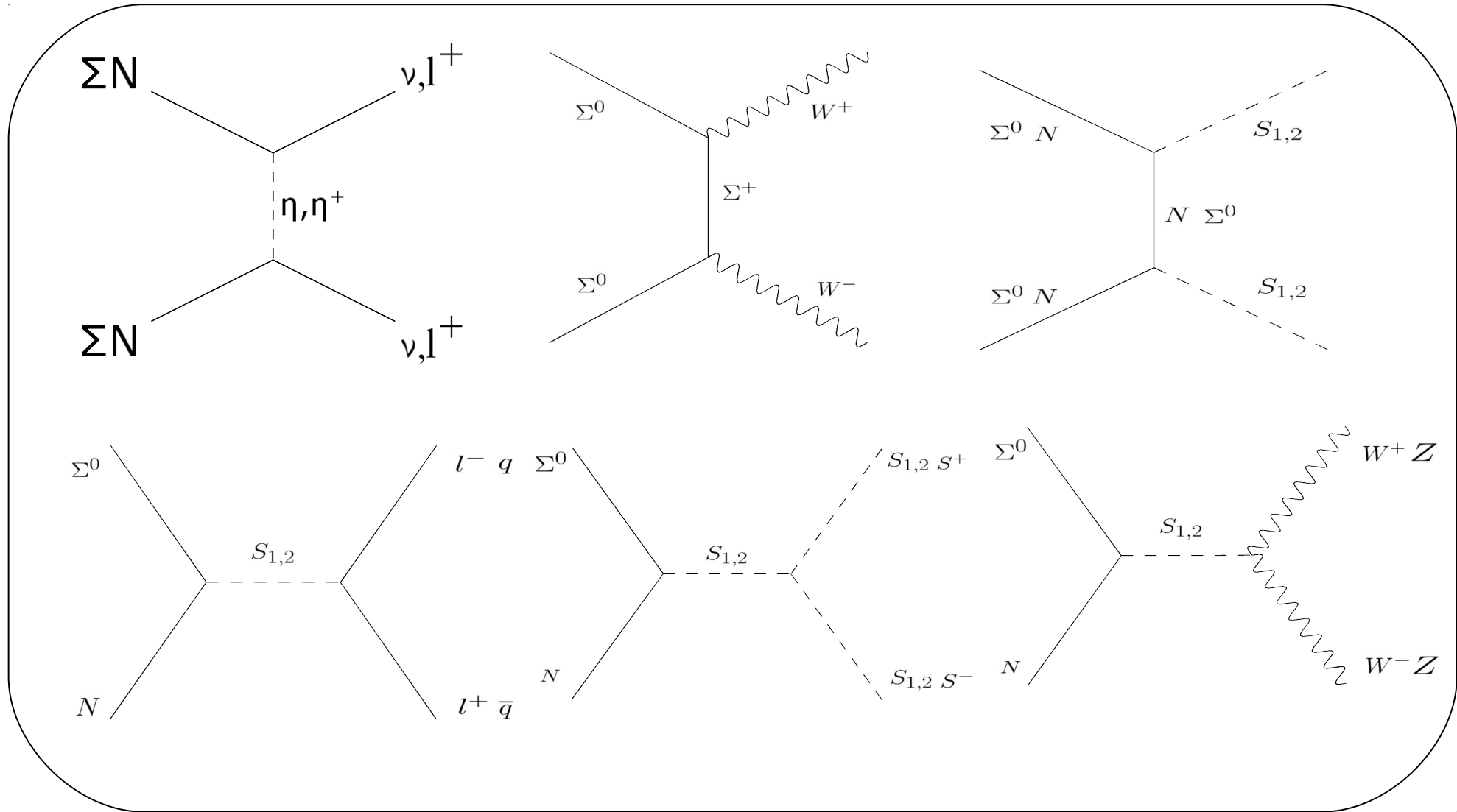
$$\mathcal{L} = -Y_{\alpha\beta} \bar{L}_\alpha e_\beta \phi - Y_{\Sigma\alpha} \bar{L}_\alpha C \Sigma^\dagger \tilde{\eta} - \frac{1}{4} M_\Sigma \text{Tr} [\bar{\Sigma}^c \Sigma] +$$

$$-Y_\Omega \text{Tr} [\bar{\Sigma} \Omega] N - Y_{N\alpha} \bar{L}_\alpha \tilde{\eta} N - \frac{1}{2} M_N \bar{N}^c N + h.c.$$

where $\alpha, \beta = 1, 2, 3$;

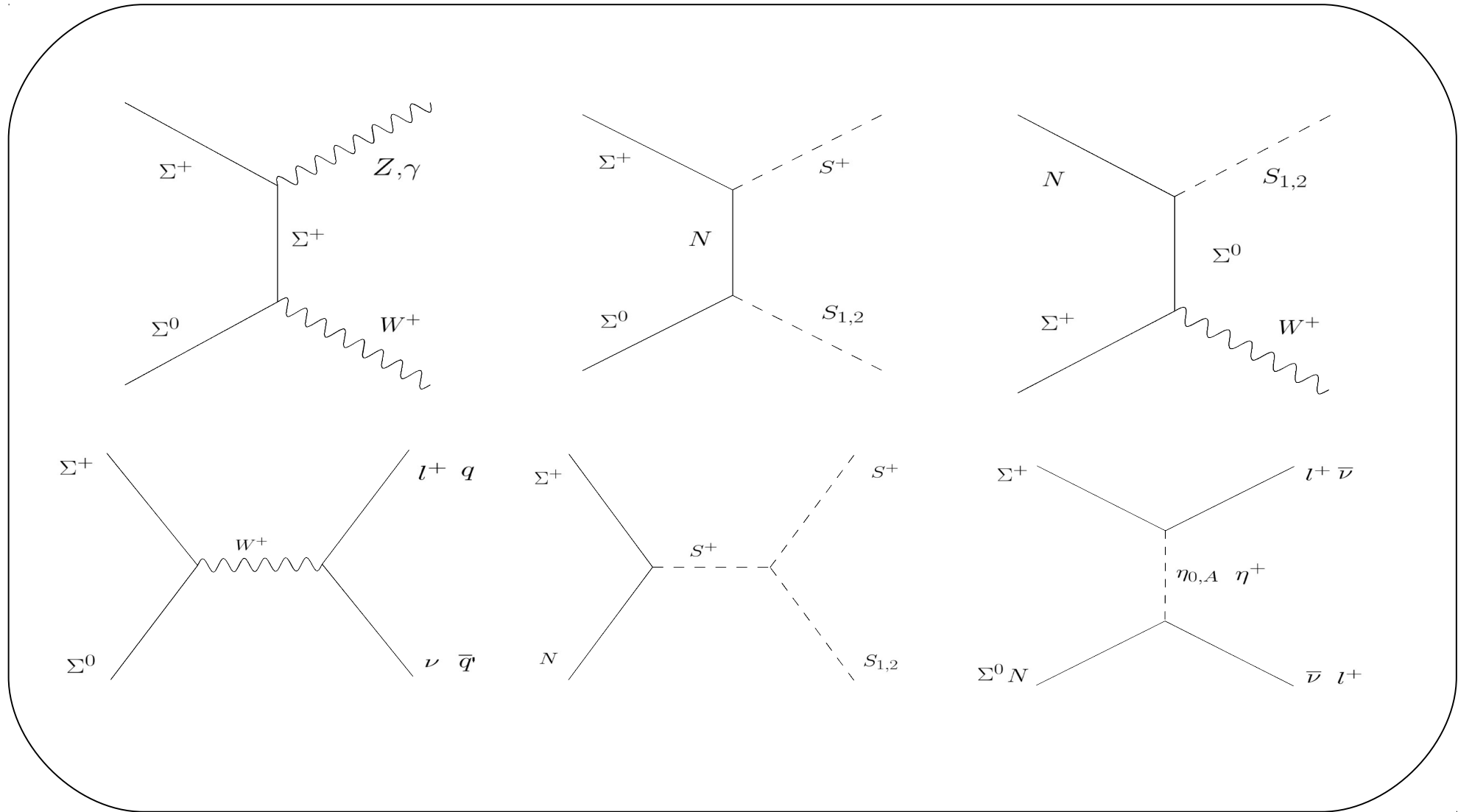
The Model

(Co-)Annihilation diagrams



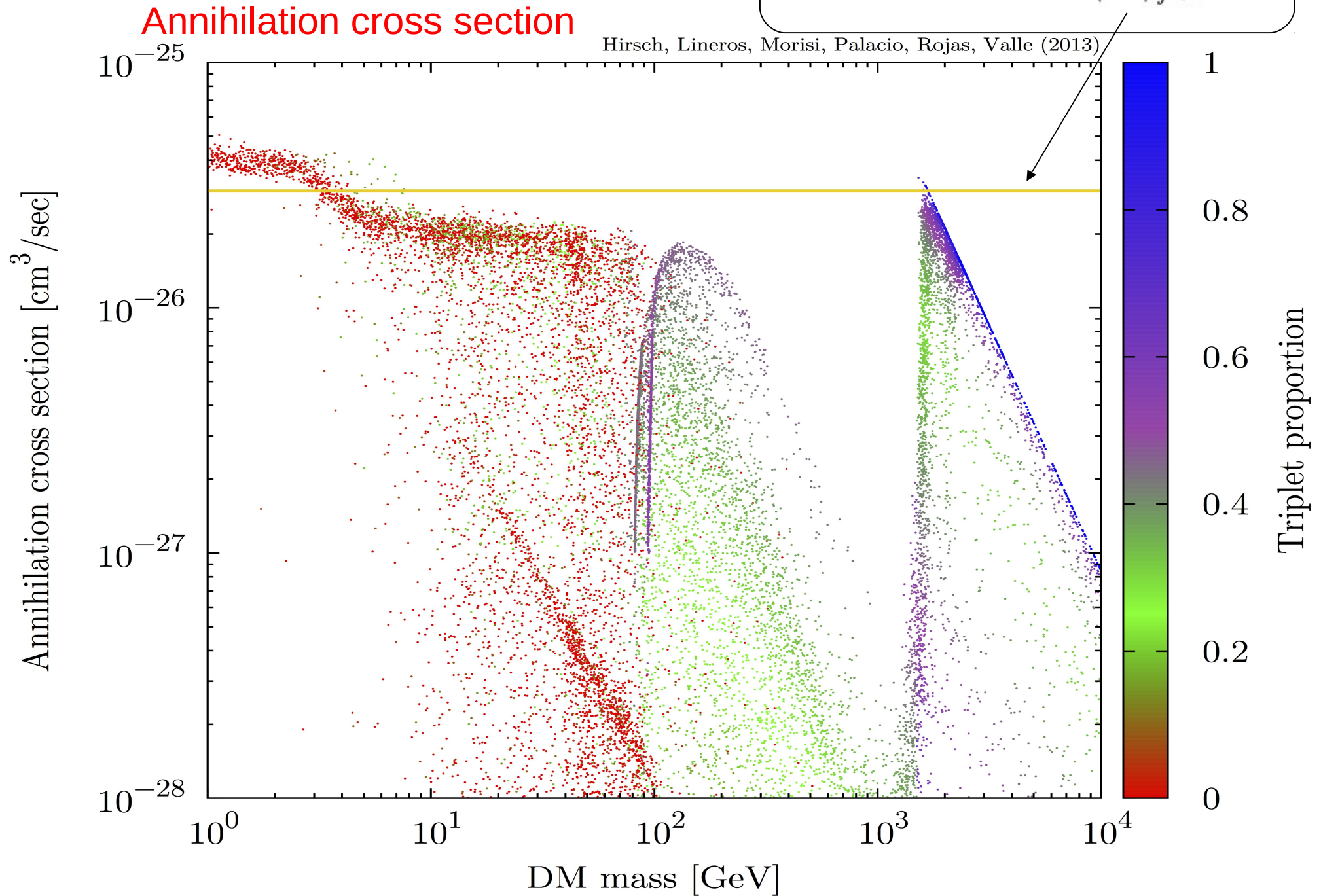
The Model

Charged co-Annihilation diagrams



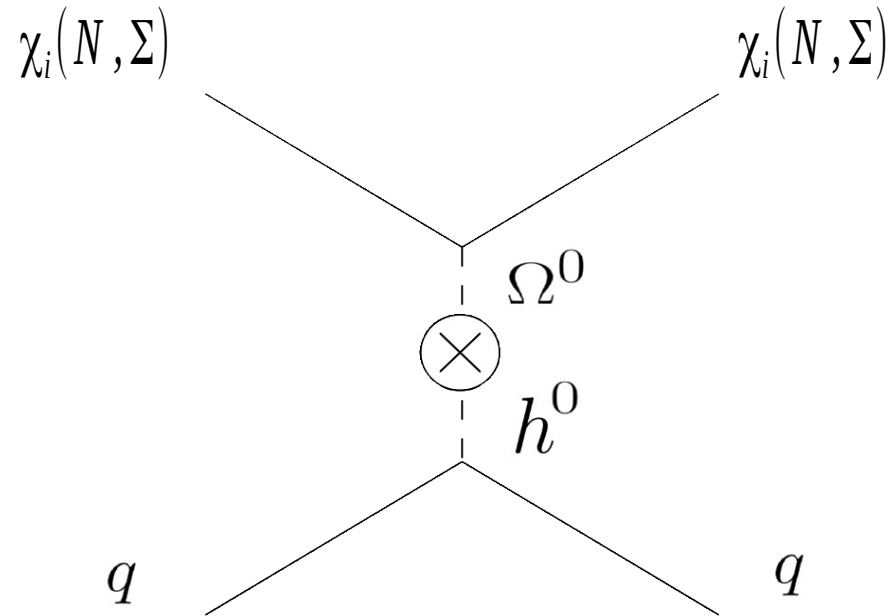
The Model

$$\Omega_{CDM}h^2 \approx 0.1 \frac{3 \cdot 10^{-26} \text{ cm}^3 \text{ s}^{-1}}{\langle \sigma v \rangle_{f.o.}}$$



The Model

Direct detection



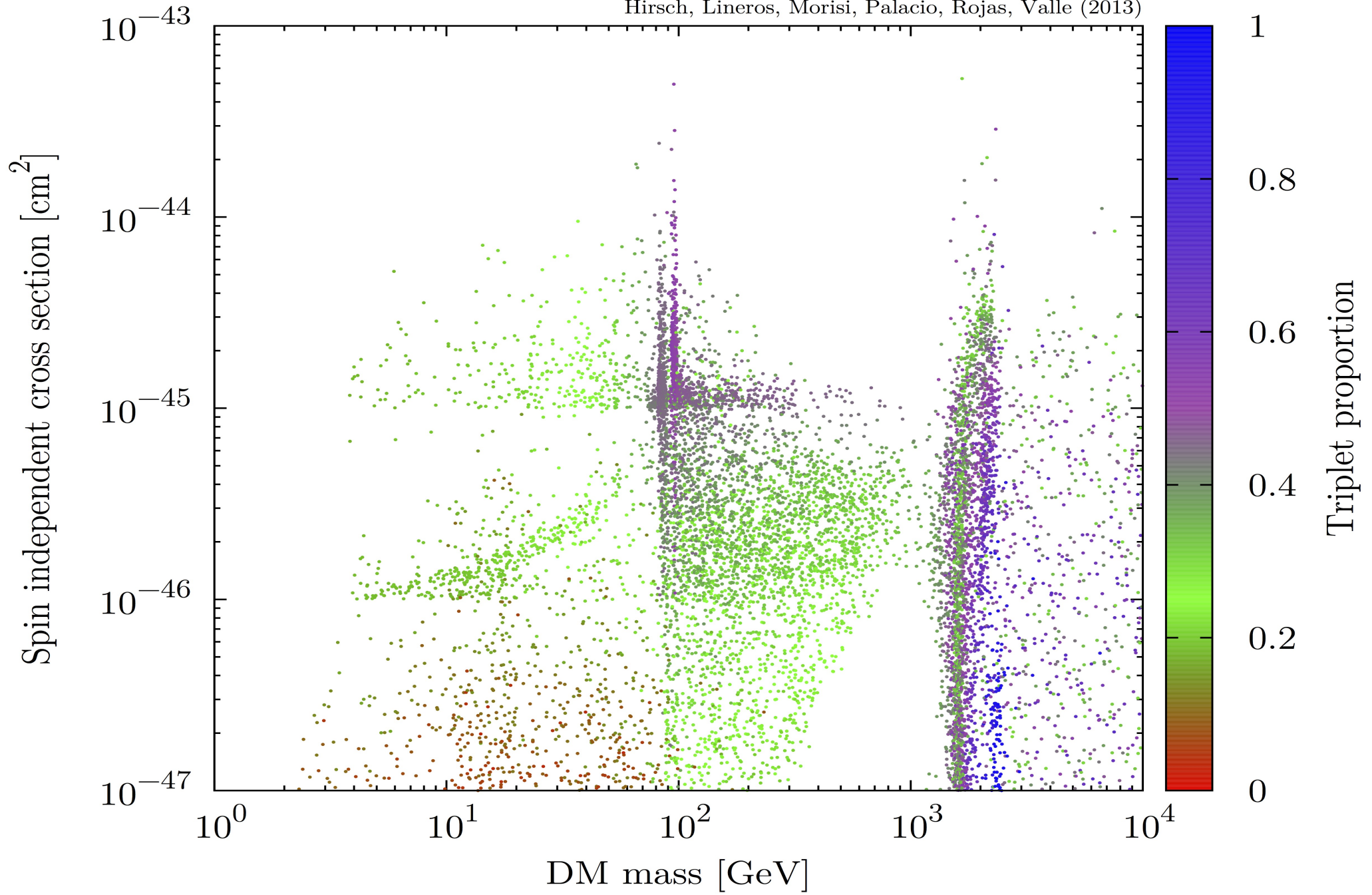
Z2

$$\begin{aligned}
 V = & -m_1^2 \phi^\dagger \phi + m_2^2 \eta^\dagger \eta + \frac{\lambda_1}{2} (\phi^\dagger \phi)^2 + \frac{\lambda_2}{2} (\eta^\dagger \eta)^2 + \lambda_3 (\phi^\dagger \phi) (\eta^\dagger \eta) \\
 & + \lambda_4 (\phi^\dagger \eta) (\eta^\dagger \phi) + \frac{\lambda_5}{2} (\phi^\dagger \eta)^2 + h.c. - \frac{M_\Omega^2}{4} \text{Tr} (\Omega^\dagger \Omega) + (\mu_1 \phi^\dagger \Omega \phi + h.c.) \\
 & + \lambda_1^\Omega \phi^\dagger \phi \text{Tr} (\Omega^\dagger \Omega) + \lambda_2^\Omega (\text{Tr} (\Omega^\dagger \Omega))^2 + \lambda_3^\Omega \text{Tr} ((\Omega^\dagger \Omega)^2) + \lambda_4^\Omega (\phi^\dagger \Omega) (\Omega^\dagger \phi) \\
 & + (\mu_2 \eta^\dagger \Omega \eta + h.c.) + \lambda_1^\eta \eta^\dagger \eta \text{Tr} (\Omega^\dagger \Omega) + \lambda_4^\eta (\eta^\dagger \Omega) (\Omega^\dagger \eta) .
 \end{aligned}$$

The Model

Spin independent cross section

Hirsch, Lineros, Morisi, Palacio, Rojas, Valle (2013)



Conclusions

Dark Matter and neutrino oscillations are the most robust evidence of physics beyond the Standard Model

We linked both phenomenas in this model: Neutrino mass-generating mechanism also stabilizes the Dark Matter.

The mixture scenario, Ω , gain the nice thinks of both pure Models: light DM with a rich phenomenology

The same mechanism that produces the fermion mixing also predicts a high interaction with quarks

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Dark Matter and neutrino oscillations are the most robust evidence of physics beyond the Standard Model

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Thanks

Back-up slides

The Scan

We used **micrOMEGAs** to do a parameter scan

G. Bélanger, F. Boudjema, A. Pukhov, A. Semenov, arXiv:1305.0237 [hep-ph]

Constraints: - Ω is an triplet of SU(2)

$$M_W = \frac{g}{2} \sqrt{v_h^2 + v_\Omega^2}$$



$$V_\Omega < 7\text{GeV}$$

- searches of new physics

Parameter	Range
M_N (GeV)	$1 - 10^5$
M_Σ (GeV)	$100 - 10^5$
m_{η^\pm} (GeV)	$100 - 10^5$
m_{η^0} (GeV)	$1 - 10^5$
M_\pm (GeV)	$100 - 10^4$
$ \lambda_i $	$10^{-4} - 1$
$ Y_i $	$10^{-4} - 1$

Scalar sector

$$\mathcal{M}_s^2 = \begin{pmatrix} \lambda_1 v_h^2 + \frac{t_h}{v_h} & -2\mu_1 v_h + 4v_h v_\Omega \left(\lambda_1^\Omega + \frac{\lambda_4^\Omega}{2} \right) \\ -2\mu_1 v_h + 4v_h v_\Omega \left(\lambda_1^\Omega + \frac{\lambda_4^\Omega}{2} \right) & \frac{\mu_1 v_h^2}{v_\Omega} + 16v_\Omega^2 (2\lambda_2^\Omega + \lambda_3^\Omega) + \frac{t_\Omega}{v_\Omega} \end{pmatrix}$$

$$M_{S1}^2 = v_h^2 \lambda_1 \cos^2(\theta_0) + 4v_h [-v_\Omega (2\lambda_1^\Omega + \lambda_4^\Omega) + \mu_1] \cos^2(\theta_0) \sin^2(\theta_0) + [16v_\Omega^2 (2\lambda_2^\Omega + \lambda_3^\Omega) + \mu_1 v_h^2 / v_\Omega] \sin^2(\theta_0)^2$$

$$M_{S2}^2 = v_h^2 \lambda_1 \sin^2(\theta_0) + 4v_h [v_\Omega (2\lambda_1^\Omega + \lambda_4^\Omega) - \mu_1] \cos^2(\theta_0) \sin^2(\theta_0) + [16v_\Omega^2 (2\lambda_2^\Omega + \lambda_3^\Omega) + \mu_1 v_h^2 / v_\Omega] \cos^2(\theta_0)^2$$

$$\text{where } \tan(2\theta_0) = \frac{4v_h [v_\Omega (2\lambda_1^\Omega + \lambda_4^\Omega) - \mu_1]}{16v_\Omega^2 (2\lambda_2^\Omega + \lambda_3^\Omega) - v_h^2 (\lambda_1 - \mu_1 / v_\Omega)}$$

Scalar sector

-Z2

$$\begin{aligned}m_{\eta 0}^2 &= m_2^2 + \frac{1}{2}(\lambda_3 + \lambda_4 + \lambda_5)v_h^2 + (2\lambda_1^\eta + \lambda_4^\eta)v_\Omega^2 - 2\mu_2 v_\Omega, \\m_{\eta A}^2 &= m_2^2 + \frac{1}{2}(\lambda_3 + \lambda_4 - \lambda_5)v_h^2 + (2\lambda_1^\eta + \lambda_4^\eta)v_\Omega^2 - 2\mu_2 v_\Omega, \\m_{\eta \pm}^2 &= m_2^2 + \frac{1}{2}\lambda_3 v_h^2 + 2\mu_2 v_\Omega + (2\lambda_1^\eta + \lambda_4^\eta)v_\Omega^2.\end{aligned}$$

λ_5 plays an important
role in v masses

The Model

$$\mathcal{L} = -Y_{\alpha\beta} \bar{L}_\alpha e_\beta \phi - Y_{\Sigma\alpha} \bar{L}_\alpha^c \Sigma^\dagger \tilde{\eta} - \frac{1}{4} M_\Sigma \text{Tr} [\bar{\Sigma}^c \Sigma] +$$

$$- Y_\Omega \text{Tr} [\bar{\Sigma} \Omega] N - Y_{N\alpha} \bar{L}_\alpha \tilde{\eta} N - \frac{1}{2} M_N \bar{N}^c N + h.c.$$

where $\alpha, \beta = 1, 2, 3$;

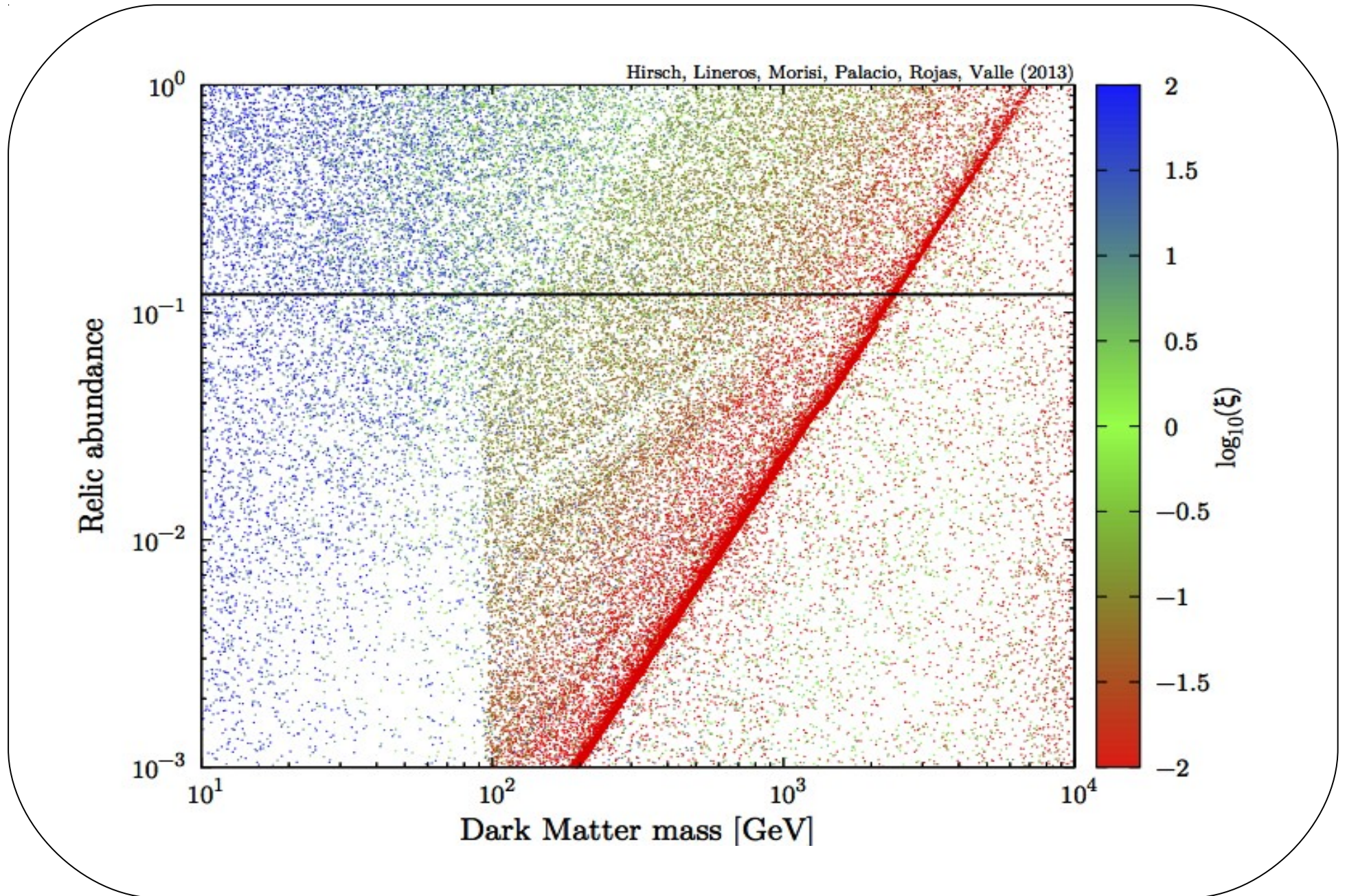
$$M_\chi = \begin{pmatrix} M_\Sigma & 2Y_\Omega v_\Omega \\ 2Y_\Omega v_\Omega & M_N \end{pmatrix}$$

$$M_{\chi 1} = \sin(\theta_f)^2 M_N + \cos(\theta_f)^2 M_\Sigma - 2v_\Omega Y_\Omega \cos(\theta_f) \sin(\theta_f)$$

$$M_{\chi 2} = \cos(\theta_f)^2 M_N + \sin(\theta_f)^2 M_\Sigma + 2v_\Omega Y_\Omega \cos(\theta_f) \sin(\theta_f)$$

where $\tan(2\theta_f) = \frac{-4Y_\omega v_\Omega}{M_\Sigma - M_N}$

Scan constrains



Scan constrains

